





# Chromo-Weibel instabilities in an expanding Quark-Gluon Plasma \*

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<sup>\*&</sup>quot; Dans la vie, rien n'est à craindre, tout est à comprendre." Marie Curie



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Hard Expanding
Loops (HEL)
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Stages of the Little Big Bang Momentum Anisotropy Weibel instabilities Scales QGP Hard (Thermal) Loops -Boltzmann -Vlasov Bjorken expansion

Plasma Instabilities

# Hard Expanding Loops (HEL)

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# Stages of the Little Big Bang

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[Gelis 2010] Illustration of the stages of a heavy ion collision.



### Momentum Anisotropy





# Weibel instabilities



[Strickland 2006]: Illustration of the mechanism of filamentation instabilities.



Stages of the Little Big Bang Momentum Anisotropy

#### Weibel instabilities

#### Scales QGP

Hard (Thermal) Loops -Boltzmann -Vlasov Bjorken expansion

Plasma Instabilities

# Scales of weakly coupled QGP

 $\blacksquare$  *T*: energy of hard particles

# ■ *gT*: thermal masses, Debye screening mass, Landau damping, plasma instabilities [Mrowczynski 1988, 1993, ..]

 $g^2T$ : magnetic confinement, color relaxation, rate for small angle scattering

 $\blacksquare g^4T$ : rate for large angle scattering,  $\eta^{-1}T^4$ 



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Hard (Thermal) Loops -

Boltzmann -

Vlasov

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# Hard (Thermal) Loops - Boltzmann - Vlasov

Assuming free streaming, one solves the gauge covariant Boltzmann-Vlasov equation

$$v \cdot D\partial f_a(\mathbf{p}, \mathbf{x}, t) = g v_\mu F_a^{\mu\nu} \partial_\nu^{(p)} f_0(\mathbf{p}, \mathbf{x}, t)$$
(1)



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coupled to Yang-Mills equation

$$D_{\mu}F_{a}^{\mu\nu} = j_{a}^{\nu} = g \int \frac{d^{3}p}{(2\pi)^{3}} \frac{p^{\mu}}{2p^{0}} \delta f_{a}(\mathbf{p}, \mathbf{x}, t)$$
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### in the HTL approximation

$$gA_{\mu} \ll |\mathbf{p}_{hard}|,$$
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Bjorken expansion

Plasma Instabilities

Vlasov

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#### in the HTL approximation

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with the Romatschke, Strickland background distribution function

$$f_0(p_\perp, \tilde{p}_\eta) = f_{\rm iso}\left([\mathbf{p}^2 + \xi(\tau)(\mathbf{p} \cdot \hat{\mathbf{n}})^2]/p_{\rm hard}^2(\tau)\right).$$
(4)



Stages of the Little Big Bang Momentum Anisotropy

Weibel instabilities

 ${\sf Scales} \ {\sf QGP}$ 

Hard (Thermal)

Loops -

Boltzmann -

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Bjorken expansion

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### **Bjorken expansion**





Loops (HEL) Stages of the Little Big Bang Momentum Anisotropy

Scales QGP

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Vlasov

Hard (Thermal)

### **Bjorken expansion**



Bjorken expansion

Weibel instabilities

Plasma Instabilities It is convenient to switch to comoving coordinates

 $t = \tau \cosh \eta$ ,  $\tau = \sqrt{t^2 - z^2}$ ,  $z = \tau \sinh \eta$ ,  $\eta = \operatorname{arctanh} \frac{z}{t}$ , (5)

#### with the corresponding metric

$$ds^{2} = d\tau^{2} - d\mathbf{x}_{\perp}^{2} - \tau^{2} d\eta^{2} \,. \tag{6}$$



Plasma

Instabilities HEL 3D+3V Check Expanding 3D+3V Abelian plasma Expanding 3D+3V non-Abelian plasma Conclusions

# **Plasma Instabilities**

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#### **Plasma Instabilities**

HEL 3D+3V Check Expanding 3D+3V Abelian plasma Expanding 3D+3V non-Abelian plasma Conclusions



HEL 3D+3V Check





[Attems, Rebhan, Strickland arXiv:1111.XXXX] Preliminary runs from the HEL 3d IC "stable nodes".



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Conclusions

Expanding 3D+3V Abelian plasma



[Attems, Rebhan, Strickland arXiv:1111.XXXX] Preliminary runs from the HEL 3d IC "transverse current".



# Expanding 3D+3V Abelian plasma

Loops (HEL) Plasma Instabilities HEL 3D+3V Check Expanding 3D+3V Abelian plasma Expanding 3D+3V non-Abelian plasma Conclusions



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Preliminary runs from the HEL 3d IC "transverse current".



### Expanding 3D+3V non-Abelian plasma





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Conclusions

#### Hard Expanding Loops (HEL)

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Conclusions

Non-abelian plasma instabilities play a significant role in a weakly coupled Quark Gluon Plasma at high T.

Chromo-Weibel instabilities are an important candidate process accelerating isotropization and thermalization of the Quark Gluon Plasma fireball maybe already at LHC.

The previous 1D HEL code has been extended to full 3D+3V parallel MPI-code with improved "transverse current" initial conditions. This significantly reduces the onset time of chromo-Weibel instabilities in the expanding setup.

Final results are being computed on the Vienna Scientific Cluster.



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# **Backup - Equation of motions**

### Yang-Mills equations

$$\tau^{-1}\partial_{\tau}\Pi_{i} = j^{i} - D_{j}F^{ji} - D_{\eta}F^{\eta i}, \qquad (7)$$
  
$$\tau\partial_{\tau}\Pi^{\eta} = j_{\eta} - D_{i}F^{i}{}_{\eta}. \qquad (8)$$

#### Gauss law constraint

$$\tau j^{\tau} = D_{\eta} \Pi^{\eta} - ig[A^i, \Pi_i], \qquad (9)$$

Canonical conjugate field momenta

$$\Pi^{i} \equiv \tau \partial_{\tau} A_{i}, \quad \Pi^{\eta} \equiv \frac{1}{\tau} \partial_{\tau} A_{\eta} .$$
 (10)



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[Rebhan, Strickland, Attems 2008] Fourier spectrum of the color-traced conjugate field momentum obtained from Abelian run with FGM initial conditions.



Plasma Instabilities HEL 3D+3V Check Expanding 3D+3V Abelian plasma Expanding 3D+3V non-Abelian plasma Conclusions Backup - Expanding 1D+3V non-Abelian plasma



[Rebhan, Strickland, Attems 2008] Fourier spectrum of the color-traced conjugate field momentum obtained from non-Abelian run with FGM initial conditions.



Plasma Instabilities HEL 3D+3V Check Expanding 3D+3V Abelian plasma

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# Backup - Expanding 1D+3V non-Abelian plasma



[Strickland 2008] Visualization of the space-time development of color correlations in a non-Abelian plasma instabilities in Bjorken expansion.



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#### $10^{5}$ $\left| \tilde{B}^{i}(\tau, \nu = 30) \right|$ $10^{4}$ ~; (ro) #0 $10^{3}$ TI (ro) = 0 $10^{2}$ $10^{1}$ $10^{0}$ 10<sup>-1</sup> $10^{-2}$ 20 30 5 10 50 100 150 200 1 $\tau/\tau_0$

[Rebhan, Steineder 2009] Influence of different initial conditions for a specific mode with  $\nu=30.$ 

# Backup - Unstable transverse modes

8th Vienna Central European Seminar