

Spectral Multiplicity of Selfadjoint Schrödinger Operators on Star-Graphs with Standard Interface Conditions

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Abstract. We analyze the singular spectrum of selfadjoint operators which arise from pasting a finite number of boundary relations with a standard interface condition. A model example for this situation is a Schrödinger operator on a star-shaped graph with continuity and Kirchhoff conditions at the interior vertex. We compute the multiplicity of the singular spectrum in terms of the spectral measures of the Weyl functions associated with the single (independently considered) boundary relations. This result is a generalization and refinement of a Theorem of I.S.Kac.

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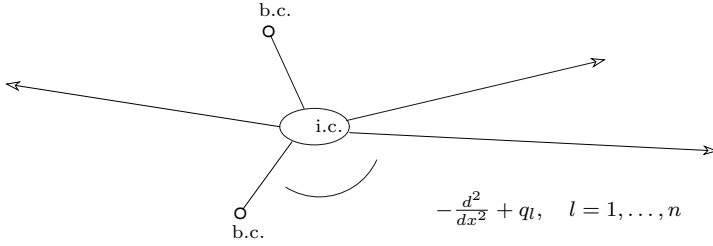
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1. Introduction

In the present paper we undertake an analysis of the singular spectrum of selfadjoint operators which are constructed by pasting a finite number of boundary triples (relations) by means of a standard interface condition.

For the purpose of explaining our results without having to introduce too much terminology, we consider a model example: A Schrödinger operator on a star-graph. Consider a star-shaped graph having finitely many edges, say, E_1, \dots, E_n . We think of the edges as (finite or infinite) intervals $E_l = [0, e_l)$, where the endpoint 0 corresponds to the interior vertex. A selfadjoint operator can be constructed from the following data:

- (1) On each edge E_l , a real-valued potential $q_l \in L^1_{\text{loc}}([0, e_l))$.
- (2) Boundary conditions at outer vertices e_l , if the potential q_l at e_l is regular or in Weyl's limit circle case.
- (3) An interface condition at the interior vertex.



The operator A one can associate with this data acts in the space $H := \prod_{l=1}^n L_2(0, e_l)$ as

$$A \begin{pmatrix} u_1 \\ \vdots \\ u_n \end{pmatrix} := \begin{pmatrix} -u_1'' \\ \vdots \\ -u_n'' \end{pmatrix} + \begin{pmatrix} q_1 u_1 \\ \vdots \\ q_n u_n \end{pmatrix}, \quad (1.1)$$

on the domain

$$\text{dom } A := \left\{ (u_1, \dots, u_n) \in \prod_{l=1}^n L_2(0, e_l) : \begin{array}{l} u_l, u_l' \text{ are absolutely continuous, } -u_l'' + q_l u_l \in L_2(0, e_l), \\ u_l \text{ satisfies the boundary condition at outer vertex (if present),} \\ u_1, \dots, u_n \text{ satisfy the interface condition at the inner vertex} \end{array} \right\}. \quad (1.2)$$

A frequently used interface condition, sometimes called the “standard condition”, is

$$u_1(0) = \dots = u_n(0) \quad \text{and} \quad \sum_{l=1}^n u_l'(0) = 0. \quad (1.3)$$

In the case “ $n = 2$ ” the condition (1.3) arises when investigating a whole-line Schrödinger operator with the classical method of Titchmarsh and Kodaira.

The task now is to describe the projection-valued spectral measure E of A in terms of the scalar spectral measures μ_l of the non-interacting operators $A_l, l = 1, \dots, n$, which are defined by the potentials q_l on the edges E_l independently (imposing Dirichlet boundary conditions at the inner vertex 0 for each of them).

A precise description of the absolutely continuous part¹ E_{ac} of E , including computation of its spectral multiplicity, is readily available. It states that E_{ac} is equivalent (in the sense of mutual absolute continuity) to the

¹Notice that the notions of absolute continuity and singularity of measures make sense also if the involved measures have different ranges. Moreover, Lebesgue decompositions of a projection-valued measure with respect to a scalar measure (in this case, the Lebesgue measure) exist.

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sum $\mu_{1,ac} + \dots + \mu_{n,ac}$ of the absolutely continuous parts of the measures μ_l . Moreover, informally speaking, the local spectral multiplicity corresponding to E_{ac} is equal to the number of overlaps of $\mu_{1,ac}, \dots, \mu_{n,ac}$ (for a precise formulation see Theorem 1.1). These facts follow from [15, Theorem 6.6], a result which can be viewed as a higher-dimensional analogue (and refinement) of one half of Aronszajn–Donoghue theory for rank one perturbations. Namely, of the part which asserts stability of absolutely continuous spectrum, cf. [2, Theorem 1], [11, Theorems 2 and 6].² Another approach proceeds via scattering theory and uses a modification of the Kato–Rosenblum theorem [5, 22], see also [38, Theorem 1.9]. Since the operator A is a finite dimensional perturbation (in the resolvent sense) of the operator $\bigoplus_{l=1}^n A_l$, wave operators exist and are complete, which in turn means that the absolutely continuous parts of these operators are unitarily equivalent.

In the present paper, we describe the singular part E_s of E , including a formula for spectral multiplicity. Our main result is Theorem 1.2 below (where we provide the formulation for the Schrödinger case; for the general situation see Theorem 4.1). Again speaking informally, it says that:

- (I) One part of E_s appears where at least two of the singular parts $\mu_{l,s}$ of the measures μ_l overlap. Where only one singular part $\mu_{l,s}$ is present, the spectrum disappears.
- (II) For the part of E_s described in (I), the local spectral multiplicity is equal to the number of overlaps of $\mu_{1,s}, \dots, \mu_{n,s}$ minus 1. In particular, the multiplicity cannot exceed $n - 1$.
- (III) The remaining part of E_s is mutually singular to each of the spectral measures μ_l and has multiplicity 1.

This theorem is a generalization and refinement of a theorem given by Kac [20].³ He considered the case of two edges and showed that the spectral multiplicity of the singular part E_s is always 1. Kac’ Theorem corresponds to the upper bound for multiplicity in (II) and simplicity of spectrum in (III). Realizing a change of boundary condition of a half-line operator as an interface condition with an “artificial second edge”, we can also reobtain the half of Aronszajn–Donoghue theory which asserts disjointness of singular spectra for different boundary conditions, see again [2, 11], or [15, Theorem 3.2, (iv)]. This corresponds to the fact in (I) that, if only one spectrum is present, it disappears.

By using the abstract framework of boundary relations, instead of just discussing a Schrödinger operator on a star-graph, we achieve a slight generalization and a significant increase of flexibility in applications (various kinds of operators, not necessarily being differential operators, can be treated). This bonus comes without additional effort, since our proofs proceed via an

²See also [15, Theorem 3.2, (i)–(iii)] for a summary.

³Full proofs are provided in [21] (in Russian). An English translation of this paper is not available, however, the proof was reproduced by D.Gilbert in [14]: The operator-theoretic half of Kac’ theorem is [14, Theorem 5.1], the measure-theoretic half is [14, Theorem 5.5, (i)]. An interesting approach to Kac’ theorem was given recently by Simon [33] who proceeds via rank-one perturbations and uses Aronszajn–Donoghue theory.

analysis of Weyl functions and associated measures, and do not rely on the concrete form of the operators on edges.

The description of the absolutely continuous part E_{ac} is not specific for the geometry of a star-graph and/or the use of standard interface conditions: the mentioned result [15, Theorem 6.6] holds for arbitrary finite-rank perturbations. Contrasting this, the description of the singular part E_s given in (I)–(III) is specific for the particular situation. This is seen, for example, from some known formulas for the maximal multiplicity of an eigenvalue of a Schrödinger operator on a graph (not necessarily a star-graph). It turns out that this number depends on the geometry of the graph (rather than rank of the perturbation), see [24] and the references therein. Another good example is a theorem due to Howland, cf. [18, 2.Theorem]. There for a certain type of finite-rank perturbation a behavior is witnessed which is fully in opposite to (I)–(III). Also we should mention that, although the considered operator A is “only” a rank-one perturbation of the direct sum of the non-interacting operators A_l , classical perturbation theory does not give much information. For example, the Kato–Rosenblum theorem deals with absolutely continuous spectrum, or the theorem [36, Satz 10.18] on the ranges of spectral projections yields information only for isolated eigenvalues. The singular (continuous and possibly embedded) spectrum is much more instable, and its behavior is much harder to control.

Let us give a brief outline of the organization of the paper. In the second part of this introductory section, we explain the structure of the spectrum of A in some more detail (old and new results). Section 2 is of preparatory nature. We set up notation and collect some results from the literature concerning: spectral multiplicity, Borel measures, and Cauchy integrals. We compiled a concise yet exhaustive review of the required machinery and included it in the paper since we believe that some of these deep and rather recent theorems are maybe not widely known among operator theorists.

In Sect. 3, we provide some facts about boundary relations and the Titchmarsh–Kodaira formula. Some of these statements are well-known and readily available in the literature, for some others explicit references are not known to us and hence proofs are given. Most important, we define the main object of our studies: the pasting of boundary relations with standard interface conditions, cf. Definition 3.16, and compute its matrix valued Weyl function in terms of the Weyl functions of the single boundary relations, cf. Proposition 3.18. Moreover, we carry out the calculations required to determine the point spectrum, cf. Theorem 3.19. Though the case of eigenvalues is of course included in our main result, we find it worth to be formulated and proved independently; it serves as an elementary accessible, yet precise, model for the behavior of singular continuous spectrum.

Section 4 forms the core of the paper. In this section we formulate and prove our main result Theorem 4.1; the major task is to get control of the singular continuous (possibly embedded) part of the singular spectrum. The proof can be outlined as follows: We further divide the singular part E_s into two summands. Namely, setting $\mu = \sum_{l=1}^n \mu_l$, we decompose E_s into the sum of a measure which is absolutely continuous with respect to μ and one which

is singular with respect to μ . First, we show that on null sets of the measure μ only simple spectrum of A may appear, and this shows item (III). Second, we consider points having certain “good” properties regarding existence of derivatives of involved measures and pointwise asymptotics of their Poisson and Cauchy integrals. For such points the multiplicity of the spectrum can be calculated, and this shows items (I) and (II) on the set of “good” points. Finally, we show that this set of “good” points in fact supports the full singular part of μ , and thereby complete the proof of items (I) and (II).

The paper closes with two appendices. In the first appendix we provide some examples which show that all possibilities permitted by (I)–(III) indeed may occur. These are not difficult to obtain and are based on classical theory and some more recent results on concrete potentials on the half-line. This section will not hold many surprises for the specialist in the field; we include it to give a fuller picture. In the second appendix we show how to reobtain from our present results the classical theorem of Aronszajn and Donoghue on singular spectra associated with different boundary conditions. Moreover, we include a short discussion of some (a few) interface conditions different from the standard condition.

There occurs an obvious open problem: Is it true that also for other finite-rank perturbations the singular continuous spectrum behaves in the same way as the point spectrum (concerning its multiplicity)? In a very general setting, one may think of investigating arbitrary finite rank perturbations; optimally obtaining a full higher-dimensional analogue of Aronszajn–Donoghue theory for singular spectra. However, this is probably wishful thinking: Keeping in mind the difficulties which arise when considering eigenvalues in the case of standard (Kirchhoff) interface conditions on a graph with a somewhat more complicated geometric structure, already a thorough investigation of this situation seems a challenging task.

At present, the answer to whatever version of the above posed question is not at all clear. The computations we use in this paper are specific for the case “star-graph+standard interface conditions”. We plan to address this problem in future work.

Recently, after the present paper was completed, we learned about a potential different approach to this circle of questions. Namely, to appropriately construct reducing subspaces, use properties of the orthogonal projections between these subspaces, and employ classical Aronszajn–Donoghue theory. This idea is due to Vladimir Kapustin who works on implementing it. We are grateful to him for discussing his ideas with us.

Detailed Description of the Structure of $\sigma(A)$

Again, for the purpose of explaining, we consider a Schrödinger operator A on a star-graph which is given by the data (1)–(3).

A first, rough, insight into the structure of the spectrum is provided by the classical Titchmarsh–Kodaira formula. We may consider the operator A as a selfadjoint extension of the symmetry S whose domain is defined by requiring that $u_l(0) = u_l'(0) = 0, l = 1, \dots, n$. This symmetry is completely non-selfadjoint, and has defect index (n, n) . The spectral multiplicity of A

cannot exceed n : There exists an $n \times n$ -matrix valued measure Ω such that the operator A is unitarily equivalent to the operator of multiplication by the independent variable in the space $L_2(\mathbb{R}, \Omega)$. A measure Ω with this property can be constructed using Weyl theory. Since A is an extension of S , there exists a matrix-valued Weyl function $M(z)$ corresponding to A . The measure Ω in the Herglotz-integral representation of M has the required properties.

Since the spectral projection of the multiplication operator in $L_2(\mathbb{R}, \Omega)$ onto a Borel set Δ is the multiplication operator with the indicator function of Δ , it follows that Ω and E are mutually absolutely continuous. If we set $\rho := \text{tr} \Omega$, then it is easy to see that Ω and ρ are mutually absolutely continuous. We call the measure ρ from this construction a *scalar spectral measure* corresponding to the operator A (this measure is of course not unique).

The same procedure can be carried out for each of the operators A_l . For the operator A_l the defect index of the minimal operator is $(1, 1)$, and one gets a unitary equivalence to the multiplication operator the space $L_2(\mathbb{R}, \mu_l)$, where μ_l is the (now scalar) measure taken from the Herglotz-integral representation of the Weyl function associated with A_l .

Let $N_A(x)$ be the *spectral multiplicity function* of A which is defined ρ -a.e. The detailed definition of N_A requires some background; we recall it in §2, see (2.3). Moreover, set $\mu := \sum_{i=1}^n \mu_i$, and

$$r(x) := \#\{l \in \{1, \dots, n\} : D_\mu \mu_l(x) > 0\}. \tag{1.4}$$

Here $D_\mu \mu_l(x)$ denotes the Radon–Nikodym derivative of μ_l with respect to μ , and the function $r(x)$ is defined μ -a.e. Note that $\sum_{l=1}^n D_\mu \mu_l = 1$ and hence $r(x) \geq 1$ for μ -a.a. points $x \in \mathbb{R}$.

A complete description of the absolutely continuous part of the spectral measure E of A follows from [15, Theorem 6.6]. Notation: We use \sim to denote mutual absolute continuity of two measures.

Theorem 1.1. [15] *Let A be a Schrödinger operator on a star-graph given by the data (1)–(3) using the standard interface condition (1.3). Denote by E the projection valued spectral measure of A , let μ be the sum of the scalar spectral measures of the non-interacting operators A_l , and let E_{ac} and μ_{ac} be their absolutely continuous parts with respect to the Lebesgue measure. Moreover, let N_A be the spectral multiplicity function of A , and let $r(x)$ be as in (1.4). Then*

- (I) $E_{ac} \sim \mu_{ac}$.
- (II) $N_A(x) = r(x)$ for E_{ac} -a.a. points $x \in \mathbb{R}$.

The following complete description of the singular part of the spectral measure E of A is the main result of this paper (formulated for the Schrödinger case; the general statement is Theorem 4.1). Notation: If X is a Borel set, we write $\mathbb{1}_X \cdot \nu$ for the measure acting as $(\mathbb{1}_X \cdot \nu)(\Delta) = \nu(X \cap \Delta)$.

Theorem 1.2. *Let A be a Schrödinger operator on a star-graph given by the data (1)–(3) using the standard interface condition (1.3). Denote by E the projection-valued spectral measure of A , let μ be the sum of the scalar spectral measures of the non-interacting operators A_l , and let E_s and μ_s be their singular parts with respect to the Lebesgue measure. Let $E_{s,ac}$ and $E_{s,s}$ be the*

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absolutely continuous and singular parts of E_s with respect to μ . Moreover, let N_A be the spectral multiplicity function of A and (as in (1.4))

$$r(x) := \#\{l \in \{1, \dots, n\} : D_\mu \mu_l(x) > 0\}.$$

Then

- (I) $E_{s,ac} \sim \mathbb{1}_{X_{>1}} \cdot \mu_s$ where $X_{>1} := r^{-1}(\{2, \dots, n\})$.
- (II) $N_A(x) = r(x) - 1$ for $E_{s,ac}$ -a.a. points $x \in \mathbb{R}$.
- (III) $N_A(x) = 1$ for $E_{s,s}$ -a.a. points $x \in \mathbb{R}$.

Notice that the Radon-Nikodym derivatives $D_\mu \mu_l$ and the number r are defined μ -a.e. The functions N_A and r should be considered as representatives of the equivalence classes under different equivalence relations. However the equality in item (II) makes sense and holds true $E_{s,ac}$ -a.e. for any choice of such representatives, because the measure $E_{s,ac}$ is absolutely continuous with respect to both E and μ . In turn the set $X_{>1}$ is defined up to a μ -zero set, but the measure $\mathbb{1}_{X_{>1}} \cdot \mu_s$ is defined uniquely.

Finally, let us make explicit the behavior of the point spectrum.

Theorem 1.3. *Let A be a Schrödinger operator on a star-graph given by the data (1)–(3) using the standard interface condition (1.3). Denote by m_l the Weyl functions of the non-interacting operators A_l , and $r(x)$ be as in (1.4). Then a point $x \in \mathbb{R}$ belongs to $\sigma_p(A)$, if and only if one of the following alternatives takes place.*

(I/II) *The point x belongs to at least two of the point spectra $\sigma_p(A_l)$. In this case the multiplicity of the eigenvalue x is equal to*

$$\#\{l = 1, \dots, n : x \in \sigma_p(A_l)\} - 1.$$

(III) *The limits $m_l(x) := \lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon)$ all exist, are real, we have $\lim_{\varepsilon \downarrow 0} \frac{1}{i\varepsilon} (m_l(x + i\varepsilon) - m_l(x)) \in [0, \infty)$, and $\sum_{j=1}^n m_j(x) = 0$. In this case x is a simple eigenvalue.*

The connection of Theorem 1.3 with the general result Theorem 1.2 is made as follows: For a point belonging to the point spectrum of at least one of the operators A_l , we have $D_\mu \mu_k(x) > 0$ if and only if $x \in \sigma_p(A_k)$ ($k \in \{1, \dots, n\}$). Hence, for such points,

$$r(x) = \#\{l = 1, \dots, n : x \in \sigma_p(A_l)\}.$$

Moreover: If x is an eigenvalue of only one operator A_l it disappears. And the set of all points x which satisfy the conditions stated in (III) is μ -zero.

2. Preliminaries

2.1. Some Terminology Concerning Measures

First of all, let us fix some measure theoretic language. We denote by \mathcal{B} the σ -algebra of all Borel sets on \mathbb{R} . All measures ν are understood to be Borel measures, and this includes the requirement that compact sets have finite measure. Whenever writing “ $\nu(X)$ ”, this implicitly includes that $X \in \mathcal{B}$. If

we speak of a *positive measure* ν , this measure needs not to be finite. For a *complex measure* ν , we denote by $|\nu|$ its *total variation*, and this is always a finite positive measure. If a complex measure takes only real values, we also speak of a *real measure*.

In some places we have to deal with sets which are not necessarily Borel sets, and with functions which are not necessarily Borel measurable. We say that X is a ν -zero set, if $X \subseteq \mathbb{R}$ and there exists a Borel set $X' \supseteq X$ such that $\nu(X') = 0$. We say that a set $X \subseteq \mathbb{R}$ is ν -full, if its complement is ν -zero. A property is said to hold ν -a.e. or for ν -a.a. points x , if the set of all points where it holds is ν -full. Moreover, we say that a partially defined function f is ν -measurable, if its domain is ν -full and there exists a Borel measurable function which coincides ν -a.e. with f . Integrals $\int_{\mathbb{R}} f d\nu$ of ν -measurable functions f are defined accordingly.

Of course, such terminology could be avoided by considering ν as a measure on the completion of the σ -algebra \mathcal{B} with respect to ν , and understanding measurability with respect to this larger σ -algebra. However, then one has to work with different σ -algebras for different measures, and this would make things technically laborious.

When ν is a (positive or complex) measure, and σ is a positive measure, we say that ν is *absolutely continuous* with respect to σ (and write $\nu \ll \sigma$), if each σ -zero set is also ν -zero. We say that ν and σ are *mutually singular* (and write $\nu \perp \sigma$), if there exists a Borel set Δ which is ν -full and σ -zero. Moreover, we say that ν and σ are *mutually absolutely continuous* (and write $\nu \sim \sigma$), if $\nu \ll \sigma$ and $\sigma \ll \nu$.

Each measure ν has a (essentially unique) decomposition into a sum $\nu = \nu_{ac} + \nu_s$ of a measure ν_{ac} with $\nu_{ac} \ll \sigma$ and a measure ν_s with $\nu_s \perp \sigma$; this is called the *Lebesgue decomposition* of ν with respect to σ . If $\nu \ll \sigma$, then there exists a (essentially unique) Borel measurable function $D_\sigma \nu$ with

$$\nu(\Delta) = \int_{\Delta} D_\sigma \nu d\sigma, \quad \Delta \in \mathcal{B}.$$

This function is called the *Radon-Nikodym derivative* of ν with respect to σ . It belongs to $L^1(\sigma)$, if ν is a complex measure, and to $L^1_{loc}(\sigma)$, if ν is a positive measure.

Let ν be a complex measure, and let $f \in L^1(\nu)$. Then we denote by $f \cdot \nu$ the measure which is absolutely continuous with respect to ν and has Radon-Nikodym derivative f , i.e.,

$$(f \cdot \nu)(\Delta) := \int_{\Delta} f d\nu, \quad \Delta \in \mathcal{B}. \quad (2.1)$$

In particular, if X is a Borel set, we have $(\mathbf{1}_X \cdot \nu)(\Delta) = \nu(X \cap \Delta)$, $\Delta \in \mathcal{B}$, where $\mathbf{1}_X$ denotes the indicator function of the set X . If ν is a positive measure, the same notation will be applied when $f \in L^1_{loc}(\nu)$, $f \geq 0$, and the product $f \cdot \nu$ will again be a positive measure.

The *support* of a measure ν is the set

$$\text{supp } \nu := \{x \in \mathbb{R} : \nu([x - \varepsilon, x + \varepsilon]) > 0, \varepsilon > 0\} = \bigcap_{\substack{A \text{ closed,} \\ \nu\text{-full}}} A.$$

This notion must be distinguished from the notion of a *minimal support* of the measure ν . By this one means a any Borel set S with $\nu(\mathbb{R} \setminus S) = 0$ and such that any set $S_0 \subseteq S$ with $\nu(S_0) = 0$ is also Lebesgue zero.

All these notions also make sense when ν is a projection valued measure (like the spectral measure of a selfadjoint operator) or a matrix valued measure (like the measure in the Herglotz integral representation of a matrix valued Herglotz function).

2.2. The Spectral Multiplicity Function

In order to define the spectral multiplicity function, which measures the local multiplicity of the spectrum, we have to provide some background material. These topics are of course classical, see, e.g., [1, 6, 31]. Let A be a (possibly unbounded) selfadjoint operator acting in some Hilbert space H , and denote by E its projection-valued spectral measure. A linear subspace G of H is called *generating for A* , if (“ \vee ” denotes the closed linear span)

$$\bigvee \{E(\Delta)G : \Delta \in \mathcal{B}\} = H$$

The *spectral multiplicity* of the operator A is defined as the minimal dimension of a generating subspace, and denoted by $\text{mult } A$. If $\text{mult } A = 1$, one also says that A has *simple spectrum*. For the sake of simplicity (and because this is all we need), we assume throughout the following that $\text{mult } A < \infty$.

There exist (see, e.g., [6, Theorem 7.3.7]) elements $g_l, l = 1, \dots, \text{mult } A$, such that the subspaces $H_l := \bigvee \{E(\Delta)g_l : \Delta \in \mathcal{B}\}$ are mutually orthogonal, the subspace $G := \text{span}\{g_1, \dots, g_{\text{mult } A}\}$ is generating, and the scalar measures ν_l defined as $\nu_l(\Delta) := (E(\Delta)g_l, g_l), \Delta \in \mathcal{B}$, satisfy

$$\nu_{\text{mult } A} \ll \dots \ll \nu_2 \ll \nu_1 \sim E. \tag{2.2}$$

Each set $\{g_1, \dots, g_{\text{mult } A}\}$ with these properties is called a *generating basis*.

If $\{g_1, \dots, g_{\text{mult } A}\}$ is a generating basis, the subspaces H_l are mutually orthogonal and each of them reduces the operator A . The operator $A|_{H_l}$ has simple spectrum and is unitarily equivalent to the operator of multiplication by the independent variable in the space $L_2(\mathbb{R}, \nu_l)$. Thus A is unitarily equivalent to the multiplication operator in the space $\prod_{l=1}^{\text{mult } A} L_2(\mathbb{R}, \nu_l)$. Consider the sets (which are defined up to E -zero sets)

$$Y_l := \{x \in \mathbb{R} : (D_{\nu_l} \nu_l)(x) > 0\}.$$

Then

$$\mathbb{R} \supseteq Y_1 \supseteq Y_2 \supseteq \dots \supseteq Y_{\text{mult } A},$$

and these sets may be considered as layers of the spectrum. Hence, it is natural to define the *spectral multiplicity function* of A as

$$N_A(x) := \#\{l : x \in Y_l\}. \tag{2.3}$$

This function is defined almost everywhere with respect to the projection valued spectral measure of A and does not depend on the choice of a generating basis $\{g_1, \dots, g_{\text{mult } A}\}$, cf. [6, Theorem 7.4.2]. The spectral multiplicity function is a unitary invariant. Notice that, since all definitions are almost everywhere with respect to E , one can always choose $Y_1 = \mathbb{R}$. Doing so, we then have $N_A(x) \geq 1$ a.e. with respect to E . However, sometimes it is more natural to use other choices of Y_1 .

If A is a selfadjoint linear relation, it can be orthogonally decomposed into a sum of a selfadjoint linear operator A_{op} and a pure multivalued selfadjoint linear relation (the pure relational part of A). Namely, let $H_0 := \{g \in H : (0; g) \in A\}$, then A decomposes as $A = A_{op} \oplus (\{0\} \times H_0)$ where A_0 is a selfadjoint operator in $H \ominus H_0$. In this situation, we define $N_A := N_{A_{op}}$. Obviously, this definition is also unitarily invariant.

Of course, the spectral multiplicity function of an eigenvalue is equal to the dimension of the corresponding eigenspace.

2.3. Symmetric Derivatives of Measures

In this subsection we recall the notion of the symmetric derivative of measures, and the formula of de la Vallée-Poussin which describes the Lebesgue decomposition of one positive Borel measure with respect to another. These topics are again classical, see, e.g., [32]. A presentation in an up-to-date language can be found, e.g., in [9].

Theorem 2.1. [9] *Let ν and σ be positive measures. Then there exists a Borel set $\mathcal{E}_{\nu, \sigma} \subseteq \text{supp } \nu \cap \text{supp } \sigma$ with*

$$\nu(\mathcal{E}_{\nu, \sigma}) = \sigma(\mathcal{E}_{\nu, \sigma}) = 0,$$

such that for each $x \in \text{supp } \sigma \setminus \mathcal{E}_{\nu, \sigma}$ the limit

$$\lim_{\varepsilon \downarrow 0} \frac{\nu([x - \varepsilon, x + \varepsilon])}{\sigma([x - \varepsilon, x + \varepsilon])}$$

exists in $[0, \infty]$ and defines a Borel measurable function.

Due to this proposition, we can naturally define a function which is partially defined, ν -measurable, and σ -measurable (but may be not Borel measurable).

Definition 2.2. Let ν and σ be positive measures. Then the *symmetric derivative* $\frac{d\nu}{d\sigma}$ of ν with respect to σ is the partially defined function

$$\frac{d\nu}{d\sigma}(x) := \begin{cases} \lim_{\varepsilon \downarrow 0} \frac{\nu([x - \varepsilon, x + \varepsilon])}{\sigma([x - \varepsilon, x + \varepsilon])}, & x \in \text{supp } \sigma \text{ and the limit exists in } [0, \infty], \\ \infty, & x \in \text{supp } \nu \setminus \text{supp } \sigma. \end{cases} \quad (2.4)$$

Note that this definition is symmetric in ν and σ in the following sense: *If a point x belongs to the domain of $\frac{d\nu}{d\sigma}$, then it also belongs to the domain of $\frac{d\sigma}{d\nu}$ and $\frac{d\sigma}{d\nu}(x) = \left(\frac{d\nu}{d\sigma}(x)\right)^{-1}$.*

The symmetric derivative $\frac{d\nu}{d\sigma}$ can be used to explicitly construct the Lebesgue decomposition of ν with respect to σ . To formulate this fact, denote

$$X_\infty(\nu, \sigma) := \left\{ x \in (\text{supp } \nu \cup \text{supp } \sigma) \setminus \mathcal{E}_{\nu, \sigma} : \frac{d\nu}{d\sigma}(x) = \infty \right\}.$$

Theorem 2.3. (de la Vallée-Poussin) *Let ν and σ be positive measures. Then*

- (i) *The function $\frac{d\nu}{d\sigma}$ belongs to $L^1_{\text{loc}}(\sigma)$. In particular, $\sigma(X_\infty(\nu, \sigma)) = 0$.*
- (ii) *For each Borel set $X \subseteq \mathbb{R}$, we have*

$$\nu(X) = \nu(X \cap X_\infty(\nu, \sigma)) + \int_X \frac{d\nu}{d\sigma}(x) d\sigma(x).$$

Let $\nu = \nu_{ac} + \nu_s$ be the Lebesgue decomposition of ν with respect to σ . Then, indeed, the formula of de la Vallée-Poussin says that

$$\nu_{ac} = \frac{d\nu}{d\sigma}(x) \cdot \sigma, \quad \nu_s = \mathbb{1}_{X_\infty(\nu, \sigma)} \cdot \nu. \tag{2.5}$$

In particular, if $\nu \ll \sigma$, then $\frac{d\nu}{d\sigma}$ is a Radon-Nikodym derivative of ν with respect to σ .

In the sequel we extensively use the following immediate consequences of the de la Vallée-Poussin theorem.

Corollary 2.4. *Let ν and σ be positive measures, and let $X \subseteq \mathbb{R}$.*

- (i) *If $\frac{d\nu}{d\sigma}(x) = 0$ for all $x \in X$, then X is ν -zero.*
- (ii) *If the set X is ν -zero, then $\frac{d\nu}{d\sigma}(x) = 0$ for σ -a.a. $x \in X$.*
- (iii) *If X is a Borel set and $\frac{d\nu}{d\sigma}(x) \in [0, \infty)$ for all $x \in X$, then $\mathbb{1}_X \cdot \nu \ll \sigma$.*

Proof. Item (i): The set

$$X_0(\nu, \sigma) := \left\{ x \in (\text{supp } \nu \cup \text{supp } \sigma) \setminus \mathcal{E}_{\nu, \sigma} : \frac{d\nu}{d\sigma}(x) = 0 \right\}$$

is a Borel set, and by Theorem 2.3 we have $\nu(X_0(\nu, \sigma)) = 0$. Since $X \subseteq \mathcal{E}_{\nu, \sigma} \cup X_0(\nu, \sigma)$, the set X is ν -zero.

Item (ii): There exists a Borel set $X' \supseteq X$ such that $\nu(X') = 0$. Theorem 2.3 gives $\int_{X'} \frac{d\nu}{d\sigma}(x) d\sigma(x) = 0$. This shows that $\frac{d\nu}{d\sigma}(x) = 0$ for σ -a.a. $x \in X'$, and hence for σ -a.a. $x \in X$.

Item (iii): Let $X' \in \mathcal{B}$ with $\sigma(X') = 0$ be given. Then also $\sigma(X' \cap X) = 0$, and hence

$$\nu(X' \cap X) = \int_{X' \cap X} \frac{d\nu}{d\sigma}(x) d\sigma(x) = 0. \quad \square$$

Corollary 2.5. *Let ν and σ be positive measures on \mathbb{R} . Let $\nu = \nu_{ac} + \nu_s$ and $\sigma = \sigma_{ac} + \sigma_s$ be the Lebesgue decompositions of ν with respect to σ and of σ with respect to ν , respectively. Then, the following hold:*

- (i) $\frac{d\nu}{d\sigma}(x) \in [0, \infty)$, σ -a.e.
- (ii) $\frac{d\nu}{d\sigma}(x) \in (0, \infty]$, ν -a.e.
- (iii) $\frac{d\nu}{d\sigma}(x) \in (0, \infty)$, ν_{ac} -a.e. and σ_{ac} -a.e.
- (iv) $\frac{d\nu}{d\sigma}(x) = \infty$, ν_s -a.e.
- (v) $\frac{d\nu}{d\sigma}(x) = 0$, σ_s -a.e.

Proof. Item (i) is immediate from Theorem 2.3, (i), and item (ii) follows by exchanging the roles of ν and σ and remembering that the symmetric derivative is symmetric in ν and σ . For (iii), note that

$$X := \left\{ x \in \mathbb{R} : \frac{d\nu}{d\sigma}(x) \in (0, \infty) \right\}^c \subseteq \mathcal{E}_{\nu, \sigma} \cup X_0(\nu, \sigma) \cup X_\infty(\nu, \sigma).$$

We have $\nu(\mathcal{E}_{\nu, \sigma}) = \nu(X_0(\nu, \sigma)) = 0$, and $\sigma(X_\infty(\nu, \sigma)) = 0$. Thus the union of these sets is ν_{ac} -zero. Exchanging the roles of ν and σ yields that X is also σ_{ac} -zero.

From (2.5), immediately, $\nu_s(X_\infty(\nu, \sigma)^c) = 0$. Hence, (iv) holds. Item (v) follows from (iv) again by exchanging the roles of ν and σ . \square

Remark 2.6. One can also define a symmetric derivative of a complex measure ν with respect to a positive measure σ by using the same limit

$$\frac{d\nu}{d\sigma}(x) := \lim_{\varepsilon \downarrow 0} \frac{\nu([x - \varepsilon, x + \varepsilon])}{\sigma([x - \varepsilon, x + \varepsilon])} \quad (2.6)$$

whenever it exists in \mathbb{C} . However, satisfactory knowledge can only be obtained when $\nu \ll \sigma$. In fact, the following holds: *If $\nu \ll \sigma$, then the limit (2.6) exists σ -a.e., and is a Radon-Nikodym derivative of ν with respect to σ .* This follows since we can decompose ν as $\nu = (\nu_{r,+} - \nu_{r,-}) + i(\nu_{i,+} - \nu_{i,-})$ with four positive and finite measures which are all absolutely continuous with respect to σ .

Remark 2.7. Sometimes the following facts are useful.

- (i) Existence and value of the symmetric derivative are local properties in the sense that

$$\frac{d\nu}{d\sigma}(x) = \frac{d(\mathbb{1}_X \cdot \nu)}{d(\mathbb{1}_X \cdot \sigma)}(x)$$

whenever X is a Borel set which contains x in its interior.

- (ii) If f is a continuous and nonnegative function on \mathbb{R} , then

$$\frac{d(f \cdot \nu)}{d\nu}(x) = f(x), \quad x \in \mathbb{R}.$$

2.4. Boundary Behavior of Cauchy Integrals

Let us recall the notion of *Herglotz functions*.⁴ In the present section our main focus lies on scalar valued functions. However, in view of our needs in Sect. 3, let us provide the definition and the integral representation for matrix valued functions.

Definition 2.8. A function $M : \mathbb{C} \setminus \mathbb{R} \rightarrow \mathbb{C}^{n \times n}$ is called a ($n \times n$ -matrix valued) *Herglotz function*, if

- (i) M is analytic and satisfies $M(\bar{z}) = M(z)^*$, $z \in \mathbb{C} \setminus \mathbb{R}$.
- (ii) For each $z \in \mathbb{C}^+$, the matrix $\text{Im } M(z) := \frac{1}{2i}(M(z) - M(z)^*)$ is positive semidefinite.

The following statement is known as the *Herglotz-integral representation*. For the scalar case, it goes back as far as to [17]. For the matrix valued case see [15, Theorem 5.4], where also an extensive list of references is provided.

⁴Often also called *Nevanlinna functions*.

Theorem 2.9. *Let M be a $n \times n$ -matrix valued Herglotz function. Then there exists a finite positive $n \times n$ -matrix valued measure⁵ Ω , a selfadjoint matrix a , and a positive semidefinite matrix b , such that*

$$M(z) = a + bz + \int_{\mathbb{R}} \frac{1+xz}{x-z} d\Omega(x), \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (2.7)$$

Conversely, each function of this form is a Herglotz function.

The data a, b, Ω in the representation (2.7) is uniquely determined by M . In fact, Ω can be recovered by the Stieltjes inversion formula, b from the non-tangential asymptotics of $M(z)$ towards $+i\infty$, and a from the real part of $M(i)$.

In the literature this integral representation is often written in the form

$$M(z) = a + bz + \int_{\mathbb{R}} \left(\frac{1}{x-z} - \frac{x}{1+x^2} \right) d\tilde{\Omega}(x), \quad z \in \mathbb{C} \setminus \mathbb{R}, \quad (2.8)$$

where $\tilde{\Omega}$ is a positive measure with $\int_{\mathbb{R}} \frac{d\tilde{\Omega}(x)}{1+x^2} < \infty$. The measures in the representations (2.7) and (2.8) are related as $\tilde{\Omega} = (1+x^2) \cdot \Omega$; here we use the notation (2.1).

In the present work we also consider the Cauchy-type integral in (2.7) for complex (scalar valued) measures.

Definition 2.10. Let ν be a complex (scalar valued) measure. Then we denote by m_ν the Cauchy-type integral

$$m_\nu(z) := \int_{\mathbb{R}} \frac{1+xz}{x-z} d\nu(x), \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (2.9)$$

Clearly, the function m_ν is analytic on $\mathbb{C} \setminus \mathbb{R}$. Moreover, note that for a real measure ν it can be written as the difference of two Herglotz functions.

Remark 2.11. Here is the reason why we decided to write the Herglotz integral representation in the form (2.7) rather than (2.8): For a positive measure Ω the multiplication $(1+x^2) \cdot \Omega$ is always defined and is again a positive measure (for scalar measures this is immediate, for matrix-valued measures use that Ω is mutually absolutely continuous with its trace measure $\rho := \text{tr } \Omega$). Contrasting this, for a complex measure ν , the multiplication $(1+x^2) \cdot \nu$ cannot anymore be interpreted as a measure, but only as a distribution of order 0. Since we want to avoid using the machinery of distributions, we decided for the representation (2.7).

For a finite positive measure ν , the imaginary part of $m_\nu(z)$ is (as a Poisson integral) well-behaved and several explicit relations between ν and the boundary behavior of $\text{Im } m_\nu(z)$ at the real line are known. In the present context, the following two pointwise relations play a role. The first one is standard, see, e.g., [26, §2.3]. Matching the literature is done using Remark 2.7, (ii).

⁵By a *positive matrix valued measure* we understand a measure which takes positive semidefinite matrices as values.

Theorem 2.12. [26] *Let ν be a finite positive measure, and denote by λ the Lebesgue measure.*

- (i) *Assume that $\frac{d\nu}{d\lambda}(x)$ exists in $[0, \infty]$. Then $\text{Im } m_\nu(z)$ has a normal boundary value at x , in fact,*

$$\lim_{\varepsilon \downarrow 0} \text{Im } m_\nu(x + i\varepsilon) = \pi(1 + x^2) \frac{d\nu}{d\lambda}(x).$$

- (ii) *Conversely, assume that $\text{Im } m_\nu(z)$ has a finite normal boundary value at x . Then $\frac{d\nu}{d\lambda}(x)$ exists.*

Let us note explicitly that no conclusion is drawn if $\text{Im } m_\nu(z)$ has an infinite normal boundary value at x .

The second result is in the same flavor, but may be less widely known. It is proved in [21, Lemma 1].

Theorem 2.13. [21] *Let ν be a real measure and σ be a finite positive measure. Assume that $\frac{d\nu}{d\sigma}(x)$ exists in \mathbb{R} , and that $\frac{d\sigma}{d\lambda}(x)$ exists (possibly equal to ∞) and is nonzero. Then*

$$\lim_{\varepsilon \downarrow 0} \frac{\text{Im } m_\nu(x + i\varepsilon)}{\text{Im } m_\sigma(x + i\varepsilon)} = \frac{d\nu}{d\sigma}(x).$$

The real part of a Cauchy integral is (as a singular integral) much harder to control than its imaginary part. We make use of the following two rather recent results, which deal with boundary values of $\text{Re } m_\nu(z)$, or even of $m_\nu(z)$ itself.

The first one says that the set of points $x \in \mathbb{R}$ for which $|\text{Re } m_\nu(z)|$ dominates $\text{Im } m_\nu(z)$ when z approaches x , is small. It has been shown in [28, Theorem 2.6] for Cauchy integrals of measures on the unit circle. The half-plane version stated below follows using the standard fractional-linear transform.

Theorem 2.14. [28] *Let ν be a finite positive measure. Then the set of all points $x \in \mathbb{R}$ for which there exists a continuous non-tangential path γ_x from \mathbb{C}_+ to x , such that*

$$\lim_{\substack{z \rightarrow x \\ z \in \gamma_x}} \frac{|\text{Re } m_\nu(z)|}{\text{Im } m_\nu(z)} = \infty,$$

is a ν -zero set.

The second result on singular integrals says that for certain points $x \in \mathbb{R}$ the Radon-Nikodym derivative of two measures can be calculated from the boundary behavior of the respective Cauchy integrals. This fact is shown in [27] for the disk, the half-plane version stated below is [29, Theorem 1.1]

Theorem 2.15. [27] *Let ν and σ be finite positive measures, assume that $\nu \ll \sigma$, and let σ_s be the singular part of σ with respect to the Lebesgue measure. Then:*

- (i) *For σ -a.a. points $x \in \mathbb{R}$ the non-tangential limit $\lim_{z \xrightarrow{\star} x} \frac{m_\nu(z)}{m_\sigma(z)}$ exists in $[0, \infty)$.*

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(ii) For σ_s -a.a. points $x \in \mathbb{R}$ we have

$$\lim_{z \xrightarrow{\downarrow} x} \frac{m_\nu(z)}{m_\sigma(z)} = \frac{d\nu}{d\sigma}(x).$$

Together these two theorems imply a statement which is essential for our present purposes.

Corollary 2.16. *Let ν and σ be finite positive measures, and assume that $\nu \ll \sigma$. Then for σ -a.a. points $x \in \mathbb{R}$ there exists a sequence $\{\varepsilon_n\}_{n \in \mathbb{N}}$ (which may depend on x), such that $\varepsilon_n \downarrow 0$ and the limit*

$$\lim_{n \rightarrow \infty} \frac{\operatorname{Re} m_\nu(x + i\varepsilon_n)}{\operatorname{Im} m_\sigma(x + i\varepsilon_n)}$$

exists and is finite.

Proof. Consider the sets

$$\mathcal{E}_1 := \left\{ x \in \mathbb{R} : \lim_{\varepsilon \downarrow 0} \frac{|\operatorname{Re} m_\sigma(x + i\varepsilon)|}{\operatorname{Im} m_\sigma(x + i\varepsilon)} = \infty \right\},$$

$$\mathcal{E}_2 := \left\{ x \in \mathbb{R} : \lim_{z \xrightarrow{\downarrow} x} \frac{m_\nu(z)}{m_\sigma(z)} \text{ does not exist in } \mathbb{C} \right\}.$$

Then $\sigma(\mathcal{E}_1 \cup \mathcal{E}_2) = 0$. Let $x \in (\mathcal{E}_1 \cup \mathcal{E}_2)^c$, and choose a sequence $\varepsilon_j \downarrow 0$ such that the limit $\lim_{j \rightarrow \infty} \frac{\operatorname{Re} m_\sigma(x + i\varepsilon_j)}{\operatorname{Im} m_\sigma(x + i\varepsilon_j)}$ exists in \mathbb{R} . Since

$$\frac{m_\nu(z)}{m_\sigma(z)} = \frac{m_\nu(z)}{\operatorname{Im} m_\sigma(z) \left(i + \frac{\operatorname{Re} m_\sigma(z)}{\operatorname{Im} m_\sigma(z)} \right)}, \quad z \in \mathbb{C} \setminus \mathbb{R},$$

it follows that

$$\lim_{j \rightarrow \infty} \frac{\operatorname{Re} m_\nu(x + i\varepsilon_j)}{\operatorname{Im} m_\sigma(x + i\varepsilon_j)} = \operatorname{Re} \left[\left(\lim_{j \rightarrow \infty} \frac{m_\nu(x + i\varepsilon_j)}{m_\sigma(x + i\varepsilon_j)} \right) \cdot \left(i + \lim_{j \rightarrow \infty} \frac{\operatorname{Re} m_\sigma(x + i\varepsilon_j)}{\operatorname{Im} m_\sigma(x + i\varepsilon_j)} \right) \right].$$

□

We need the above stated facts in a slightly more general situation. Namely, for arbitrary Herglotz functions rather than Cauchy-type integrals. This is easy to deduce.

Corollary 2.17. *The above statements 2.12, 2.13, 2.14, 2.15, (ii), and 2.16 remain true when the Cauchy-type integrals m_ν and m_σ are substituted by arbitrary scalar valued Herglotz functions having the measures ν and σ in their Herglotz integral representation.*

Proof. Throughout this proof, let m, \tilde{m} be Herglotz functions, and write $m(z) = a + bz + m_\nu(z)$ and $\tilde{m}(z) = \tilde{a} + \tilde{b}z + m_\sigma(z)$.

Theorem 2.12: We have

$$\lim_{\varepsilon \downarrow 0} \operatorname{Im} (a + b(x + i\varepsilon)) = 0, \quad x \in \mathbb{R}.$$

Hence the limit $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m(x + i\varepsilon)$ exists if and only if $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m_\nu(x + i\varepsilon)$ does. Moreover, if these limits exist, they coincide.

Theorem 2.13: Since $\frac{d\sigma}{d\lambda}(x)$ exists and is nonzero, we have

$$\lim_{\varepsilon \downarrow 0} \operatorname{Im} m_\sigma(x + i\varepsilon) > 0.$$

Thus

$$\begin{aligned} \lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} m(x + i\varepsilon)}{\operatorname{Im} \tilde{m}(x + i\varepsilon)} &= \lim_{\varepsilon \downarrow 0} \left[\left(\frac{b\varepsilon}{\operatorname{Im} m_\sigma(x + i\varepsilon)} + \frac{\operatorname{Im} m_\nu(x + i\varepsilon)}{\operatorname{Im} m_\sigma(x + i\varepsilon)} \right) \cdot \right. \\ &\quad \left. \cdot \left(1 + \frac{\tilde{b}\varepsilon}{\operatorname{Im} m_\sigma(x + i\varepsilon)} \right)^{-1} \right] = \lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} m_\nu(x + i\varepsilon)}{\operatorname{Im} m_\sigma(x + i\varepsilon)}. \end{aligned}$$

Theorem 2.14: The set of all points $x \in \mathbb{R}$ with $\frac{d\nu}{d\lambda}(x) \in (0, \infty]$ is ν -full. Hence, we may restrict all considerations to points x belonging to this set, and hence assume that $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m_\nu(x + i\varepsilon) > 0$. Now use the estimate

$$\begin{aligned} \frac{|\operatorname{Re} m_\nu(z)|}{\operatorname{Im} m_\nu(z)} - \frac{|a| + |b| \cdot |z|}{\operatorname{Im} m_\nu(z)} &\leq \frac{|\operatorname{Re} m(z)|}{\operatorname{Im} m(z)} \left(1 + \frac{b \operatorname{Im} z}{\operatorname{Im} m_\nu(z)} \right) \\ &\leq \frac{|\operatorname{Re} m_\nu(z)|}{\operatorname{Im} m_\nu(z)} + \frac{|a| + |b| \cdot |z|}{\operatorname{Im} m_\nu(z)}. \end{aligned}$$

Theorem 2.15, (ii): For σ_s -a.a. points $x \in \mathbb{R}$ we have $\frac{d\sigma}{d\lambda}(x) = \infty$, and hence $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m_\sigma(x + i\varepsilon) = \infty$. Thus also $\lim_{\varepsilon \downarrow 0} |m_\sigma(x + i\varepsilon)| = \infty$, and it follows that

$$\lim_{z \xrightarrow{\text{q}} x} \frac{m(z)}{\tilde{m}(z)} = \lim_{z \xrightarrow{\text{q}} x} \left[\left(\frac{a + bz}{m_\sigma(z)} + \frac{m_\nu(z)}{m_\sigma(z)} \right) \cdot \left(1 + \frac{\tilde{a} + \tilde{b}z}{m_\sigma(z)} \right)^{-1} \right] = \lim_{z \xrightarrow{\text{q}} x} \frac{m_\nu(z)}{m_\sigma(z)}.$$

Corollary 2.16: For σ -a.a. points $x \in \mathbb{R}$ we have $\frac{d\sigma}{d\lambda}(x) \in (0, \infty]$, and hence $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m_\sigma(x + i\varepsilon) > 0$. Thus also $\lim_{\varepsilon \downarrow 0} \operatorname{Im} \tilde{m}(x + i\varepsilon) > 0$, and it follows that

$$\lim_{n \rightarrow \infty} \frac{\operatorname{Re} m(x + i\varepsilon_n)}{\operatorname{Im} \tilde{m}(x + i\varepsilon_n)} = \frac{a + bx}{\lim_{n \rightarrow \infty} \operatorname{Im} m_\sigma(x + i\varepsilon_n)} + \lim_{n \rightarrow \infty} \frac{\operatorname{Re} m_\nu(x + i\varepsilon_n)}{\operatorname{Im} m_\sigma(x + i\varepsilon_n)}.$$

□

Convention: When referring to one of the above statements 2.12, 2.13, 2.14, 2.15, (ii), 2.16, we mean their general versions provided in Corollary 2.17.

3. Boundary Relations

Throughout the following we use without further notice the language and theory of linear relations. In particular, we will think of a linear operator T interchangeably as a map or as a linear relation (i.e., identify the operator T with its graph). Notationally, we interchangeably write $y = Tx$ or $(x; y) \in T$.

Our standard references for the theory of boundary triples are [7] and the survey article [8]. There also some basic notations and results about linear relations can be found.

3.1. Boundary Relations and Weyl Families

Boundary relations provide a general framework to study symmetric operators and their extensions. Let us recall their definition, see, e.g., [8, Definition 3.1].

Definition 3.1. Let S be a closed symmetric linear relation in a Hilbert space H , and let B be an auxiliary Hilbert space. A linear relation $\Gamma \subseteq H^2 \times B^2$ is called a *boundary relation for S^** , if

(BR1) The domain of Γ is contained in S^* and is dense there.

(BR2) For each two elements $((f; g); (\alpha; \beta)), ((f'; g'); (\alpha'; \beta')) \in \Gamma$ the abstract Green's identity

$$(g, f')_H - (f, g')_H = (\beta, \alpha')_B - (\alpha, \beta')_B$$

holds.

(BR3) The relation Γ is maximal with respect to the properties (BR1) and (BR2).

If the auxiliary space B is finite-dimensional, the theory of boundary relations becomes significantly simpler. Since this is all we need in the present paper, we will in most cases assume that $\dim B < \infty$.

A central notion is the Weyl family associated with a boundary relation, cf. [8, Definition 3.4].

Definition 3.2. Let $\Gamma \subseteq H^2 \times B^2$ be a boundary relation for S^* . Then, for each $z \in \mathbb{C} \setminus \mathbb{R}$, we define a linear relation $M(z)$ as

$$M(z) := \{(\alpha; \beta) \in B^2 : \exists f \in H \text{ with } ((f; zf); (\alpha; \beta)) \in \Gamma\}.$$

This family of relations is called the *Weyl family of Γ* . If $\text{mul } M(z) = \{0\}$ for all z , one also refers to M as the *Weyl function of Γ* .

Definition 3.3. A family of boundary relations $M(z)$ in the Hilbert space B is called a *Nevanlinna family*, if

- (i) for each $z \in \mathbb{C}^+$, the relation $M(z)$ is maximal dissipative;
- (ii) $M(z)^* = M(\bar{z})$, $z \in \mathbb{C} \setminus \mathbb{R}$;
- (iii) for some $w \in \mathbb{C}^+$ the operator valued function $z \mapsto (M(z) + w)^{-1}$ is holomorphic in \mathbb{C}^+ .

The basic representation theorem for Weyl families reads as follows, see, e.g., [8, Theorem 3.6] or [7, Theorem 3.9].

Theorem 3.4. [8] *Let $\Gamma \subseteq H^2 \times B^2$ be a boundary relation for S^* , and let M be its Weyl family. Then M is a Nevanlinna family. Conversely, each Nevanlinna family can be represented as the Weyl family of a boundary relation for the adjoint of some symmetric linear relation S . Moreover, S can be chosen to be completely non-selfadjoint (simple).*

This representation theorem is accompanied by the following uniqueness result, see corresponding part in the proof of [7, Theorem 3.9].

Theorem 3.5. [7] *Let S_j be closed symmetric simple linear relations in Hilbert spaces H_j , and $\Gamma_j \subseteq H_j^2 \times B^2$ be boundary relations for S_j^* , let M_j be their*

Weyl families, $j = 1, 2$. If $M_1 = M_2$, then there exists a unitary operator U of H_1 onto H_2 such that

$$\Gamma_2 = \left\{ ((Uf; Ug); (\alpha; \beta)) : ((f; g); (\alpha; \beta)) \in \Gamma_1 \right\}. \quad (3.1)$$

Two boundary relations which are related as in (3.1) are called *unitarily equivalent*.

The following properties, which a boundary relation may or may not possess, play a role in the present paper.

Definition 3.6. Let $\Gamma \subseteq H^2 \times B^2$ be a boundary relation for S^* .

(i) Γ is called of *function type*, if

$$\text{mul } \Gamma \cap (\{0\} \times B) = \{0\}.$$

(ii) Γ is called a *boundary function*, if

$$\text{mul } \Gamma = \{0\}.$$

It is an important fact that these properties reflect in properties of the Weyl family associated with Γ . For the following statement see, e.g., [8, Proposition 3.7] and [7, Lemma 4.1].

Theorem 3.7. [7] Let $\dim B < \infty$ and $\Gamma \subseteq H^2 \times B^2$ be a boundary relation for S^* . Then

- (i) Γ is of function type, if and only if $\text{mul } M(z) = \{0\}, z \in \mathbb{C} \setminus \mathbb{R}$.
- (ii) If Γ is a boundary function, then $M(z)$ is an invertible operator in B for every $z \in \mathbb{C} \setminus \mathbb{R}$.

Given a boundary relation Γ for S^* which is of function type, we can single out a particular selfadjoint extension of S . Namely, let $\pi_1 : B^2 \rightarrow B$ be the projection onto the first component, and set

$$A := \ker [\pi_1 \circ \Gamma]. \quad (3.2)$$

The fact that the relation A is selfadjoint, follows from [8, Proposition 3.16], since the auxiliary space B is finite-dimensional.

For later use, let us mention the following facts which follow from [8, Proposition 3.2].

Remark 3.8. Let $\dim B < \infty$ and $\Gamma \subseteq H^2 \times B^2$ be a boundary relation for S^* .

(i) The relation S has finite and equal defect indices $n_{\pm}(S)$. Moreover,

$$\dim B = n_{\pm}(S) + \dim \text{mul } \Gamma.$$

(ii) Assume that $\dim B = 1$. Then either $\text{mul } \Gamma = \{0\}$ and S has defect index $(1, 1)$, or $\dim \text{mul } \Gamma = 1$ and S is selfadjoint.

(iii) Assume that $\dim B = 1, \text{mul } \Gamma \neq \{0\}$ and that S is simple. Then either $\Gamma = \{0\}^2 \times (\{0\} \times \mathbb{C})$ or $\Gamma = \{0\}^2 \times \{(w; mw), w \in \mathbb{C}, m \in \mathbb{R}\}$, and the Weyl function is equal to m , a real constant.⁶

⁶The case that $\text{mul } \Gamma = \{0\} \times \mathbb{C}$ informally corresponds to the ‘‘Weyl function’’ $m \equiv \infty$ (formally to the Weyl family $m \equiv \{0\} \times \mathbb{C}$).

Next, we recall four methods to construct new boundary relations from given ones. The first one is just taking orthogonal sums, the second is making a change of basis. Both are easy to verify (and common knowledge); we skip the details.

Lemma 3.9. *Let $n \in \mathbb{N}$, and let for each $l \in \{1, \dots, n\}$ a boundary relation $\Gamma_l \subseteq H_l^2 \times B_l^2$ for S_l^* be given. Define*

$$\prod_{l=1}^n \Gamma_l := \left\{ \left(\left(\left(\begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix}; \begin{pmatrix} g_1 \\ \vdots \\ g_n \end{pmatrix} \right); \left(\begin{pmatrix} \alpha_1 \\ \vdots \\ \alpha_n \end{pmatrix}; \begin{pmatrix} \beta_1 \\ \vdots \\ \beta_n \end{pmatrix} \right) \right) : \right. \\ \left. ((f_l; g_l); (\alpha_l; \beta_l)) \in \Gamma_l, l = 1, \dots, n \right\} \subseteq \left(\prod_{l=1}^n H_l \right)^2 \times \left(\prod_{l=1}^n B_l \right)^2.$$

Then $\prod_{l=1}^n \Gamma_l$ is a boundary relation for $\prod_{l=1}^n S_l^*$. Its Weyl family is given as

$$\prod_{l=1}^n M_l(z) := \left\{ \left(\left(\begin{pmatrix} \alpha_1 \\ \vdots \\ \alpha_n \end{pmatrix}; \begin{pmatrix} \beta_1 \\ \vdots \\ \beta_n \end{pmatrix} \right) : (\alpha_l; \beta_l) \in M_l(z), l = 1, \dots, n \right\} \subseteq \left(\prod_{l=1}^n B_l \right)^2.$$

The relation $\prod_{l=1}^n \Gamma_l$ is of function type (a boundary function) if and only if all relations Γ_l are.

Lemma 3.10. *Let $\Gamma \subseteq H^2 \times B^2$ be a boundary relation for S^* of function type with Weyl function M . Moreover, let $U : B \rightarrow B$ be unitary, and define*

$$\Gamma_1 := \left\{ ((f; g); (U\alpha; U\beta)) : ((f; g); (\alpha; \beta)) \in \Gamma \right\} \subseteq H^2 \times B^2.$$

Then Γ_1 is a boundary relation for S^* which is of function type, and the Weyl function M_1 is given as

$$M_1(z) = UM(z)U^{-1}.$$

The relation Γ_1 is a boundary function, if and only if Γ is.

The third construction shows how to realize the sum of two Weyl functions as a Weyl function, cf. [8, Corollary 4.5].

Theorem 3.11. [8] *Assume that $\dim B < \infty$. For $j = 1, 2$, let $\Gamma_j \subseteq H_j^2 \times B^2$ be boundary relations for S_j^* of function type with corresponding Weyl functions M_j . Define*

$$\Gamma := \left\{ \left(\left(\begin{pmatrix} f_1 \\ f_2 \end{pmatrix}; \begin{pmatrix} g_1 \\ g_2 \end{pmatrix} \right); (\alpha; \beta_1 + \beta_2) \right) : \right. \\ \left. ((f_j; g_j); (\alpha; \beta_j)) \in \Gamma_j, j = 1, 2 \right\} \subseteq (H_1 \times H_2)^2 \times B^2,$$

and

$$S := \left\{ \left(\begin{pmatrix} f_1 \\ f_2 \end{pmatrix}; \begin{pmatrix} g_1 \\ g_2 \end{pmatrix} \right) : \exists \beta \in B \text{ with} \right. \\ \left. ((f_1; g_1); (0; \beta)) \in \Gamma_1, ((f_2; g_2); (0; -\beta)) \in \Gamma_2 \right\} \subseteq (H_1 \times H_2)^2.$$

Then S is a closed symmetric relation in $H_1 \times H_2$, and Γ is a boundary relation for S^* . The relation Γ is of function type, and its Weyl function is given as

$$M(z) = M_1(z) + M_2(z).$$

With the fourth procedure we construct a new boundary relation via a fractional linear transform. A proof can be found in [8, Proposition 3.11]. Before we can formulate this, let us introduce one more notation. We denote by $J_{\mathbb{C}^n}$ the $2n \times 2n$ -matrix

$$J_{\mathbb{C}^n} := i \begin{pmatrix} 0 & I_{\mathbb{C}^n} \\ -I_{\mathbb{C}^n} & 0 \end{pmatrix}.$$

Then $J_{\mathbb{C}^n}$ defines a non-degenerated inner product on \mathbb{C}^{2n} . Let w be a $2n \times 2n$ -matrix. Then w is $J_{\mathbb{C}^n}$ -unitary (i.e., unitary with respect to the inner product induced by $J_{\mathbb{C}^n}$) if and only if

$$w^* J_{\mathbb{C}^n} w = J_{\mathbb{C}^n}. \quad (3.3)$$

Theorem 3.12. [8] *Let $\Gamma \subseteq H^2 \times (\mathbb{C}^n)^2$ be a boundary relation for S^* . Let w be a $J_{\mathbb{C}^n}$ -unitary $2n \times 2n$ -matrix, and write w in block form as $w = (w_{ij})_{i,j=1}^2$ with $n \times n$ -blocks $w_{ij}, i, j = 1, 2$. Then the composition*

$$\Gamma_1 := w \circ \Gamma = \left\{ ((f; g); (w_{11}\alpha + w_{12}\beta; w_{21}\alpha + w_{22}\beta)) : ((f; g); (\alpha; \beta)) \in \Gamma \right\}.$$

is a boundary relation for S^* , and its Weyl family $M_1(z)$ is given as

$$M_1(z) = \left\{ (w_{11}\alpha + w_{12}\beta; w_{21}\alpha + w_{22}\beta) : (\alpha; \beta) \in M(z) \right\}, \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (3.4)$$

3.2. The Titchmarsh–Kodaira Formula

Let Γ be a boundary relation of function type, and let A be the selfadjoint relation (3.2). Then the data a, b, Ω in the integral representation (2.7) of the Weyl function M associated with Γ can be used to construct a functional model for $A = \ker[\pi_1 \circ \Gamma]$ acting as the multiplication operator in an L^2 -space (to be exact, the relational analogue of the multiplication operator). For a Schrödinger operator on the half line (meaning limit-circle on one end and limit-point on the other) or on the whole line (limit-point at both ends) this is a classical fact. In the first case (where the spectral measure is scalar) this goes back to the initial considerations of Weyl, cf. [37], in the second case (where the spectral measure is 2×2 -matrix valued) to the independent works of Titchmarsh and Kodaira, cf. [35] and [23]. See also [12, XII.5. Theorems 13 and 14], where differential expressions of arbitrary order are studied.

In the present context this functional model plays an important role, since it allows us to compute the spectral multiplicity function by computing

the rank of a certain matrix; the precise statement being Proposition 3.15 below.

First, let us recall the appropriate notion of a “multiplication operator”. Let $\Omega = (\Omega_{ij})_{i,j=1}^n$ be a positive $n \times n$ -matrix valued measure, and denote by ρ the (scalar) trace-measure $\rho := \sum_{i=1}^n \Omega_{ii}$. Since, for each Borel set Δ , the matrix $\Omega(\Delta)$ is positive semidefinite, we have $|\Omega_{ij}(\Delta)| \leq \sqrt{\Omega_{ii}(\Delta)\Omega_{jj}(\Delta)}$, $i, j = 1, \dots, n$, and this yields that $\Omega \sim \rho$. Hence, the Radon-Nikodym derivative $D_\rho \Omega = (D_\rho \Omega_{ij})_{i,j=1}^n$ is well-defined and ρ -a.e. positive.

Consider now the set of all ρ -a.e. finite functions $f : \mathbb{R} \mapsto \mathbb{C}^n$, such that each component is ρ -measurable and such that

$$\int_{\mathbb{R}} (f(x), D_\rho \Omega(x) f(x))_{\mathbb{C}^n} d\rho(x) < \infty.$$

The space $L_2(\Omega)$ is the space of equivalence classes of such functions under the equivalence relation

$$f \sim g \text{ if and only if } \int_{\mathbb{R}} (f(x) - g(x), D_\rho \Omega(x)(f(x) - g(x)))_{\mathbb{C}^n} d\rho(x) = 0.$$

When endowed with the inner product⁷

$$(f, g)_\Omega := \int_{\mathbb{R}} (f(x), D_\rho \Omega(x) g(x))_{\mathbb{C}^n} d\rho(x), \quad f, g \in L^2(\Omega),$$

this space becomes a Hilbert space. The operator of multiplication A_t by the independent variable t in this space is selfadjoint, see [19] or [12, XIII.5.Theorem 10].

Moreover, for a positive semidefinite $n \times n$ -matrix b , denote by G_b the space $\text{ran } b$ endowed with the inner product defined as

$$(bx, by)_{G_b} := (bx, y)_{\mathbb{C}^n}, \quad x, y \in \mathbb{C}^n.$$

Definition 3.13. Let Ω be a positive $n \times n$ -matrix valued measure, and let b be a positive semidefinite $n \times n$ -matrix. Then we set

$$H_{\Omega,b} := L^2(\Omega) \oplus G_b, \quad A_{\Omega,b} := A_t \oplus (\{0\} \times G_b).$$

Clearly, $H_{\Omega,b}$ is a Hilbert space and $A_{\Omega,b}$ is a selfadjoint linear relation in $H_{\Omega,b}$.

Let us now provide the afore mentioned functional model for $A = \ker[\pi_1 \circ \Gamma]$. The essence of this result is the model constructed in [10, Proposition 5.2], and it is easily deduced from this.

Proposition 3.14. *Let $\Gamma \subseteq H^2 \times \mathbb{C}^{2n}$ be a boundary relation of function type with Weyl function M , let a, b, Ω be the data in the integral representation (2.7) of M , and set $\tilde{\Omega} := (1 + t^2) \cdot \Omega$. Then the selfadjoint relation $A := \ker[\pi_1 \circ \Gamma]$ is unitarily equivalent to the relation $A_{\tilde{\Omega},b}$.*

⁷The right-hand side of this formula does not depend on the choice of the representative of the equivalence class; as usual we slightly abuse notation.

Proof. The following facts are well-known (for an explicit proof see, e.g., [3, Proof of Theorem 3.2]): The kernel of $\text{Im } M(z)$ for $z \in \mathbb{C} \setminus \mathbb{R}$ is independent of z . Denote its codimension by n_0 . There exists a constant unitary matrix U such that $U^{-1} \ker \text{Im } M(i) = \{0\} \oplus \mathbb{C}^{n-n_0}$ and $U^{-1}(M(z) - a)U = M_0(z) \oplus 0$ (the block form with respect to the decomposition $\mathbb{C}^n = \mathbb{C}^{n_0} \oplus \mathbb{C}^{n-n_0}$). At the same time $U^{-1}\tilde{\Omega}U = \tilde{\Omega}_0 \oplus 0$, $U^{-1}bU = b_0 \oplus 0$ and

$$M_0(z) = b_0z + \int_{\mathbb{R}} \left(\frac{1}{x-z} - \frac{x}{1+x^2} \right) d\tilde{\Omega}_0(x)$$

is the Herglotz-integral representation for the function M_0 . The latter is such that $\text{Im } M_0(z)$ is invertible for each $z \in \mathbb{C} \setminus \mathbb{R}$. We apply [10, Proposition 5.2] to obtain a functional model for M_0 : The relation

$$S := \left\{ (f(t) \oplus 0; tf(t) \oplus y) \in H_{\tilde{\Omega}_0, b_0}^2 : \int_{\mathbb{R}} d\tilde{\Omega}_0(t)f(t) + y = 0 \right\}$$

is closed symmetric and simple, its adjoint is given by

$$S^* := \left\{ (f(t) \oplus x; g(t) \oplus y) \in H_{\tilde{\Omega}_0, b_0}^2 : \exists h \in \mathbb{C}^{n_0} : g(t) - tf(t) = -h, x = b_0h \right\}.$$

Obviously, the element h in this formula is uniquely determined by the element $(f(t) \oplus x; g(t) \oplus y)$ of S^* . Hence, we may define

$$\begin{aligned} \Gamma_0((f(t) \oplus x; g(t) \oplus y)) &:= h, \\ \Gamma_1((f(t) \oplus x; g(t) \oplus y)) &:= y + \int_{\mathbb{R}} d\tilde{\Omega}_0(t) \frac{tg(t) + f(t)}{1+t^2}. \end{aligned}$$

Then it follows from [10] that the relation

$$\begin{aligned} \mathring{\Gamma} := & \left\{ \left((f(t) \oplus x; g(t) \oplus y); (\Gamma_0(f(t) \oplus x; g(t) \oplus y); \Gamma_1(f(t) \oplus x; g(t) \oplus y)) \right) : \right. \\ & \left. (f(t) \oplus x; g(t) \oplus y) \in S^* \right\} \end{aligned}$$

is a boundary function for S^* . The Weyl function of $\mathring{\Gamma}$ is equal to M_0 , and $\ker \Gamma_0 = A_{\tilde{\Omega}_0, b_0}$. Obviously $\Gamma_U := \mathring{\Gamma} \oplus (\{0\}^2 \times \{(w; 0), w \in \mathbb{C}^{n-n_0}\})$ is a boundary relation of function type for S^* with the Weyl function $M_0 \oplus 0$ and $\ker[\pi_1 \circ \Gamma_U] = A_{\tilde{\Omega}_0, b_0}$. Next, $\hat{\Gamma} := \{((f; g); (U\alpha; U\beta)) : ((f; g); (\alpha; \beta)) \in \Gamma_U\}$ is another boundary relation of function type for S^* with the Weyl function $M - a$ and $\ker[\pi_1 \circ \hat{\Gamma}] = A_{\tilde{\Omega}_0, b_0}$. Finally, the selfadjoint constant a is the Weyl function of the boundary relation $\{0\}^2 \times \{(w; aw), w \in \mathbb{C}^n\}$ acting in $\{0\}^2 \times \mathbb{C}^{2n}$. Using Theorem 3.11, we obtain a boundary relation $\tilde{\Gamma}$ having M as its Weyl function. Explicitly computing $\tilde{\Gamma}$ shows that the relation $\ker[\pi_1 \circ \tilde{\Gamma}]$ coincides with $A_{\tilde{\Omega}_0, b_0}$. Uniqueness part of Theorem 3.5 ensures that $\ker[\pi_1 \circ \tilde{\Gamma}]$ is unitarily equivalent to $\ker[\pi_1 \circ \tilde{\Gamma}] = A_{\tilde{\Omega}_0, b_0}$. Obviously, the relation $A_{\tilde{\Omega}_0, b_0}$ can be identified with the relation $A_{\tilde{\Omega}_0 \oplus 0, b_0 \oplus 0} = A_{U^{-1}\tilde{\Omega}U, U^{-1}bU}$, which is unitarily equivalent to $A_{\tilde{\Omega}, b}$. \square

Now we come to the promised way to compute the spectral multiplicity function. For the case “ $n = 2$ ” this fact is proved and used in [20], see also [14]. It is of course not hard to believe that it holds for arbitrary $n \geq 2$, however, we are not aware of an explicit reference, and therefore provide a complete proof.

Proposition 3.15. *Let $\Gamma \subseteq H^2 \times B^2$ be a boundary relation of function type. Denote by M its Weyl function, and set $A := \ker[\pi_1 \circ \Gamma]$. Let a, b, Ω be the data in the integral representation (2.7) of M , let ρ be the trace measure of Ω , and let ω be the symmetric derivative*

$$\omega := \frac{d\Omega}{d\rho}. \tag{3.5}$$

Then (N_A is the spectral multiplicity function of A)

$$N_A = \text{rank } \omega, \quad \rho\text{-a.e.} \tag{3.6}$$

Proof. From the readily established by Proposition 3.14 unitary equivalence, we see that it is enough to compute the spectral multiplicity function of the multiplication operator A_t in the space $L^2(\Omega)$. To do this, the idea is to construct a measurable (ρ -measurable, or Borel measurable on a compliment of some Borel ρ -zero set) diagonalization of $\omega(x)$. Once this is done, it is easy to give a unitarily equivalent form of A (and a particular generating basis) from which $N_A(x)$ can be read off. The essential tool in the proof is Hammersley’s theorem on the measurability of the zeros of a random polynomial, cf. [16, Theorem 4.1].⁸

By Hammersley’s theorem there exist measurable functions ξ_1, \dots, ξ_n such that

$$\det [\omega(x) - t] = (-1)^n \prod_{j=1}^n (t - \xi_j(x)), \quad x \in \mathbb{R}.$$

Since $\omega(x)$ is nonnegative, we have $\xi_j(x) \geq 0, j = 1, \dots, n$. By pointwise rearranging (which can be done in a measurable way) we can redefine the functions ξ_j , such that in addition

$$0 \leq \xi_n(x) \leq \xi_{n-1}(x) \leq \dots \leq \xi_1(x), \quad x \in \mathbb{R}.$$

What follows is basic linear algebra. For $j, k \in \{1, \dots, n\}$ and $\{i_1, \dots, i_k\} \subseteq \{1, \dots, n\}, i_1 < \dots < i_k$, set

$$M_{i_1, \dots, i_k}^{j, k} := \left\{ x \in \mathbb{R} : \text{rank} [\omega(x) - \xi_j(x)] = k, \right. \\ \left. \det (w_{i_i i_m}(x) - \xi_j(x) \delta_{i_i i_m})_{l, m=1}^k \neq 0 \right\}.$$

The determinant of a matrix is a polynomial of the entries, and hence is measurable. The rank of a matrix depends, as the maximal order of an

⁸See also [4, Theorem 2.2], where a short proof based on von Neumann’s measurable selection theorem is given.

invertible square minor, measurably on the entries of the matrix. It follows that $M_{i_1, \dots, i_k}^{j,k}$ is a Borel set. Also the set

$$M_{\emptyset}^{j,0} := \{x \in \mathbb{R} : \omega(x) = \xi_j(x)\}$$

is a Borel set.

Let $j \in \{1, \dots, n\}$ be fixed. For each $x \in M_{i_1, \dots, i_k}^{j,k}$ the submatrix $(w_{i_l i_m}(x) - \xi_j(x) \delta_{i_l i_m})_{l,m=1}^k$ of $[\omega(x) - \xi_j(x)]$ is invertible. Applying Cramer's rule, we find a basis of the eigenspace $\ker[\omega(x) - \xi_j(x)]$ which depends measurably on $x \in M_{i_1, \dots, i_k}^{j,k}$. Applying the Gram-Schmidt orthogonalization procedure, we obtain an orthonormal basis which also depends measurably on $x \in M_{i_1, \dots, i_k}^{j,k}$. For $x \in M_{\emptyset}^{j,0}$ the canonical basis of \mathbb{C}^n is an orthonormal basis of $\ker[\omega(x) - \xi_j(x)]$.

Clearly, for each j , the sets $M_{i_1, \dots, i_k}^{j,k}, k \in \{0, \dots, n\}, \{i_1, \dots, i_k\} \subseteq \{1, \dots, n\}$ together cover the whole line. Hence, we can produce a disjoint covering of \mathbb{R} with each set of the covering being a Borel subset of some intersection

$$\bigcap_{j=1}^n M_{i_{j,1}, \dots, i_{j,k_j}}^{j,k_j}.$$

By the above paragraph, we can thus find a measurable orthonormal basis in \mathbb{C}^n which consists of eigenvectors of ω . The corresponding basis transform $U(x)$ is a measurable function and diagonalizes $\omega(x)$:

$$U(x)^{-1} \omega(x) U(x) = D(x)$$

with

$$D(x) := \begin{pmatrix} \xi_1(x) & & \\ & \ddots & \\ & & \xi_n(x) \end{pmatrix}.$$

The map $f \mapsto U^{-1}f$ is an isometric isomorphism of $L^2(\Omega)$ onto $L^2(D \cdot \rho)$, and establishes a unitary equivalence between the respective multiplication operators. One can regard $\xi_l \cdot \rho$ as ν_l from (2.2), so $D_{\nu_l} \nu_l(x) = \xi_l(x)$. Therefore the spectral multiplicity function computes as

$$\#\{l \in \{1, \dots, n\} : \xi_l(x) > 0\} = \text{rank } D(x) = \text{rank } \omega(x). \quad \square$$

3.3. Pasting of Boundary Relations with Standard Interface Conditions

Let $n \geq 2$, and let for each $l \in \{1, \dots, n\}$ a closed symmetric relation S_l in a Hilbert space H_l and a boundary relation $\Gamma_l \subseteq H_l^2 \times \mathbb{C}^2$ for S_l^* be given. Moreover, denote

$$A_l := \ker [\pi_1 \circ \Gamma_l]. \quad (3.7)$$

Consider the Hilbert space $H := \prod_{l=1}^n H_l$, the linear relation $S := \prod_{l=1}^n S_l$ acting in this space, and the orthogonal sum $\tilde{\Gamma} = \prod_{l=1}^n \Gamma_l$, cf. Lemma 3.9.

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Now we define another boundary relation Γ by using in Theorem 3.12 the $J_{\mathbb{C}^n}$ -unitary matrix $w = (w_{ij})_{i,j=1}^2$ whose blocks w_{ij} are given as

$$w_{11} := \begin{pmatrix} -1 & 0 & \cdots & 0 & 1 \\ 0 & -1 & \cdots & 0 & 1 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & -1 & 1 \\ 0 & 0 & \cdots & 0 & 0 \end{pmatrix}, \quad w_{12} := \begin{pmatrix} 0 & 0 & \cdots & 0 & 0 \\ 0 & 0 & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & 0 & 0 \\ 1 & 1 & \cdots & 1 & 1 \end{pmatrix}, \quad (3.8)$$

$$w_{21} := \begin{pmatrix} 0 & 0 & \cdots & 0 & 0 \\ 0 & 0 & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & 0 & 0 \\ 0 & 0 & \cdots & 0 & -1 \end{pmatrix}, \quad w_{22} := \begin{pmatrix} -1 & 0 & \cdots & 0 & 0 \\ 0 & -1 & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & -1 & 0 \\ 0 & 0 & \cdots & 0 & 0 \end{pmatrix}. \quad (3.9)$$

A straightforward computation shows that this matrix w indeed satisfies (3.3). Explicitly, the relation Γ is given as

$$\Gamma := w \circ \tilde{\Gamma} = \left\{ \left(\left(\left(\begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix} \right); \begin{pmatrix} g_1 \\ \vdots \\ g_n \end{pmatrix} \right) \right); \left(\left(\begin{pmatrix} -\alpha_1 + \alpha_n \\ \vdots \\ -\alpha_{n-1} + \alpha_n \\ \beta_1 + \cdots + \beta_n \end{pmatrix} \right); \begin{pmatrix} -\beta_1 \\ \vdots \\ -\beta_{n-1} \\ -\alpha_n \end{pmatrix} \right) \right) \right\}; \\ ((f_l; g_l); (\alpha_l; \beta_l)) \in \Gamma_l, \quad l = 1, \dots, n \}. \quad (3.10)$$

Definition 3.16. Let $n \geq 2$, and let for each $l \in \{1, \dots, n\}$ a closed symmetric relation S_l in a Hilbert space H_l and a boundary relation $\Gamma_l \subseteq H_l^2 \times \mathbb{C}^2$ for S_l^* be given. Assume that

- (Hyp1) Each relation S_l is simple.
- (Hyp2) Each boundary relation Γ_l is of function type.
- (Hyp3) There exists an index $l \in \{1, \dots, n\}$, such that Γ_l is a boundary function.

Obviously, the relation Γ depends on the order in which the relations $\Gamma_1, \Gamma_2, \dots, \Gamma_n$ are taken, but the selfadjoint relation $A := \ker[\pi_1 \circ \Gamma]$ does not. Thus we call A the relation constructed by *pasting the family* $\{\Gamma_l : l = 1, \dots, n\}$, *with standard interface conditions*.

The hypothesis (Hyp1) ensures that knowledge about the spectrum of A can be deduced from the associated Weyl family (in fact, the Weyl function, see below). The hypotheses (Hyp2), (Hyp3), are required in order to avoid trivial cases (remember Remark 3.8).

To justify our choice of terminology, let us return to our model example.

Example 3.17. Let a Schrödinger operator on a star-shaped graph be given by the data (1)–(3). Let S_l be the minimal operator on the l -th edge, i.e.

$$\begin{aligned} \text{dom } S_l &:= \left\{ u \in L^2(0, e_l) : u, u' \text{ absolutely continuous,} \right. \\ &\quad \left. -u'' + q_l u \in L_2(0, e_l), \right. \\ &\quad \left. u(0) = u'(0) = 0, u \text{ satisfies b.c. at } e_l, \text{ if present} \right\}. \\ S_l u &:= -u'' + q_l u, \quad u \in \text{dom } S_l. \end{aligned}$$

Moreover, define

$$\Gamma_l := \{((u; v); (u(0); u'(0)) : (u; v) \in S_l^*\}.$$

Then Γ_l is a boundary relation for S_l^* . The selfadjoint extension $A_l = \ker[\pi_1 \circ \Gamma_l]$ is just the Schrödinger operator given by the potential q_l with Dirichlet boundary conditions at 0.

Now consider the boundary relation Γ defined by (3.10). The operator $A := \ker[\pi_1 \circ \Gamma]$, is nothing but the operator defined by (1.1), (1.2), using the standard interface condition (1.3).

In order to understand the spectrum of a pasting with standard interface conditions, we will analyze the Weyl function of the boundary relation Γ . Using Theorem 3.12, this Weyl function can be computed explicitly in terms of the Weyl functions of the boundary relations Γ_l .

Proposition 3.18. *Let Γ_l be as in Definition 3.16, and let Γ be the boundary relation given by (3.10). Denote by m_l the Weyl function of Γ_l , and set $m := \sum_{l=1}^n m_l$. Then we have $\text{mul } \Gamma \cap (\{0\} \times \mathbb{C}^n) = \{0\}$, and the Weyl function M of Γ is given as*

$$M = \frac{1}{m} \begin{pmatrix} m_1(m-m_1) & -m_2 m_1 & \cdots & -m_{n-1} m_1 & -m_1 \\ -m_1 m_2 & m_2(m-m_2) & \cdots & -m_{n-1} m_2 & -m_2 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ -m_1 m_{n-1} & -m_2 m_{n-1} & \cdots & m_{n-1}(m-m_{n-1}) & -m_{n-1} \\ -m_1 & -m_2 & \cdots & -m_{n-1} & -1 \end{pmatrix}. \quad (3.11)$$

Proof. Consider an element of $\text{mul } \Gamma \cap (\{0\} \times \mathbb{C}^n)$. By the definition of Γ there exist $(\alpha_l; \beta_l) \in \text{mul } \Gamma_l, l = 1, \dots, n$, such that this element is equal to

$$\left(\left(\begin{pmatrix} -\alpha_1 + \alpha_n \\ \vdots \\ -\alpha_{n-1} + \alpha_n \\ \beta_1 + \cdots + \beta_n \end{pmatrix}; \begin{pmatrix} -\beta_1 \\ \vdots \\ -\beta_{n-1} \\ -\alpha_n \end{pmatrix} \right) \right) = \left(\begin{pmatrix} 0 \\ \vdots \\ 0 \\ 0 \end{pmatrix}; \begin{pmatrix} -\beta_1 \\ \vdots \\ -\beta_{n-1} \\ -\alpha_n \end{pmatrix} \right).$$

By (Hyp3), there exists an index l_0 with $\alpha_{l_0} = \beta_{l_0} = 0$. Now it follows that $\alpha_l = 0$ for all l , and (Hyp2) implies that also $\beta_l = 0$ for all l . This shows that $\text{mul } \Gamma \cap (\{0\} \times \mathbb{C}^n) = \{0\}$.

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The Weyl function of the boundary relation $\tilde{\Gamma} = \prod_{l=1}^n \Gamma_l$ is

$$\tilde{M} = \begin{pmatrix} m_1 & 0 & \cdots & 0 \\ 0 & m_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & m_n \end{pmatrix}.$$

Computation gives:

$$w_{11} + w_{12}\tilde{M} = \begin{pmatrix} -1 & \cdots & 0 & 1 \\ \vdots & \ddots & \vdots & \vdots \\ 0 & \cdots & -1 & 1 \\ m_1 & \cdots & m_{n-1} & m_n \end{pmatrix},$$

$$w_{21} + w_{22}\tilde{M} = \begin{pmatrix} m_1 & \cdots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots \\ 0 & \cdots & -m_{n-1} & 0 \\ 0 & \cdots & 0 & -1 \end{pmatrix},$$

and

$$\det(w_{11} + w_{12}\tilde{M}) = (-1)^{n-1} \sum_{l=1}^n m_l.$$

Since there exists at least one index l such that m_l is not a real constant, this determinant does not vanish throughout $\mathbb{C} \setminus \mathbb{R}$.

Next, let M be the matrix defined by (3.11). It is easy to check that M satisfies $M(w_{11} + w_{12}\tilde{M}) = w_{21} + w_{22}\tilde{M}$, and this implies that

$$M = (w_{21} + w_{22}\tilde{M})(w_{11} + w_{12}\tilde{M})^{-1}.$$

Theorem 3.12 now yields that M is indeed the Weyl function of Γ . \square

3.4. The Point Spectrum

It is elementary to locate the point spectrum of a pasting.

Theorem 3.19. *Let $n \geq 2$, and let for each $l \in \{1, \dots, n\}$ a closed symmetric relation S_l in a Hilbert space H_l and a boundary relation $\Gamma_l \subseteq H_l^2 \times \mathbb{C}^2$ for S_l^* be given. Assume that these data are subject to (Hyp1)–(Hyp3), and consider the selfadjoint operator A constructed by pasting $\{\Gamma_l : l = 1, \dots, n\}$ with standard interface conditions.*

Let $x \in \mathbb{R}$. Then $x \in \sigma_p(A)$ if and only if one of the following two alternatives takes place.

(I/II) *The point x belongs to at least two of the point spectra $\sigma_p(A_l)$. In this case the multiplicity $N_A(x)$ of x as an eigenvalue is equal to*

$$\begin{aligned} N_A(x) &= \#\{l \in \{1, \dots, n\} : x \in \sigma_p(A_l)\} - 1 \\ &= \#\{l \in \{1, \dots, n\} : \lim_{\varepsilon \downarrow 0} \varepsilon \operatorname{Im} m_l(x + i\varepsilon) > 0\} - 1. \end{aligned}$$

(III) *The limits $m_l(x) := \lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon)$ all exist, are real, we have $\lim_{\varepsilon \downarrow 0} \frac{1}{i\varepsilon} (m_l(x + i\varepsilon) - m_l(x)) \in [0, \infty)$, and $\sum_{l=1}^n m_l(x) = 0$. In this case x is a simple eigenvalue.*

First we prove a technical statement which is an immediate consequence of (Hyp1), (Hyp2).

Lemma 3.20. *Let S be a closed symmetric simple relation in a Hilbert space H , let $\Gamma \subseteq H^2 \times \mathbb{C}^2$ be a boundary relation for S^* of function type. Set $A := \ker[\pi_1 \circ \Gamma]$.*

- (i) *For each $(f; g) \in A$ there exists a unique number $\beta \in \mathbb{C}$ such that $((f; g); (0; \beta)) \in \Gamma$.*
- (ii) *Let $\Upsilon : A \rightarrow \mathbb{C}$ be defined as $\Upsilon(f; g) = \beta$ where $((f; g); (0; \beta)) \in \Gamma$, and let $x \in \mathbb{R}$. Then the restriction of Υ to the set $\{(f; xf) : f \in \ker(A - x)\}$ is injective.*
- (iii) *Let $x \in \mathbb{R}$ and $\alpha \in \mathbb{C} \setminus \{0\}$. Then there exists at most one element of the form $((f; xf); (\alpha; \beta))$ which belongs to Γ .*

Proof. Existence of β is the definition of A . Uniqueness follows since $\text{mul } \Gamma \cap (\{0\} \times \mathbb{C}) = \{0\}$. This shows (i) and that the map Υ in (ii) is well-defined. Since Γ is minimal, S is completely non-selfadjoint. We clearly have $\ker \Upsilon \subseteq S$, and hence $\ker \Upsilon \cap \{(f; xf) : f \in \ker(A - x)\} = \{0\}$. This shows (ii).

To show (iii), assume that $((f; xf); (\alpha; \beta)), ((f'; xf'); (\alpha; \beta')) \in \Gamma$ and that $f \neq f'$. Without loss of generality assume that $f \neq 0$. By minimality of Γ this implies that (remember Remark 3.8)

$$\text{mul } \Gamma = \{0\}, \quad \dim S^*/S = 2, \quad \dim \mathcal{N}_x \leq 1, \quad x \in \mathbb{R},$$

where $\mathcal{N}_x := \{f \in H : (f; zf) \in S^*\}$. Let λ be such that $f' = \lambda f$, then $((0; 0); ((\lambda - 1)\alpha; \lambda\beta - \beta')) \in \Gamma$. Since $\lambda \neq 1$ and $\alpha \neq 0$, we obtain $\text{mul } \Gamma \neq \{0\}$ a contradiction. We conclude that $f = f'$. Since $\text{mul } \Gamma \cap (\{0\} \times \mathbb{C}) = \{0\}$, it follows that also $\beta = \beta'$. \square

Proof of Theorem 3.19. Step 1; A preliminary observation: In this step, we show that

$$\left(\begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix}; \begin{pmatrix} g_1 \\ \vdots \\ g_n \end{pmatrix} \right) \in A \implies \exists! \alpha, \beta_l : ((f_l; g_l); (\alpha; \beta_l)) \in \Gamma_l$$

Moreover, for these numbers β_l , it holds that $\beta_1 + \dots + \beta_n = 0$.

By the definition of A and Γ , cf. (3.10), there exist $\alpha_1, \dots, \alpha_n$ and β_1, \dots, β_n such that

$$((f_l; g_l); (\alpha_l; \beta_l)) \in \Gamma_l, \quad -\alpha_1 + \alpha_n = \dots = -\alpha_{n-1} + \alpha_n = \beta_1 + \dots + \beta_n = 0.$$

This proves the existence part (set $\alpha := \alpha_1$). For uniqueness, assume that α' and $\beta'_1, \dots, \beta'_n$ are such that $((f_l; g_l); (\alpha'; \beta'_l)) \in \Gamma_l$. By (Hyp3) there exists an index $l_0 \in \{1, \dots, n\}$ with $\text{mul } \Gamma_{l_0} = \{0\}$. It follows that $\alpha' = \alpha$ and $\beta'_{l_0} = \beta_{l_0}$. Due to (Hyp2) it follows that for all indices $\beta'_l = \beta_l$.

Step 2; Two examples of eigenvectors: Let $x \in \mathbb{R}$. First, consider the space

$$\mathcal{L}_x := \left\{ \begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix} \in \prod_{l=1}^n \ker(A_l - x) : \sum_{l=1}^n \Upsilon_l(f_l; x f_l) = 0 \right\}.$$

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Since each map Υ_l is injective, we have

$$\dim \mathcal{L}_x = \begin{cases} 0, & \text{if for every } l \ker(A_l - x) = \{0\}, \\ \sum_{l=1}^n \dim \ker(A_l - x) - 1, & \text{if for some } l \ker(A_l - x) \neq \{0\}. \end{cases}$$

It is clear that $\mathcal{L}_x \subseteq \ker(A - x)$.

Second, assume that there exist elements $f_l \in \ker(S_l^* - x)$ and $\beta_l \in \mathbb{C}$ such that $((f_l; x f_l); (1; \beta_l)) \in \Gamma_l$ and $\beta_1 + \dots + \beta_n = 0$. Then, clearly, $(f_1, \dots, f_n) \in \ker(A - x)$.

Step 3; Determining the eigenspace: Let $x \in \mathbb{R}$ and $(f_1, \dots, f_n) \in \ker(A - x) \setminus \{0\}$ be given, and let α and β_1, \dots, β_n be the unique numbers with $((f_l; x f_l); (\alpha; \beta_l)) \in \Gamma_l, l = 1, \dots, n$. We distinguish the two cases that $\alpha = 0$ and $\alpha \neq 0$.

Assume that $\alpha = 0$. Then $f_l \in \ker(A_l - x)$ and $\beta_l = \Upsilon_l(f_l; x f_l)$. Hence, in this case, $(f_1, \dots, f_n) \in \mathcal{L}_x$.

Assume that $\alpha \neq 0$. Clearly, $f_l \in \mathcal{N}_{l,x}$ for all l . Moreover, whenever $\text{mul } \Gamma_l = \{0\}$, we must have $f_l \neq 0$. Let us show that

$$\ker(A_l - x) = \{0\}, \quad l = 1, \dots, n.$$

If $f_l = 0$, then $\text{mul } \Gamma_l \neq \{0\}$, and hence $H_l = \{0\}$. If $f_l \neq 0$ and $\ker(A_l - x) \neq \{0\}$, then $f_l \in \ker(A_l - x)$ since $\dim \mathcal{N}_{l,x} \leq 1$. Thus there exists β'_l with $((f_l; x f_l); (0; \beta'_l)) \in \Gamma_l$, and it follows that $\text{mul } \Gamma_l \neq \{0\}$. This contradicts the fact that $f_l \neq 0$.

Next we show (still assuming $\alpha \neq 0$) that

$$\ker(A - x) = \text{span} \{(f_1, \dots, f_n)\}.$$

Let $g = (g_1, \dots, g_n) \in \ker(A - x) \setminus \{0\}$ be given, and let α' and $\beta'_1, \dots, \beta'_n$ be the unique numbers with $((g_l; x g_l); (\alpha'; \beta'_l)) \in \Gamma_l, l = 1, \dots, n$. If $\alpha' = 0$, we would have $\ker(A_l - x) \neq \{0\}$ for at least one index l . This contradicts what we showed in the previous paragraph, and we conclude that $\alpha' \neq 0$. Set $\lambda := \frac{\alpha'}{\alpha}$. Then $((\lambda f_l; x \lambda f_l); (\alpha'; \lambda \beta_l)) \in \Gamma_l$, and it follows from Lemma 3.20, (iii), that $g_l = \lambda f_l$.

Putting together these facts with what we showed in Step 2, we obtain that for each real point x one of the following three alternatives holds:

- (i) $\ker(A - x) = \{0\}$.
- (ii) There exist at least two indices l with $\ker(A_l - x) \neq \{0\}$.
- (iii) We have $\ker(A_l - x) = \{0\}, l = 1, \dots, n$, and $\mathcal{N}_{l,x} \neq \{0\}$ whenever $\text{mul } \Gamma_l = \{0\}$.

If the alternative (ii) takes place, then $\ker(A - x) = \mathcal{L}_x$. If (iii) takes place, then $\dim \ker(A - x) = 1$.

Step 4; Asymptotics of m_l : The characterizations stated in the theorem now follow from standard Weyl function theory (see, e.g., [15, §2,3] for the relation of Weyl functions, measures, and spectral properties). First, a point $x \in \mathbb{R}$ is an eigenvalue of A_l if and only

$$\lim_{\varepsilon \downarrow 0} \varepsilon \text{Im } m_l(x + i\varepsilon) > 0.$$

Next, assume that $\text{mul } \Gamma_l = \{0\}$. Then we have $\mathcal{N}_{l,x} \setminus \ker(A_l - x) \neq \emptyset$ if and only if

$$m_l(x) := \lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon) \in \mathbb{R}, \quad \lim_{\varepsilon \downarrow 0} \frac{1}{i\varepsilon} (m_l(x + i\varepsilon) - m_l(x)) \in (0, \infty).$$

If $\text{mul } \Gamma_l \neq \{0\}$, then m_l is identically equal to a real constant, and hence trivially $\lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon)$ exists in \mathbb{R} and $\lim_{\varepsilon \downarrow 0} \frac{1}{i\varepsilon} (m_l(x + i\varepsilon) - m_l(x)) = 0$. Conversely, if these two relations hold, the function m_l must be a real constant.

Finally, we need to relate the limit $m_l(x)$ with the boundary relation Γ_l under the assumption that this limit at all exists, and that $\lim_{\varepsilon \downarrow 0} \frac{1}{i\varepsilon} (m_l(x + i\varepsilon) - m_l(x)) \in [0, \infty)$. For $\varepsilon > 0$, let $f_{l,\varepsilon}$ and $\beta_{l,\varepsilon}$ be the unique elements with

$$((f_{l,\varepsilon}; (x + i\varepsilon)f_{l,\varepsilon}); (1; \beta_{l,\varepsilon})) \in \Gamma_l.$$

Then, by the definition of m_l , we have $m_l(x + i\varepsilon) = \beta_{l,\varepsilon}$. The abstract Green's identity gives

$$\varepsilon \|f_{l,\varepsilon}\|^2 = \text{Im } m_l(x + i\varepsilon),$$

and hence $\|f_{l,\varepsilon}\|$ remains bounded when ε approaches 0. Let f_l be the weak limit of $f_{l,\varepsilon}$ for $\varepsilon \downarrow 0$. Since Γ_l is a closed linear relation, it is weakly closed, and we obtain

$$((f_l; x f_l); (1; m_l(x))) \in \Gamma_l.$$

This finishes the proof. \square

4. Computation of Rank for Singular Spectrum

The following theorem is our main result, and this section is entirely devoted to its proof. Concerning terminology for boundary relations, remember Definition 3.6.

Theorem 4.1. *Let $n \geq 2$, and let for each $l \in \{1, \dots, n\}$ a closed symmetric simple relation S_l in a Hilbert space H_l and a boundary relation $\Gamma_l \subseteq H_l^2 \times \mathbb{C}^2$ for S_l^* be given. Assume that each Γ_j is of function type, and that at least one Γ_l is a boundary function. Let μ_l be the measure in the Herglotz-integral representation of the Weyl function of Γ_l , set $\mu := \sum_{l=1}^n \mu_l$, and let μ_s be the singular part of μ with respect to the Lebesgue measure.*

Consider the selfadjoint operator A constructed by pasting $\{\Gamma_l : l = 1, \dots, n\}$ with standard interface conditions. Denote by E the projection valued spectral measure of A , let E_s be its singular part with respect to the Lebesgue measure, and let $E_{s,ac}$ and $E_{s,s}$ be the absolutely continuous and singular parts of E_s with respect to μ . Moreover, let N_A be the spectral multiplicity function of A and

$$r(x) := \#\{l \in \{1, \dots, n\} : D_\mu \mu_l(x) > 0\}.$$

Then the following hold:

- (I) $E_{s,ac} \sim \mathbf{1}_{X_{>1}} \cdot \mu_s$ where $X_{>1} := r^{-1}(\{2, \dots, n\})$.
- (II) $N_A(x) = r(x) - 1$ for $E_{s,ac}$ -a.a. points $x \in \mathbb{R}$.
- (III) $N_A(x) = 1$ for $E_{s,s}$ -a.a. points $x \in \mathbb{R}$.

Remark 4.2. We may assume without loss of generality that all boundary relations Γ_l with possible exception of Γ_n are boundary functions: First, reordering the boundary relations Γ_l obviously does not change the relation A . Second, it is easy to see that the pasting with standard interface conditions of a collection as given in the theorem is always unitary equivalent to the pasting with standard interface conditions of a collection with at most one boundary relation being a proper relation (take instead of k pure relations $\{((0;0);(w;\beta_l w)), w \in \mathbb{C}\}, l = l_1, \dots, l_k$, one relation $\{((0;0);(w;\sum_{j=1}^k \beta_{l_j} w)), w \in \mathbb{C}\}$).

Let us recall some notation: The Weyl functions of the boundary relations Γ_l are denoted as m_l , and we set $m := \sum_{l=1}^n m_l$. The boundary relation Γ is as in 3.10, and we denote by $M(z) = (M_{ij}(z))_{i,j=1}^n$ its Weyl function. Explicitly, the function M is given by (3.11). Let $\Omega = (\Omega_{ij})_{i,j=1}^n$ be the $n \times n$ -matrix valued measure in the Herglotz-integral representation (2.9) of M , let ρ be the trace measure $\rho := \text{tr } \Omega$, and let $\omega = (\omega_{ij})_{i,j=1}^n$ be the symmetric derivative of Ω with respect to ρ , i.e.

$$\omega_{ij}(x) := \frac{d\Omega_{ij}}{d\rho}(x), \quad i, j = 1, \dots, n,$$

whenever these derivatives exist.

Moreover, remember 3.15 which says that the actual task is to compute $\text{rank } \omega(x)$.

Stage 1: μ -singular part.

In this part, we prove the following statement.

Proposition 4.3. *If the set $X \subseteq \mathbb{R}$ is μ -zero, then for ρ -a.a. points $x \in X$ the symmetric derivative $\omega(x)$ exists and $\text{rank } \omega(x) = 1$.*

The proof is split into two parts. First, the case when the right lower entry M_{nn} of M dominates.

Lemma 4.4. *Let $x \in \mathbb{R}$, and assume that*

- (i) *The symmetric derivative $\omega(x)$ exists.*
- (ii) $\frac{d\rho}{d\lambda}(x) = \infty$.
- (iii) $\omega_{nn}(x) > 0$.

Then the limits $m_l(x) := \lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon), l = 1, \dots, n$, exist, are real, and

$$m(x) := \lim_{\varepsilon \downarrow 0} m(x + i\varepsilon) = 0.$$

The rank of the matrix $\omega(x)$ is equal to one.

Proof. Step 1. Existence of limits. The present hypotheses imply that

$$\frac{d\Omega_{nn}}{d\lambda}(x) = \frac{d\Omega_{nn}}{d\rho}(x) \cdot \frac{d\rho}{d\lambda}(x) = \infty, \tag{4.1}$$

$$\frac{d\Omega_{ij}}{d\Omega_{nn}}(x) = \frac{d\Omega_{ij}}{d\rho}(x) \cdot \frac{d\rho}{d\Omega_{nn}}(x) = \frac{\omega_{ij}(x)}{\omega_{nn}(x)}. \tag{4.2}$$

Applying Theorem 2.12 with Ω_{nn} gives $\lim_{\varepsilon \downarrow 0} \operatorname{Im} M_{nn}(x + i\varepsilon) = \infty$. However, $M_{nn} = -\frac{1}{m}$, and we thus have

$$\lim_{\varepsilon \downarrow 0} \left| \operatorname{Im} \frac{1}{m(x + i\varepsilon)} \right| = \infty.$$

In particular, $\lim_{\varepsilon \downarrow 0} m(x + i\varepsilon) = 0$. Since $\operatorname{Im} m_l$ is nonnegative throughout the upper half-plane, this implies that also $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m_l(x + i\varepsilon) = 0$, $l = 1, \dots, n$.

In order to capture the behavior of the real parts, we apply Theorem 2.13 with the measures Ω_{ij} and Ω_{nn} . This gives

$$\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} M_{ij}(x + i\varepsilon)}{\operatorname{Im} M_{nn}(x + i\varepsilon)} = \frac{\omega_{ij}(x)}{\omega_{nn}(x)}.$$

Let $l \in \{1, \dots, n-1\}$. For each $z \in \mathbb{C}^+$, we have

$$\frac{\operatorname{Im} M_{ln}(z)}{\operatorname{Im} M_{nn}(z)} = \frac{\operatorname{Im} \frac{m_l(z)}{m(z)}}{\operatorname{Im} \frac{1}{m(z)}} = \operatorname{Re} m_l(z) + \frac{\operatorname{Im} m_l(z) \operatorname{Re} \frac{1}{m(z)}}{\operatorname{Im} \frac{1}{m(z)}}, \quad (4.3)$$

and

$$\left| \operatorname{Im} m_l(z) \operatorname{Re} \frac{1}{m(z)} \right| \leq \operatorname{Im} m_l(z) \left| \frac{1}{m(z)} \right| \leq \frac{\operatorname{Im} m_l(z)}{\operatorname{Im} m(z)} \leq 1.$$

Hence, the limit of $\operatorname{Re} m_l$ exists, in fact,

$$\lim_{\varepsilon \downarrow 0} \operatorname{Re} m_l(x + i\varepsilon) = \frac{\omega_{ln}(x)}{\omega_{nn}(x)}.$$

Since we already know that imaginary parts tend to zero, thus

$$\lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon) = \frac{\omega_{ln}(x)}{\omega_{nn}(x)}, \quad l = 1, \dots, n-1.$$

Since m tends to zero, it follows that $\lim_{\varepsilon \downarrow 0} m_n(x + i\varepsilon) = -\sum_{l=1}^{n-1} \frac{\omega_{ln}(x)}{\omega_{nn}(x)}$.

Step 2. Computing rank. Let $l \in \{1, \dots, n-1\}$. For each $z \in \mathbb{C}^+$, we have

$$\begin{aligned} \operatorname{Im} M_{ll}(z) &= \operatorname{Im} \frac{m_l(z)[m(z) - m_l(z)]}{m(z)} \\ &= \operatorname{Im} \frac{1}{m(z)} \operatorname{Re} (m_l(z)[m(z) - m_l(z)]) \\ &\quad + \operatorname{Re} \frac{1}{m(z)} \operatorname{Im} (m_l(z)[m(z) - m_l(z)]), \end{aligned} \quad (4.4)$$

and

$$\begin{aligned}
 & \left| \operatorname{Re} \frac{1}{m(z)} \operatorname{Im} (m_l(z)[m(z) - m_l(z)]) \right| \\
 & \leq \left| \operatorname{Re} \frac{1}{m(z)} \cdot \operatorname{Im} m_l(z) \cdot \operatorname{Re} (m(z) - m_l(z)) \right| \\
 & \quad + \left| \operatorname{Re} \frac{1}{m(z)} \cdot \operatorname{Re} m_l(z) \cdot \operatorname{Im} (m(z) - m_l(z)) \right| \\
 & \leq \left| \operatorname{Re} \frac{1}{m(z)} \right| \cdot \operatorname{Im} m(z) \cdot \sum_{j=1}^n |\operatorname{Re} m_j(z)| \\
 & \leq \frac{\operatorname{Im} m(z)}{|m(z)|} \cdot \sum_{j=1}^n |\operatorname{Re} m_j(z)| \leq \sum_{j=1}^n |\operatorname{Re} m_j(z)|.
 \end{aligned}$$

Since $m_j(x + i\varepsilon)$ approaches a real limit when $\varepsilon \downarrow 0$, and $m(x + i\varepsilon)$ tends to zero, and $|\operatorname{Im} \frac{1}{m(x+i\varepsilon)}|$ tends to infinity, we obtain (here “ $o(1)$ ” is understood as a term which may depend on ε but tends to zero if $\varepsilon \downarrow 0$)

$$\operatorname{Im} M_{ll}(x + i\varepsilon) = -\operatorname{Im} \frac{1}{m(x + i\varepsilon)} \cdot (m_l^2(x) + o(1)), \quad l = 1, \dots, n-1.$$

Arguing analogously, we obtain

$$\operatorname{Im} M_{lk}(x + i\varepsilon) = -\operatorname{Im} \frac{1}{m(x + i\varepsilon)} \cdot (m_l(x)m_k(x) + o(1)),$$

$$l, k = 1, \dots, n-1, \quad l \neq k,$$

$$\operatorname{Im} M_{ln}(x + i\varepsilon) = -\operatorname{Im} \frac{1}{m(x + i\varepsilon)} \cdot (m_l(x) + o(1)), \quad l = 1, \dots, n-1.$$

Therefore

$$\operatorname{Im} \operatorname{tr} M(x + i\varepsilon) = -\operatorname{Im} \frac{1}{m(x + i\varepsilon)} \left(1 + \sum_{l=1}^{n-1} m_l^2(x) + o(1) \right),$$

and hence, referring again to Theorem 2.13,

$$\omega(x) = \frac{1}{1 + \sum_{l=1}^{n-1} m_l^2(x)} \times \begin{pmatrix} m_1^2(x) & m_2(x)m_1(x) & \cdots & m_{n-1}(x)m_1(x) & m_1(x) \\ m_1(x)m_2(x) & m_2^2(x) & \cdots & m_{n-1}(x)m_2(x) & m_2(x) \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ m_1(x)m_{n-1}(x) & m_2(x)m_{n-1}(x) & \cdots & m_{n-1}^2(x) & m_{n-1}(x) \\ m_1(x) & m_2(x) & \cdots & m_{n-1}(x) & 1 \end{pmatrix}$$

Obviously, the rank of this matrix is 1. □

Second, the case that the right lower entry of M does not dominate. In this case, a more refined argument is necessary. First, two technical observations.

Lemma 4.5. *Let $D \subseteq \mathbb{C}$ be a connected set with $x \in \bar{D}$. Moreover, let $f, g : D \rightarrow \mathbb{C}^+$ be continuous functions. If $\lim_{t \rightarrow x} \frac{f(t)}{g(t)} = -1$, then*

$$\lim_{t \rightarrow x} \frac{\operatorname{Im} f(t)}{\operatorname{Re} f(t)} = \lim_{t \rightarrow x} \frac{\operatorname{Im} g(t)}{\operatorname{Re} g(t)} = 0. \quad (4.5)$$

Proof. For $z \in \mathbb{C}^+$, let $\arg z$ denote the branch of the argument of z in $[0, \pi]$. Then $\arg f$ and $\arg g$ are continuous functions. We have

$$\lim_{t \rightarrow x} [\arg f(t) - \arg g(t)] = \pi \pmod{2\pi},$$

and hence either

$$\lim_{t \rightarrow x} \arg f(t) = \pi \quad \text{and} \quad \lim_{t \rightarrow x} \arg g(t) = 0,$$

or

$$\lim_{t \rightarrow x} \arg f(t) = 0 \quad \text{and} \quad \lim_{t \rightarrow x} \arg g(t) = \pi.$$

In both cases, (4.5) follows. \square

Lemma 4.6. *Let $\alpha \in \mathbb{R}$, and let $\{f_j\}_{j \in \mathbb{N}}$ and $\{g_j\}_{j \in \mathbb{N}}$ be sequences of complex numbers. Assume that*

- (i) *For each $j \in \mathbb{N}$ we have $\operatorname{Im} g_j \neq 0$.*
- (ii) $\lim_{j \rightarrow \infty} \frac{f_j}{g_j} = \alpha$.
- (iii) $\frac{\operatorname{Re} g_j}{\operatorname{Im} g_j} = O(1)$ as $j \rightarrow \infty$.

Then $\lim_{j \rightarrow \infty} \frac{\operatorname{Im} f_j}{\operatorname{Im} g_j} = \alpha$.

Proof. Let $\varepsilon > 0$. Then, for sufficiently large indices j , we have

$$\left| \frac{f_j}{g_j} - \alpha \right| < \varepsilon.$$

This implies that $|\operatorname{Im} f_j - \alpha \operatorname{Im} g_j| \leq |f_j - \alpha g_j| < \varepsilon |g_j|$, and hence

$$\left| \frac{\operatorname{Im} f_j}{\operatorname{Im} g_j} - \alpha \right| < \varepsilon \frac{|g_j|}{|\operatorname{Im} g_j|}.$$

Due to our assumption (iii), the right-hand side of this estimate can be made arbitrarily small for large indices j . \square

Now we are ready to settle the case that M_{nn} does not dominate.

Lemma 4.7. *Let $x \in \mathbb{R}$, and assume that*

- (i) *The symmetric derivative $\omega(x)$ exists.*
- (ii) $\frac{d\rho}{d\lambda}(x) = \infty$.
- (iii) $\omega_{nn}(x) = 0$.
- (iv) $\frac{d\mu}{d\rho}(x) = 0$.
- (v) *There exists no $k \in \{1, \dots, n\}$ with $\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Re} M_{kk}(x+i\varepsilon)}{\operatorname{Im} \operatorname{tr} M(x+i\varepsilon)} = \infty$.*

Then the rank of the matrix $\omega(x)$ is equal to one.

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Proof. Since $w_{nn}(x) = 0$, there exists an index $k \in \{1, \dots, n-1\}$ such that $\omega_{kk}(x) > 0$. From this, the condition (ii), and Theorem 2.13, it follows that

$$\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} M_{kk}(x + i\varepsilon)}{\operatorname{Im} \operatorname{tr} M(x + i\varepsilon)} = \omega_{kk}(x) > 0. \quad (4.6)$$

Throughout the proof we fix an index k with this property.

Step 1: In this step we deduce that

$$\lim_{\varepsilon \downarrow 0} \frac{m(x + i\varepsilon)}{m_k(x + i\varepsilon)} = 0. \quad (4.7)$$

By Theorem 2.13, the present hypotheses (iv) and (ii) imply that

$$\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} m(x + i\varepsilon)}{\operatorname{Im} \operatorname{tr} M(x + i\varepsilon)} = 0.$$

Therefore, using (4.6), one has $\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} m(x+i\varepsilon)}{\operatorname{Im} M_{kk}(x+i\varepsilon)} = 0$. We compute

$$\begin{aligned} \operatorname{Im} M_{kk} &= \operatorname{Im} \frac{m_k(m - m_k)}{m} = \frac{\operatorname{Im}[m_k(m - m_k)\bar{m}]}{|m|^2} \\ &= \frac{\operatorname{Im}[m_k(m - m_k)(\bar{m} - \bar{m}_k + \bar{m}_k)]}{|m|^2} \\ &= \left| \frac{m - m_k}{m} \right|^2 \operatorname{Im} m_k + \left| \frac{m_k}{m} \right|^2 \operatorname{Im}(m - m_k). \end{aligned} \quad (4.8)$$

From this we have

$$\frac{\operatorname{Im} m}{\operatorname{Im} M_{kk}} = \frac{1}{\left| \frac{m_k}{m} \right|^2 \frac{\operatorname{Im}(m - m_k)}{\operatorname{Im} m} + \left| \frac{m - m_k}{m} \right|^2 \frac{\operatorname{Im} m_k}{\operatorname{Im} m}}.$$

Using the estimate

$$\begin{aligned} \left| \frac{m_k}{m} \right|^2 \frac{\operatorname{Im}(m - m_k)}{\operatorname{Im} m} + \left| \frac{m - m_k}{m} \right|^2 \frac{\operatorname{Im} m_k}{\operatorname{Im} m} &\leq \left| \frac{m_k}{m} \right|^2 + \left| \frac{m - m_k}{m} \right|^2 \\ &\leq 2 \left| \frac{m_k}{m} \right|^2 + 2 \left| \frac{m_k}{m} \right| + 1 \end{aligned}$$

we see that the limit relation $\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} m(x+i\varepsilon)}{\operatorname{Im} M_{kk}(x+i\varepsilon)} = 0$ implies (4.7).

In addition, further rewriting (4.7) as $\lim_{\varepsilon \downarrow 0} \frac{m(x+i\varepsilon) - m_k(x+i\varepsilon)}{m_k(x+i\varepsilon)} = -1$, we get from Lemma 4.5 that

$$\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} m_k(x + i\varepsilon)}{\operatorname{Re} m_k(x + i\varepsilon)} = 0, \quad \lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im}[m(x + i\varepsilon) - m_k(x + i\varepsilon)]}{\operatorname{Re}[m(x + i\varepsilon) - m_k(x + i\varepsilon)]} = 0.$$

Step 2: In this step we show that for each $l \in \{1, \dots, n-1\}$

$$a_l(x) := \lim_{\varepsilon \downarrow 0} \frac{m_l(x + i\varepsilon)}{m_k(x + i\varepsilon)} \quad \text{exists in } \mathbb{R}.$$

Let $l \in \{1, \dots, n-1\}$ be given. By (i), (ii), and Theorem 2.13, we have $\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} M_{ll}(x+i\varepsilon)}{\operatorname{Im} \operatorname{tr} M(x+i\varepsilon)} = \omega_{ll}(x)$. Using (4.6), thus $\lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} M_{ll}(x+i\varepsilon)}{\operatorname{Im} M_{kk}(x+i\varepsilon)} = \frac{\omega_{ll}(x)}{\omega_{kk}(x)}$.

From the computation (4.8), it follows that

$$\frac{\operatorname{Im} M_{ll}(z)}{\operatorname{Im} M_{kk}(z)} = \frac{\left| \frac{m_l(z)}{m_k(z)} \right|^2 \frac{\operatorname{Im}(m(z) - m_l(z))}{\operatorname{Im} m(z)} + \left| \frac{m_l(z)}{m_k(z)} - \frac{m(z)}{m_k(z)} \right|^2 \frac{\operatorname{Im} m_l(z)}{\operatorname{Im} m(z)}}{\frac{\operatorname{Im}(m(z) - m_k(z))}{\operatorname{Im} m(z)} + \left| 1 - \frac{m(z)}{m_k(z)} \right|^2 \frac{\operatorname{Im} m_k(z)}{\operatorname{Im} m(z)}}.$$

Due to (4.7), the denominator tends to 1 when $z = x + i\varepsilon$ and $\varepsilon \downarrow 0$. Therefore, the numerator has the limit $\frac{\omega_{ll}(x)}{\omega_{kk}(x)}$, i.e.

$$\lim_{\varepsilon \downarrow 0} \left[\left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 + \left(\left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} - \frac{m(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 - \left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 \right) \frac{\operatorname{Im} m_l(x+i\varepsilon)}{\operatorname{Im} m(x+i\varepsilon)} \right] = \frac{\omega_{ll}(x)}{\omega_{kk}(x)}. \quad (4.9)$$

Let us show that $\lim_{\varepsilon \downarrow 0} \left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 = \frac{\omega_{ll}(x)}{\omega_{kk}(x)}$. First, if $\frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)}$ were not bounded as $\varepsilon \downarrow 0$, then there would exist a sequence $\{\varepsilon_j\}_{j \in \mathbb{N}}$ with $\varepsilon_j \downarrow 0$, such that $\lim_{j \rightarrow \infty} \left| \frac{m_l(x+i\varepsilon_j)}{m_k(x+i\varepsilon_j)} \right| = \infty$. We have

$$\left| \frac{m_l}{m_k} - \frac{m}{m_k} \right|^2 - \left| \frac{m_l}{m_k} \right|^2 = \left| \frac{m_l}{m_k} \right|^2 \cdot \left(\left| 1 - \frac{m}{m_l} \right| - 1 \right),$$

and it follows that

$$\begin{aligned} & \left| \frac{m_l(x+i\varepsilon_j)}{m_k(x+i\varepsilon_j)} - \frac{m(x+i\varepsilon_j)}{m_k(x+i\varepsilon_j)} \right|^2 - \left| \frac{m_l(x+i\varepsilon_j)}{m_k(x+i\varepsilon_j)} \right|^2 \\ &= o\left(\left| \frac{m_l(x+i\varepsilon_j)}{m_k(x+i\varepsilon_j)} \right|^2 \right) \quad \text{as } j \rightarrow \infty. \end{aligned} \quad (4.10)$$

Since $0 < \frac{\operatorname{Im} m_l(x+i\varepsilon)}{\operatorname{Im} m(x+i\varepsilon)} \leq 1$, it follows that the expression on the left side of (4.9) would also be unbounded, a contradiction. This means that $\frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)}$ remains bounded when $\varepsilon \downarrow 0$. We have

$$\begin{aligned} & \left| \left| \frac{m_l}{m_k} - \frac{m}{m_k} \right|^2 - \left| \frac{m_l}{m_k} \right|^2 \right| = \left| \left| \frac{m_l}{m_k} - \frac{m}{m_k} \right| - \left| \frac{m_l}{m_k} \right| \right| \left(\left| \frac{m_l}{m_k} - \frac{m}{m_k} \right| + \left| \frac{m_l}{m_k} \right| \right) \\ & \leq \left| \frac{m}{m_k} \right| \left(\left| \frac{m_l}{m_k} - \frac{m}{m_k} \right| + \left| \frac{m_l}{m_k} \right| \right), \end{aligned}$$

and it follows that

$$\left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} - \frac{m(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 - \left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 = o(1) \quad \text{as } \varepsilon \downarrow 0.$$

Now we get from (4.9) that

$$\lim_{\varepsilon \downarrow 0} \left| \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right|^2 = \frac{\omega_{ll}(x)}{\omega_{kk}(x)}.$$

Next, rewrite

$$\operatorname{Im} \left(\frac{m_l}{m_k} \right) = \frac{\operatorname{Im} m_l}{|m_k|} \cdot \frac{\operatorname{Re} m_k}{|m_k|} - \frac{\operatorname{Re} m_l}{|m_k|} \cdot \frac{\operatorname{Im} m_k}{|m_k|}.$$

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Since

$$\begin{aligned} \frac{|\operatorname{Re} m_k|}{|m_k|} &\leq 1, & \frac{|\operatorname{Re} m_l|}{|m_k|} &\leq \left| \frac{m_l}{m_k} \right| = O(1), \\ \frac{\operatorname{Im} m_l}{|m_k|}, \frac{\operatorname{Im} m_k}{|m_k|} &\leq \frac{\operatorname{Im} m}{|m_k|} \leq \left| \frac{m}{m_k} \right| = o(1), \end{aligned}$$

we have $\lim_{\varepsilon \downarrow 0} \operatorname{Im} \left(\frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right) = 0$. Therefore

$$\text{either } \lim_{\varepsilon \downarrow 0} \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} = \sqrt{\frac{\omega_{ll}(x)}{\omega_{kk}(x)}} \quad \text{or} \quad \lim_{\varepsilon \downarrow 0} \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} = -\sqrt{\frac{\omega_{ll}(x)}{\omega_{kk}(x)}}.$$

In both cases the limit $\lim_{\varepsilon \downarrow 0} \frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)}$ exists and is real.

Step 3; Computing rank: Let $l, p \in \{1, \dots, n-1\}$. If $l \neq p$, then

$$\lim_{\varepsilon \downarrow 0} \frac{M_{lp}(x+i\varepsilon)}{M_{kk}(x+i\varepsilon)} = \lim_{\varepsilon \downarrow 0} \frac{\frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \cdot \frac{m_p(x+i\varepsilon)}{m_k(x+i\varepsilon)}}{1 - \frac{m(x+i\varepsilon)}{m_k(x+i\varepsilon)}} = a_l(x)a_p(x). \quad (4.11)$$

If $l = p$, then

$$\lim_{\varepsilon \downarrow 0} \frac{M_{ll}(x+i\varepsilon)}{M_{kk}(x+i\varepsilon)} = \lim_{\varepsilon \downarrow 0} \frac{\frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} \cdot \left(\frac{m_l(x+i\varepsilon)}{m_k(x+i\varepsilon)} - \frac{m(x+i\varepsilon)}{m_k(x+i\varepsilon)} \right)}{1 - \frac{m(x+i\varepsilon)}{m_k(x+i\varepsilon)}} = a_l^2(x). \quad (4.12)$$

From (v) it follows that there exists a sequence $\{\varepsilon_j\}_{j \in \mathbb{N}}$ with $\varepsilon_j \downarrow 0$, such that $\lim_{j \rightarrow \infty} \frac{\operatorname{Re} M_{kk}(x+i\varepsilon_j)}{\operatorname{Im} M_{kk}(x+i\varepsilon_j)} = O(1)$ as $j \rightarrow \infty$. Applying Lemma 4.6, we get from (4.11) and (4.12) that

$$\frac{\operatorname{Im} M_{lp}(x+i\varepsilon_j)}{\operatorname{Im} M_{kk}(x+i\varepsilon_j)} = a_l(x)a_p(x), \quad l, p = 1, \dots, n-1.$$

Since $\omega(x)$ is positive semidefinite, (iii) implies that

$$\omega_{ij}(x) = 0, \quad i = n \text{ or } j = n.$$

Altogether,

$$\omega(x) = \frac{1}{\sum_{l=1}^{n-1} a_l^2(x)} \begin{pmatrix} a_1^2(x) & a_2(x)a_1(x) & \cdots & a_{n-1}(x)a_1(x) & 0 \\ a_1(x)a_2(x) & a_2^2(x) & \cdots & a_{n-1}(x)a_2(x) & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ a_1(x)a_{n-1}(x) & a_2(x)a_{n-1}(x) & \cdots & a_{n-1}^2(x) & 0 \\ 0 & 0 & \cdots & 0 & 0 \end{pmatrix}.$$

The rank of this matrix obviously cannot exceed 1. However, $a_k(x) = 1$, and hence it is nonzero. \square

Having available Lemma 4.4 and Lemma 4.7, it is not difficult to prove Proposition 4.3.

Proof of Proposition 4.3. Let a μ -zero set $X \subseteq \mathbb{R}$ be given. It is enough to show that the conditions (i), (ii) appearing in Lemma 4.4 and Lemma 4.7, and the conditions (iv), (v) appearing in Lemma 4.7 are satisfied ρ -a.e. on X .

The fact that the symmetric derivative $\omega(x)$ exists ρ -a.e., has been noted in Remark 2.6. Denote by μ_{ac} and ρ_{ac} the absolutely continuous parts of μ and ρ with respect to the Lebesgue measure. By Theorem 1.1, $\rho_{ac} \sim \mu_{ac} \ll \mu$. Thus the set X is ρ_{ac} -zero. Corollary 2.5, (iv), says that for ρ_s -a.a. points $x \in \mathbb{R}$ one has $\frac{d\rho}{d\lambda}(x) = \infty$. Since X is ρ_{ac} -zero, one has $\frac{d\rho}{d\lambda}(x) = \infty$ not only for ρ_s -a.a., but even for ρ -a.a. $x \in X$.

Corollary 2.4, (ii), shows that for ρ -a.a. points $x \in X$ we have $\frac{d\mu}{d\rho}(x) = 0$. Corollary 2.16 applied with the measure ρ and the measures that correspond to the Herglotz functions M_{ll} gives

$$\frac{|\operatorname{Re} M_{ll}(x + i\varepsilon)|}{\operatorname{Im} \operatorname{tr} M(x + i\varepsilon)} \rightarrow \infty \text{ as } \varepsilon \downarrow 0, \quad l = 1, \dots, n-1,$$

for ρ -a.a. $x \in \mathbb{R}$. □

Stage 2: μ -Absolutely Continuous Part.

Consider the sets

$$X_{reg} := \left\{ x \in \mathbb{R} : \frac{d\mu}{d\lambda}(x) = \infty, \frac{d\rho}{d\mu}(x) \in [0, \infty), \right. \\ \left. \frac{|\operatorname{Re} m(x + i\varepsilon)|}{\operatorname{Im} m(x + i\varepsilon)} \rightarrow \infty \text{ as } \varepsilon \downarrow 0, \right. \\ \left. \forall l = 1, \dots, n : \frac{d\mu_l}{d\mu}(x) \text{ exists, and } \lim_{\varepsilon \downarrow 0} \frac{m_l(x + i\varepsilon)}{m(x + i\varepsilon)} = \frac{d\mu_l}{d\mu}(x) \right\}. \\ X_{reg}^1 := \left\{ x \in X_{reg} : \exists l \in \{1, \dots, n\} : \frac{d\mu_l}{d\mu}(x) = 1 \right\}, \\ X_{reg}^{>1} := \left\{ x \in X_{reg} : \nexists l \in \{1, \dots, n\} : \frac{d\mu_l}{d\mu}(x) = 1 \right\}.$$

Note that, since $\sum_{l=1}^n \frac{d\mu_l}{d\mu}(x) = 1$ and $\frac{d\mu_l}{d\mu}(x) \in [0, 1]$, one of the following alternatives takes place:

- (1) There exists one index k_0 with $\frac{d\mu_{k_0}}{d\mu}(x) = 1$, and for all other indices $k \neq k_0$ we have $\frac{d\mu_k}{d\mu}(x) = 0$.
- (>1) For all indices k we have $\frac{d\mu_k}{d\mu}(x) < 1$, and there exist at least two indices k with $\frac{d\mu_k}{d\mu}(x) > 0$.

In this part we show the following statement.

Proposition 4.8. *The following hold:*

- (i) *The set X_{reg}^1 is ρ -zero.*
- (ii) *For ρ -a.a. points $x \in X_{reg}$ the symmetric derivative $\omega(x)$ exists and*

$$\operatorname{rank} \omega(x) = r(x) - 1.$$

First, an elementary fact which we use to compute rank.

Spectral Multiplicity of Schrödinger Operators

Lemma 4.9. *Let $n \in \mathbb{N}$, let $b_1, b_2, \dots, b_n, d \in \mathbb{R} \setminus \{0\}$, and consider the matrix*

$$M_d := \begin{pmatrix} b_1(d - b_1) & -b_2b_1 & \cdots & -b_nb_1 \\ -b_1b_2 & b_2(d - b_2) & \cdots & -b_nb_2 \\ \vdots & \vdots & \ddots & \vdots \\ -b_1b_n & -b_2b_n & \cdots & b_n(d - b_n) \end{pmatrix}.$$

If $d = \sum_{l=1}^n b_l$, then $\text{rank } M_d = n - 1$. Otherwise $\text{rank } M_d = n$.

Proof. We have:

$$\begin{aligned} \text{rank } M_d &= \text{rank} \begin{pmatrix} d - b_1 & -b_1 & \cdots & -b_1 \\ -b_2 & d - b_2 & \cdots & -b_2 \\ \vdots & \vdots & \ddots & \vdots \\ -b_n & -b_n & \cdots & d - b_n \end{pmatrix} \\ &= \text{rank} \begin{pmatrix} d & 0 & \cdots & -b_1 \\ 0 & d & \cdots & -b_2 \\ \vdots & \vdots & \ddots & \vdots \\ -d & -d & \cdots & d - b_n \end{pmatrix} = \text{rank} \begin{pmatrix} 1 & 0 & \cdots & -b_1 \\ 0 & 1 & \cdots & -b_2 \\ \vdots & \vdots & \ddots & \vdots \\ -1 & -1 & \cdots & d - b_n \end{pmatrix} \\ &= \text{rank} \begin{pmatrix} 1 & 0 & \cdots & -b_1 \\ 0 & 1 & \cdots & -b_2 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & d - \sum_{l=1}^n b_l \end{pmatrix} = \begin{cases} n - 1, & d - \sum_{l=1}^n b_l = 0, \\ n, & d - \sum_{l=1}^n b_l \neq 0. \end{cases} \end{aligned}$$

The next two lemmata contain the essential arguments. □

Lemma 4.10. *For each $x \in X_{reg}$ we have*

$$\frac{d\rho}{d\mu}(x) = \sum_{k=1}^{n-1} \frac{d\mu_k}{d\mu}(x) \left(1 - \frac{d\mu_k}{d\mu}(x)\right).$$

Proof. Since $\frac{d\mu}{d\lambda}(x) = \infty$, Theorem 2.12 gives $\lim_{\varepsilon \downarrow 0} \text{Im } m(x + i\varepsilon) = \infty$. This implies that

$$\lim_{\varepsilon \downarrow 0} \frac{\text{Im } M_{nn}(x + i\varepsilon)}{\text{Im } m(x + i\varepsilon)} = \lim_{\varepsilon \downarrow 0} \frac{1}{|m(x + i\varepsilon)|^2} = 0.$$

Let $l \in \{1, \dots, n - 1\}$. Then we have

$$\lim_{\varepsilon \downarrow 0} \frac{M_{ll}(x + i\varepsilon)}{m(x + i\varepsilon)} = \frac{d\mu_l}{d\mu}(x) \left(1 - \frac{d\mu_l}{d\mu}(x)\right).$$

Choose a sequence $\{\varepsilon_j\}_{j \in \mathbb{N}}, \varepsilon_j \downarrow 0$, such that $\frac{|\operatorname{Re} m(x+i\varepsilon_j)|}{\operatorname{Im} m(x+i\varepsilon_j)}$ remains bounded when $j \rightarrow \infty$. Lemma 4.6 yields that also

$$\lim_{j \rightarrow \infty} \frac{\operatorname{Im} M_{ll}(x+i\varepsilon_j)}{\operatorname{Im} m(x+i\varepsilon_j)} = \frac{d\mu_l}{d\mu}(x) \left(1 - \frac{d\mu_l}{d\mu}(x)\right), \quad l = 1, \dots, n. \quad (4.13)$$

Now we use the sequence $\{\varepsilon_j\}_{j \in \mathbb{N}}$ to evaluate $\frac{d\rho}{d\mu}(x)$ by means of Theorem 2.13. This gives

$$\frac{d\rho}{d\mu}(x) = \lim_{j \rightarrow \infty} \frac{\operatorname{Im} \operatorname{tr} M(x+i\varepsilon_j)}{\operatorname{Im} m(x+i\varepsilon_j)} = \sum_{l=1}^{n-1} \frac{d\mu_l}{d\mu}(x) \left(1 - \frac{d\mu_l}{d\mu}(x)\right). \quad (4.14)$$

□

Corollary 4.11. *Let $x \in X_{reg}$. Then*

$$\frac{d\rho}{d\mu}(x) \begin{cases} = 0, & x \in X_{reg}^1, \\ > 0, & x \in X_{reg}^{>1}. \end{cases}$$

Proof. Let $x \in X_{reg}$. Then we have $x \in X_{reg}^1$ if the above alternative (1) takes place, and $x \in X_{reg}^{>1}$ if (> 1) takes place. Hence, for $x \in X_{reg}^1$ we have $\sum_{l=1}^{n-1} \frac{d\mu_k}{d\mu}(x) \left(1 - \frac{d\mu_k}{d\mu}(x)\right) = 0$, and for $x \in X_{reg}^{>1}$ this sum is positive. □

Lemma 4.12. *Let $x \in X_{reg}^{>1}$, and assume that*

- (i) *The symmetric derivative $\omega(x)$ exists.*
- (ii) $\frac{d\rho}{d\lambda}(x) = \infty$.

Then

$$\operatorname{rank} \omega(x) = \#\left\{l \in \{1, \dots, n\} : \frac{d\mu_l}{d\mu}(x) > 0\right\} - 1. \quad (4.15)$$

Proof. By Theorem 2.13, the present assumptions (i) and (ii) ensure that

$$\omega(x) = \lim_{\varepsilon \downarrow 0} \frac{\operatorname{Im} M(x+i\varepsilon)}{\operatorname{Im} \operatorname{tr} M(x+i\varepsilon)}. \quad (4.16)$$

It is easy to show that the last row and column of the matrix $\omega(x)$ vanishes: By Theorem 2.12, our assumption (ii) gives $\lim_{\varepsilon \downarrow 0} \operatorname{Im} \operatorname{tr} M(x+i\varepsilon) = \infty$. We already saw in the proof of the last lemma that $\lim_{\varepsilon \downarrow 0} \operatorname{Im} m(x+i\varepsilon) = \infty$, and it follows that $\lim_{\varepsilon \downarrow 0} \operatorname{Im} M_{nn}(x+i\varepsilon) = 0$ and $\omega_{nn}(x) = 0$. Since $\omega(x)$ is positive semidefinite, all entries $\omega_{ij}(x)$ with $i = n$ or $j = n$ must vanish.

To shorten notation, set $d_l(x) := \frac{d\mu_l}{d\mu}(x)$. Let $\{\varepsilon_j\}_{j \in \mathbb{N}}$ be the same sequence as in the proof of the previous lemma. Then not only (4.13) holds, but also

$$\lim_{j \rightarrow \infty} \frac{\operatorname{Im} M_{lk}(x+i\varepsilon_j)}{\operatorname{Im} m(x+i\varepsilon_j)} = -d_l(x)d_k(x), \quad l, k = 1, \dots, n-1, l \neq k.$$

Referring to (4.16) and (4.14), we obtain

$$\omega(x) = \frac{1}{\sum_{l=1}^{n-1} d_l(x)(1-d_l(x))} \times \begin{pmatrix} d_1(x)(1-d_1(x)) & -d_2(x)d_1(x) & \cdots & -d_{n-1}(x)d_1(x) & 0 \\ -d_1(x)d_2(x) & d_2(x)(1-d_2(x)) & \cdots & -d_{n-1}(x)d_2(x) & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ -d_1(x)d_{n-1}(x) & -d_2(x)d_{n-1}(x) & \cdots & d_{n-1}(x)(1-d_{n-1}(x)) & 0 \\ 0 & 0 & \cdots & 0 & 0 \end{pmatrix}.$$

Applying Lemma 4.9 with the matrix obtained from $\omega(x)$ by deleting all rows and columns which contain only zeros, gives

$$\text{rank } \omega(x) = \begin{cases} \#\{l : 1 \leq l \leq n-1, d_l(x) > 0\}, & \sum_{l=1}^{n-1} d_l(x) \neq 1, \\ \#\{l : 1 \leq l \leq n-1, d_l(x) > 0\} - 1, & \sum_{l=1}^{n-1} d_l(x) = 1. \end{cases}$$

The condition $\sum_{l=1}^{n-1} d_l(x) \neq 1$ is equivalent to $d_n(x) \neq 0$, and the formula (4.15) follows. \square

Proof of Proposition 4.8. Corollary 4.11 and Corollary 2.4, (i), show that X_{reg}^1 is ρ -zero. Denote

$$X_{reg}^+ := \left\{ x \in X_{reg}^{>1} : \omega(x) \text{ exists, } \frac{d\rho}{d\lambda}(x) = \infty \right\}.$$

Then Lemma 4.12 says that

$$\text{rank } \omega(x) = \#\left\{ l \in \{1, \dots, n\} : \frac{d\mu_l}{d\mu}(x) > 0 \right\} - 1, \quad x \in X_{reg}^+. \quad (4.17)$$

We have

$$\begin{aligned} X_{reg} \setminus X_{reg}^+ &\subseteq \{x \in \mathbb{R} : \omega(x) \text{ does not exist}\} \cup \{x \in \mathbb{R} : \frac{d\rho}{d\lambda}(x) \in [0, \infty)\} \\ &\cup \{x \in \mathbb{R} : \frac{d\rho}{d\lambda}(x) \text{ does not exist}\} \cup X_{reg}^1. \end{aligned}$$

The first set in this union is ρ -zero by Remark 2.6. By Corollary 2.5, (iv), the second set is ρ_s -zero. The third set is ρ -zero by Theorem 2.1, and the last by the already proved item (i). Together, and due to the fact that the set X_{reg} itself is Lebesgue-zero, we see that $X_{reg} \setminus X_{reg}^+$ is ρ -zero.

We have

$$X_{reg} \subseteq \left\{ x \in \mathbb{R} \setminus \mathcal{E}_{\rho, \mu} : \frac{d\rho}{d\mu}(x) \in [0, \infty) \right\} \cup \mathcal{E}_{\rho, \mu}.$$

Using Corollary 2.4, (iii), we obtain that the intersection of every μ -zero set with X_{reg} is ρ -zero. Hence, $D_\mu \mu_l(x) = \frac{d\mu_l}{d\mu}(x)$ for ρ -a.a. $x \in X_{reg}$, and (4.17) implies that

$$\text{rank } \omega(x) = r(x) - 1, \quad \text{for } \rho\text{-a.a. } x \in X_{reg}. \quad \square$$

Stage 3: Finishing the proof of the main theorem.

Having available Propositions 4.3 and 4.8, it is not difficult anymore to complete the proof of Theorem 4.1.

Proof of Theorem 4.1. Let $\mu = \mu_{ac} + \mu_s$ and $\rho = \rho_{ac} + \rho_s$ be the decompositions of μ and ρ with respect to the Lebesgue measure.

The set X_{reg} is μ_{ac} -zero, because (as we know from Corollary 2.5, (iii)) the set $\{x \in \mathbb{R} : \frac{d\mu}{d\lambda}(x) = \infty\}$ is μ_{ac} -zero. The following properties hold:

- (1) $\frac{d\mu}{d\lambda}(x) = \infty$ μ_s -a.e. (due to Corollary 2.5, (iv)).
- (2) $\frac{d\rho}{d\mu}(x) \in [0, \infty)$ and the limit $\frac{d\mu_l}{d\mu}(x)$ exists μ -a.e. (due to Corollary 2.5, (i), and Theorem 2.1).
- (3) $\frac{|\operatorname{Re} m(x+i\varepsilon)|}{\operatorname{Im} m(x+i\varepsilon)} \not\rightarrow \infty$ as $\varepsilon \downarrow 0$ μ -a.e. (due to Theorem 2.14).
- (4) $\lim_{\varepsilon \downarrow 0} \frac{m_l(x+i\varepsilon)}{m(x+i\varepsilon)} = \frac{d\mu_l}{d\mu}(x)$ μ_s -a.e. (due to Theorem 2.15, (ii)).

Hence, the set X_{reg} is μ_s -full. Thus we can choose a Borel set $X \subseteq X_{reg}$ with $\mu_s(\mathbb{R} \setminus X) = 0$ and $\mu_{ac}(X) = 0$ (since X_{reg} is Lebesgue-zero).

Let $\mathcal{E} := \bigcup_{l=1}^n \mathcal{E}_{\mu_l, \mu}$ where $\mathcal{E}_{\mu_l, \mu}$ are exceptional sets as in Theorem 2.1. By passing from X to $X \setminus \mathcal{E}$, we may assume that the functions $\frac{d\mu_l}{d\mu}|_X$ are Borel measurable. Set

$$X_{s,ac} := X \cap X_{reg}^{>1}.$$

It is a Borel set. The set $X_{>1} = r^{-1}(\{2, \dots, n\})$ in the formulation of the theorem is determined only up to a μ -zero (and hence $\rho_{s,ac}$ -zero) set. Since the symmetric derivative $\frac{d\mu_l}{d\mu}$ coincides μ -a.e. with the Radon-Nikodym derivative $D_{\mu} \mu_l$ (i.e., is a representative of the class of equivalent functions), we may use

$$X_{>1} := \left\{ x \in \mathbb{R} \setminus \mathcal{E} : \text{for at least two indices } \frac{d\mu_l}{d\mu}(x) > 0 \right\}.$$

Then $X_{>1} \cap X = X_{s,ac}$.

Now we can determine the Lebesgue decomposition of ρ_s with respect to μ . Choose a Borel set Y with $\rho_s(Y^c) = 0$ and $\lambda(Y) = 0$, so that $\rho_s = \mathbb{1}_Y \cdot \rho$. Our candidates for the Lebesgue decomposition are:

$$\rho_{s,ac} := \mathbb{1}_{X_{s,ac}} \cdot \rho_s, \quad \rho_{s,s} := \mathbb{1}_{X^c} \cdot \rho_s = \mathbb{1}_{X^c \cap Y} \cdot \rho.$$

Since $X \setminus X_{s,ac} \subseteq X_{reg}^1$, by Proposition 4.8, (i), we have $\rho(X \setminus X_{s,ac}) = 0$, so $\rho_s = \rho_{s,ac} + \rho_{s,s}$. Furthermore, for every $x \in X_{s,ac}$ we have $\frac{d\rho}{d\mu}(x) \in (0, \infty)$, and hence by Corollary 2.4, (iii),

$$\rho_{s,ac} \sim \mathbb{1}_{X_{s,ac}} \cdot \mu = \mathbb{1}_{X_{>1}} \cdot \mathbb{1}_X \cdot \mu = \mathbb{1}_{X_{>1}} \cdot \mu_s. \quad (4.18)$$

Finally, from $\mu_s(X^c) = 0$ it follows that $\rho_{s,s} \perp \mu_s$ and therefore $\rho_{s,s} \perp \mu$. Thus, indeed, $\rho_{s,ac}$ and $\rho_{s,s}$ are the absolutely continuous and singular part of ρ_s with respect to μ .

From the fact that $E \sim \rho$ and (4.18) it follows that Item (I) of Theorem 4.1 holds. Since $X \subseteq X_{reg}$ and $\rho_{s,ac}(X^c) = 0$, Proposition 4.8 implies that

$$\operatorname{rank} \omega(x) = r(x) - 1 \text{ for } \rho_{s,ac} - \text{a.a. points } x \in \mathbb{R},$$

which gives item (II). Due to the fact that $\mu(X^c \cap Y) = 0$, Proposition 4.3 implies

$$\operatorname{rank} \omega(x) = 1 \text{ for } \rho_{s,s} - \text{a.a. points } x \in \mathbb{R},$$

and this is item (III). \square

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Appendix A. Some Examples

In this appendix we provide four examples in order to show that all possibilities for the spectrum which are admitted by Theorem 4.1 indeed may occur. We realize these examples on the level of Schrödinger operators. Due to the general inverse theorem stated as the second part of Theorem 3.4, it would be somewhat simpler to realize them on the level of boundary relations. However, in order to remain in a more intuitive setting, we decided to stick to the Schrödinger case. Also we should say it very clearly that our emphasis in this appendix is on examples and methods rather than on maximal generality.

Let us first recall in some detail how a half-line Schrödinger operator can be considered as a boundary relation. This is of course a (if not “the”) standard example for boundary relations, see [13] and the references therein. Formulated in our present language it reads as follows.

Remark A.1. Let q be a real and locally integrable potential defined on $(0, \infty)$, and assume that 0 is a regular endpoint and that Weyl’s limit point case prevails at the endpoint ∞ . Denote by T_{max} the maximal differential operator generated in $L^2(0, \infty)$ by the differential expression $-\frac{d^2}{dx^2} + q$. For $\alpha \in \mathbb{R}$, denote by $\Gamma_{(\alpha)}$ the relation

$$\Gamma_{(\alpha)} := \left\{ \left((u; T_{max}u); (u(0) \cos \alpha + u'(0) \sin \alpha; -u(0) \sin \alpha + u'(0) \cos \alpha) \right) : u \in \text{dom } T_{max} \right\} \subseteq L^2(0, \infty)^2 \times \mathbb{C}^2$$

Then it is easy to see that $\Gamma_{(\alpha)}$ is a boundary relation (in fact, “boundary function”) for the operator T_{max} : For $\alpha = 0$ see [7, Example 1.3]. For other values of α note that $\Gamma_{(\alpha)}$ and $\Gamma_{(0)}$ are related by

$$\Gamma_{(\alpha)} = w_{\alpha} \circ \Gamma_{(0)}$$

with the $J_{\mathbb{C}}$ -unitary matrix

$$w_{\alpha} := \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix},$$

and apply Theorem 3.12.

Apparently, the selfadjoint operator

$$A_{\alpha} := \ker [\pi_1 \circ \Gamma_{(\alpha)}]$$

is nothing but the selfadjoint restriction of T_{max} given by the boundary condition

$$u(0) \cdot \cos \alpha + u'(0) \cdot \sin \alpha = 0.$$

First we consider $E_{s,ac}$ and E_{ac} and give an example that an arbitrary number of overlaps (embedded into absolutely continuous spectrum or not) can be produced. This is very simple; we elaborate it only for the sake of illustration and completeness.

Example A.2. Let λ_1, λ_2 be measures such that (λ is the Lebesgue measure)

$$\lambda_1 \perp \lambda_2, \quad \lambda_j \perp \lambda, \quad \text{supp } \lambda_j = [0, 1], \quad \lambda_j(\{0\}) = \lambda_j(\{1\}) = 0, \quad j = 1, 2.$$

For $n, m \in \mathbb{Z}, n < m$, set

$$\lambda_j^{(n,m)}(\Delta) := \sum_{l=n}^{m-1} \lambda_j(\Delta - l), \quad \Delta \text{ Borel set.}$$

Moreover, let f be the function

$$f(x) := \begin{cases} \frac{2}{3\pi} x^{\frac{3}{2}}, & x \geq \frac{9}{2}, \\ 0, & \text{otherwise.} \end{cases}$$

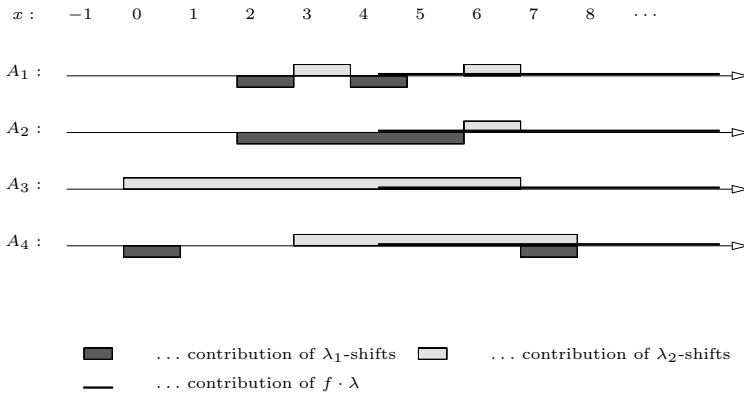
We consider the measures μ_1, \dots, μ_4 defined as

$$\begin{aligned} \mu_1 &:= \lambda_1^{(2,3)} + \lambda_1^{(4,5)} + \lambda_2^{(3,4)} + \lambda_2^{(6,7)} + f \cdot \lambda, \\ \mu_2 &:= \lambda_1^{(2,6)} + \lambda_2^{(6,7)} + f \cdot \lambda, \\ \mu_3 &:= + \lambda_2^{(0,7)} + f \cdot \lambda, \\ \mu_4 &:= \lambda_1^{(0,1)} + \lambda_1^{(7,8)} + \lambda_2^{(3,8)} + f \cdot \lambda. \end{aligned}$$

Now we appeal to the version of the Gelfand-Levitan theorem which applies to the Dirichlet boundary condition, cf. [25, §2.9]. The hypotheses of this result are obviously fulfilled, and we obtain a potentials q_1, \dots, q_4 on the half-line, such that μ_j is the measure in the integral representation of the Titchmarsh-Weyl coefficient constructed from the potential q_j with Dirichlet boundary conditions.

Let A_1, \dots, A_4 be the corresponding non-interacting operators, and let A be their pasting with standard interface conditions. The support sets of the spectral measures of A_1, \dots, A_4 can be pictured as follows:

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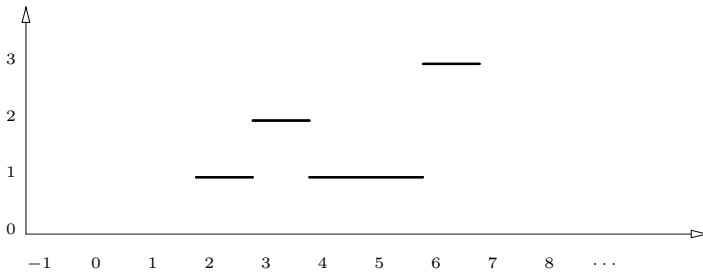


Support sets of $E_{s,ac}$ and E_{ac} including multiplicities are:

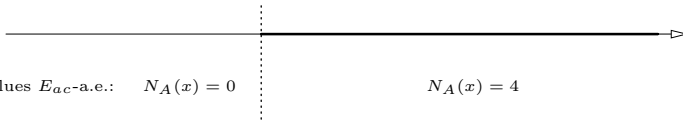
Support set of $E_{s,ac}$:



$N_A(x)$: values $E_{s,ac}$ -a.e.



Support set of E_{ac} :



This example also demonstrates that from the singular spectra $\sigma_s(A_j)$ (which are the closures of the above pictured support sets), one cannot draw any conclusions about the singular spectrum $\sigma_s(A)$ or the spectral multiplicity function N_A .

Next, we investigate with $E_{s,s}$, i.e., the possible appearance of new singular spectrum. This is not so straightforward. In order to make explicit computations, we consider the situation which resembles a single half-line operator.

Remark A.3. Let q be real and locally integrable potential defined on $(0, \infty)$, and assume that 0 is a regular endpoint and that Weyl's limit point case

prevails at the endpoint ∞ . Let $n \geq 2$, and extend the potential q to the star-graph with n edges by symmetry (i.e., consider the same potential on all edges). Matching the notation of Theorem 4.1, we thus have

$$\Gamma_l := \Gamma_{(0)}, \quad l = 1, \dots, n,$$

where $\Gamma_{(0)}$ is defined for q as in Remark A.1. Then the Weyl functions m_l are all equal, namely equal to the classical Weyl function $m_{(0)}$ constructed from the potential q . Thus, $m = n \cdot m_{(0)}$.

From Proposition 3.18 we see that the $n \times n$ -matrix valued Weyl function M corresponding to the boundary relation Γ constructed by pasting $\Gamma_1, \dots, \Gamma_n$ with standard interface conditions is given as

$$M = \frac{1}{n} \begin{pmatrix} (n-1)m_{(0)} & -m_{(0)} & \cdots & -m_{(0)} & -1 \\ -m_{(0)} & (n-1)m_{(0)} & \cdots & -m_{(0)} & -1 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ -m_{(0)} & -m_{(0)} & \cdots & (n-1)m_{(0)} & -1 \\ -1 & -1 & \cdots & -1 & -\frac{1}{m_{(0)}} \end{pmatrix}.$$

In particular,

$$\text{tr } M = \frac{1}{n} \left[(n-1)^2 m_{(0)} - \frac{1}{m_{(0)}} \right].$$

Denote by μ_0 and σ_0 the measures in the integral representations of $m_{(0)}$ and $-\frac{1}{m_{(0)}}$, respectively. Then, for the trace measure ρ constructed from M we have

$$\rho \sim \mu_0 + \sigma_0.$$

Using the notation of Theorem 4.1, we thus have

$$\mu = n \cdot \mu_{(0)}, \quad r(x) = n \text{ for } \mu\text{-a.a. points } x \in \mathbb{R},$$

and hence (denoting by $\mu_{(0),s}$ and $\sigma_{0,s}$ the singular parts of $\mu_{(0)}$ and σ_0 with respect to the Lebesgue measure)

- (I) $E_{s,ac} \sim \mu_{(0),s}$,
- (II) $N_A(x) = n - 1$ for $E_{s,ac}$ -a.a. points $x \in \mathbb{R}$,

$$E_{s,s} \sim \sigma_{0,s}. \tag{A.1}$$

Using appropriately chosen potentials q , we can now provide examples that new singular spectrum (embedded or not) does appear, or that no new singular spectrum appears.

Example A.4 (Appearance of new spectrum, partially embedded). In [30, Theorem 3.5] a class of potentials is given, such that for every boundary condition the corresponding Schrödinger operator A_α satisfies (F denotes a certain Smith-Volterra-Cantor-type set with positive Lebesgue measure)

$$\sigma_{sc}(A_\alpha) = [0, \infty), \quad \sigma_{ac}(A_\alpha) = F^2, \quad \sigma_p(A_\alpha) \cap (0, \infty) = \emptyset.$$

Using measure theoretic terms, we may thus say that there exist minimal supports of the corresponding singular continuous parts $\mu_{(\alpha),sc}$ of the spectral

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measures $\mu_{(\alpha)}$ which are Lebesgue-zero sets, mutually disjoint, and dense in $[0, \infty)$.

For such potentials we see from (A.1) that $E_{s,s} \neq 0$, i.e., new singular spectrum appears (precisely on a minimal support of $\mu_{(\frac{\pi}{2}),s}$). The part of this new spectrum located on the positive half-line is singular continuous. Moreover, since $\emptyset \neq \text{supp } E_{ac} = F^2 \subseteq [0, \infty)$, some part of it is embedded into the absolutely continuous spectrum. The spectrum originating from overlaps (which happens precisely on a minimal support of $\mu_{(0)}$) shares these properties.

Example A.5 (Appearance of new spectrum, not embedded). Consider the potential

$$q(x) := \begin{cases} k, & |x - \exp(2k\frac{3}{2})| < \frac{1}{2}, \\ 0, & \text{otherwise.} \end{cases}$$

This potential was studied in [34], and it turned out that for every boundary condition the corresponding selfadjoint operator A_α satisfies

$$\sigma_{ac}(A_\alpha) = \emptyset, \quad \sigma_{sc}(A_\alpha) = [0, \infty).$$

Moreover, depending on the boundary condition either $\sigma_p(A_\alpha)$ is empty or consists of one negative eigenvalue. Expressed in measure theoretic terms, this means that there exist minimal supports of the corresponding spectral measures $\mu_{(\alpha)}$ which are all Lebesgue-zero sets, are mutually disjoint, whose intersection with $[0, \infty)$ is dense in $[0, \infty)$, which contain at most one point on the negative half-line, and that $\mu_{(\alpha)}$ has no point masses in $[0, \infty)$.

For this potential we see from (A.1) that $E_{s,s} \neq 0$, i.e., new singular spectrum appears. This new part of the spectrum is singular continuous and embedded into the spectrum originating from overlaps (the closures of minimal supports of $E_{s,ac}$ and $E_{s,s}$ are both equal to $[0, \infty)$).

Example A.6 (Non-appearance of new spectrum). We follow the idea given in [11] to construct examples, and use the type of measures discussed in [11, Example 1] and the Gelfand-Levitan theorem. However, we need to refer to the version of the Gelfand-Levitan theorem which applies to the Dirichlet boundary condition, cf. [25, §2.9].

Let us first recall the argument made in [11, Example 1].⁹ Let f be a continuous, bounded, and positive function on \mathbb{R} , let λ_0 be a positive measure with $\int_{\mathbb{R}} \frac{d\lambda_0(t)}{1+t^2} < \infty$, and consider the measure

$$\nu := \lambda_0 + f \cdot \lambda,$$

where λ is the Lebesgue measure. Denote by m the corresponding Herglotz-function

$$m(z) := \int_{\mathbb{R}} \left(\frac{1}{t-z} - \frac{t}{1+t^2} \right) d\nu(t), \quad z \in \mathbb{C} \setminus \mathbb{R}.$$

⁹Again, we do not aim for maximal generality.

Then, since $\nu \geq f \cdot \lambda$, we have

$$\liminf_{\varepsilon \downarrow 0} \operatorname{Im} m(x + i\varepsilon) \geq \pi f(x), \quad x \in \mathbb{R}.$$

In particular, the function $m(z)$ never approaches a real boundary value when z tends to a real point (along a perpendicular ray).

Consider the Herglotz functions

$$m_\tau(z) := \frac{\tau m(z) - 1}{m(z) + \tau}, \quad \tau \in \mathbb{R},$$

and let ν_τ be the measure in the integral representation of m_τ . Then, by Aronszajn–Donoghue (cf. [15, Theorem 3.2,(3.17)]), the measures ν_τ are all absolutely continuous with respect to the Lebesgue measure.

Next, we make an appropriate choice of f and λ_0 , so to allow an application of the Gelfand-Levitan theorem. Set, for example,

$$f(x) := \frac{2}{3\pi} \cdot \begin{cases} e^{x-1}, & x < 1, \\ x^{\frac{3}{2}}, & x \geq 1, \end{cases}$$

and let λ_0 be compactly supported. Then the hypotheses of the Gelfand-Levitan theorem are obviously fulfilled, and we obtain a potential q on the half-line, such that the measure ν is the measure in the integral representation of the Weyl function constructed from the potential q with Dirichlet boundary conditions.

Symmetrically extending the potential q constructed in the above paragraph, yields examples with

- (i) $E_s = 0$ (choose $\lambda_0 = 0$),
- (ii) $E_{s,ac} \neq 0$ but $E_{s,s} = 0$ (choose λ_0 to be singular).

Appendix B. Other Boundary/Interface Conditions

In the first statement of this section we show how Theorem 4.1 can be used to deduce the classical result of Aronszajn–Donoghue that the singular parts of the spectral measures corresponding to different boundary conditions in a half-line problem are mutually singular. This approach is of course more complicated than the original one, and hence should not be viewed as a “new proof of an old result”. The reasons why we still find it worth to be elaborated are: (1) reobtaining previously known results gives a hint that the new result is not unnecessarily weak, and (2) it demonstrates the usage of boundary relations with a nontrivial multivalued part.

Corollary B.1. (Aronszajn–Donoghue) *Let q be a real and locally integrable potential defined on $(0, \infty)$, assume that 0 is a regular endpoint and that Weyl’s limit point case prevails at the endpoint ∞ . Moreover, for $\alpha \in [0, \pi)$, denote by A_α the selfadjoint operator given by the differential expression $-\frac{d^2}{dx^2} + q$ and the boundary condition*

$$u(0) \cdot \cos \alpha + u'(0) \cdot \sin \alpha = 0.$$

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Then, for each two $\alpha_1, \alpha_2 \in [0, \pi), \alpha_1 \neq \alpha_2$, the singular parts of the spectral measures of A_{α_1} and A_{α_2} are mutually singular.

Proof. Let T_{max} be the maximal operator associated with differential expression $-\frac{d^2}{dx^2} + q$, and let Γ_{α_j} be the boundary relations constructed in Remark A.1. Moreover, denote by $\mu_{(\alpha_j)}, j = 1, 2$, the measure in the integral representation of the Weyl function of Γ_{α_j} .

Set $\beta := \alpha_1 - \alpha_2$, and denote by $\mathring{\Gamma}$ the boundary relation

$$\mathring{\Gamma} := \{((0; 0); (-w \sin \beta; w \cos \beta)) : w \in \mathbb{C}\} \subseteq \{0\}^2 \times \mathbb{C}^2. \quad (\text{B.1})$$

Then the pasting Γ of $\mathring{\Gamma}$ and $\Gamma_{(\alpha_1)}$ with standard interface conditions is given as

$$\begin{aligned} \Gamma = & \left\{ \left(\left(\left(\begin{pmatrix} 0 \\ u \end{pmatrix} \right); \begin{pmatrix} 0 \\ T_{max} u \end{pmatrix} \right) \right); \right. \\ & \left. \left(\begin{pmatrix} w \sin \beta + [u(0) \cos \alpha_1 + u'(0) \sin \alpha_1] \\ w \cos \beta + [-u(0) \sin \alpha_1 + u'(0) \cos \alpha_1] \end{pmatrix} \right); \begin{pmatrix} -w \cos \beta \\ -[u(0) \cos \alpha_1 + u'(0) \sin \alpha_1] \end{pmatrix} \right) \right\}; \\ & w \in \mathbb{C}, u \in \text{dom } T_{max} \left\} \subseteq (\{0\} \times L^2(0, \infty))^2 \times \mathbb{C}^2. \end{aligned}$$

A short computation shows that (we identify $\{0\} \times L^2(0, \infty)$ with $L^2(0, \infty)$) $\ker[\pi_1 \circ \Gamma] = A_{\alpha_2}$. The Weyl function \mathring{m} of $\mathring{\Gamma}$ is equal to the real constant $-\cot \beta$; note here that $\beta \in (-\pi, \pi) \setminus \{0\}$. The measure $\mathring{\mu}$ in its integral representation is thus equal to 0. Using the notation of Theorem 4.1, we have

$$\mu = \mu_{(\alpha_1)}, \quad r(x) = 1 \text{ for } \mu\text{-a.a. points } x \in \mathbb{R},$$

and hence $E_{s,ac} = 0$. However, $E \sim \mu_{(\alpha_2)}$, and we see that $\mu_{(\alpha_2),s} = \mu_{(\alpha_2),s,s} \perp \mu_{(\alpha_1)}$. Therefore $\mu_{(\alpha_1),s} \perp \mu_{(\alpha_2),s}$. \square

Finally, we provide some knowledge on other interface conditions than the standard ones. In the context of this example, it is however important to add two remarks:

- (1) The formula for spectral multiplicity will not be given in terms of the spectral measures of the non-interacting operators, but in terms of the corresponding Weyl functions. Hence, the below result cannot be viewed as a strict analogue of Theorem 4.1 for other interface conditions.
- (2) The applied method provides knowledge only about a particular (small) class of interface conditions. It does not lead to a treatment of arbitrary interface conditions on a star-graph, and even less to a formula for multiplicity on graphs with a more complicated geometry.

Proposition B.2. *Let $n \geq 2$ and real valued locally integrable potentials $q_l, l = 1, \dots, n$, on the half-line be given, such that q_1 is regular at 0 and in Weyl's limit point case at ∞ . Let $a_1, \dots, a_n \in [0, 2\pi)$ and $b \in (0, \pi)$ be given, and consider the selfadjoint matrix Schrödinger operator A defined on $L^2(0, \infty)^n$*

by the differential expression $-\frac{d^2}{dx^2} + V$ with the diagonal matrix potential

$$V(x) := \begin{pmatrix} q_1(x) & & \\ & \ddots & \\ & & q_n(x) \end{pmatrix}$$

and the interface conditions

$$\begin{aligned} u_1(0) \cos a_1 + u_1'(0) \sin a_1 = \dots = u_n(0) \cos a_n + u_n'(0) \sin a_n, \\ \sum_{l=1}^n [u_l(0) \cos(a_l - b) + u_l'(0) \sin(a_l - b)] = 0. \end{aligned} \tag{B.2}$$

Denote the Titchmarsh-Weyl coefficient constructed from the potential q_l (with Dirichlet boundary conditions) as m_l , and set

$$S_l := \begin{cases} \{x \in \mathbb{R} : \lim_{\varepsilon \downarrow 0} m_l(x + i\varepsilon) = -\cot a_l\}, & a_l \notin \{0, \pi\}, \\ \{x \in \mathbb{R} : \lim_{\varepsilon \downarrow 0} \operatorname{Im} m_l(x + i\varepsilon) = \infty\}, & a_l \in \{0, \pi\}. \end{cases}$$

Then

$$N_A(x) = \begin{cases} \#\{l \in \{1, \dots, n\} : x \in S_l\} - 1, & E_{s,ac}\text{-a.e.}, \\ 1, & E_{s,s}\text{-a.e.} \end{cases}$$

Proof. First of all, let us explicitly state how the operator A acts:

$$A \begin{pmatrix} u_1 \\ \vdots \\ u_n \end{pmatrix} := -\frac{d^2}{dx^2} \begin{pmatrix} u_1 \\ \vdots \\ u_n \end{pmatrix} + \begin{pmatrix} q_1 u_1 \\ \vdots \\ q_n u_n \end{pmatrix},$$

on the domain

$$\begin{aligned} \operatorname{dom} A := \left\{ (u_1, \dots, u_n) \in \prod_{l=1}^n L_2(0, \infty) : \right. \\ \left. \begin{aligned} &u_l, u_l' \text{ are absolutely continuous, } -u_l'' + q_l u_l \in L_2(0, \infty), \\ &u_1, \dots, u_n \text{ satisfy the interface conditions (B.2)} \end{aligned} \right\}. \end{aligned}$$

For $l = 1, \dots, n$ let Γ_l be the boundary relation which is defined from the potential q_l as $\Gamma_{(a_l)}$, cf. Remark A.1. Moreover, let $\mathring{\Gamma}$ be the boundary relation defined in (B.1). We consider the pasting of $\Gamma_1, \dots, \Gamma_n, \mathring{\Gamma}$ with standard interface conditions. Then an element (u_1, \dots, u_n) belongs to the domain of the operator¹⁰ $\ker[\pi_1 \circ \Gamma]$ if and only if

$$\exists w \in \mathbb{C} : \begin{cases} u_l(0) \cos a_l + u_l'(0) \sin a_l = -w \sin b, & l = 1, \dots, n, \\ \sum_{l=1}^n [-u_l(0) \sin a_l + u_l'(0) \cos a_l] = -w \cos b. \end{cases}$$

Eliminating w from these equations, yields the assertion. \square

Remark B.3. Some observations are in order:

¹⁰Similar to the previous example: we identify $L^2(0, \infty)^n \times \{0\}$ with $L^2(0, \infty)^n$.

- (i) It is interesting to notice that the support set and the multiplicity function corresponding to $E_{s,ac}$ do not depend on the choice of the parameter b .
- (ii) The fact that the value “ $b = 0$ ” is excluded is natural. For this value the conditions (B.2) reduce to

$$u_1(0) \cos a_1 + u_1'(0) \sin a_1 = \cdots = u_n(0) \cos a_n + u_n'(0) \sin a_n = 0,$$

i.e., the operator A is equal to the direct sum of the non-interacting operators A_l defined by the potentials q_l using the boundary condition $u_l(0) \cos a_l + u_l'(0) \sin a_l = 0$.

- (iii) The case that “ $b = \frac{\pi}{2}$ ” could be treated somewhat simpler. For this value the operator A coincides with $\ker[\pi_1 \circ \Gamma]$ where Γ is the pasting of $\Gamma_1, \dots, \Gamma_n$ with standard interface conditions.

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