in collaboration with Alexander Haber and Andreas Schmitt

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### Motivation

Cold dense matter in strong magnetic fields is found in ....

magnetars

on the surface  $B \sim 10^{15} \,\mathrm{G}$ R.C. Duncan and C. Thompson, Astrophys.J. 392, L9 (1992) in the core up to  $B \sim 10^{18} \,\mathrm{G}$ D. Lai and S. Shapiro, Astrophys.J. 383, 745 (1991)

• binary neutron star mergers Gravitational wave signature sensitive to equation of state J.S. Read et al., Phys.Rev. D88, 044042 (2013) The magneto-rotational instability may increase the initial B-field by one order of magnitude D.M. Siegel et al., Phys.Rev. D87, 121302 (2013)





### Extended Linear Sigma Model (eLSM) part I

#### Mesons:

D. Paraganlija et al., Phys.Rev. D82, 054024 (2010) & Phys.Rev. D87, 014011 (2013) & Phys.Rev. D87, 014011 (2013)

The eLSM is based on

- degrees of freedom of the QCD vacuum: hadrons
- symmetries of QCD: chiral and dilatation invariance

Relevant ingredients  $(N_f = 2)$ :

- chiral partner of pions  $\sigma$ : a finite VEV breaks chiral symmetry fluctuations correspond to the state  $f_0(1370)$
- the lightest scalar resonance  $f_0(500)$  corresponds to an  $SU(2)_L \times SU(2)_R$ singlet state  $\chi$ tetraquark or pion-pion resonance
- the  $SU(2)_L \times SU(2)_R$  singlet vector meson  $\omega_\mu$  corresponding to  $\omega(782)$

### eLSM part I

#### Tree level free energy:

$$\Omega_{\rm mes}^{\rm tree} = -\frac{1}{2}m\sigma^2 + \frac{\lambda}{4}\sigma^4 - \epsilon\sigma - \frac{1}{2}m_\omega^2\omega_0^2 + \frac{1}{2}m_\chi^2\chi^2 - g\chi\sigma^2,$$

with gap equation for  $\chi$  and  $\omega_0$ 



#### Including nucleons in the mirror assignment:

- C. DeTar and T. Kunihiro, Phys.Rev. D39, 2805 (1989)
- S. Gallas, F. Giacosa, and D.Rischke, Phys.Rev. D82, 014004 (2010)

Isospin-douplet spinor  $\psi_1$  and its parity partner  $\psi_2$  obeying

$$\psi_{1,R/L} \to U_{R/L}\psi_{1,R/L}, \quad \psi_{2,R/L} \to U_{L/R}\psi_{2,R/L},$$

allow for a chirally symmetric mass term

$$m_0(\overline{\psi}_2\gamma_5\psi_1-\overline{\psi}_1\gamma_5\psi_2).$$

dilatation invariance  $\to m_0$  generated dynamically, e.g.  $m_0 \propto \chi$ . Interaction Lagrangian:

$$\mathcal{L}^{\mathrm{int}} = -\sum_{i} \overline{\psi}_{i} (\frac{g_{i}\sigma}{2} + g_{\omega}\gamma^{0}\omega_{0})\psi_{i} - a\chi(\overline{\psi}_{2}\gamma_{5}\psi_{1} - \overline{\psi}_{1}\gamma_{5}\psi_{2})$$

Effective fermion mass and chemical potential:

$$\begin{aligned} \mathcal{M}_{\mathcal{N},\mathcal{N}^*} &= \pm \frac{g_1 - g_2}{2} \sigma + \sqrt{\left(\frac{g_1 + g_2}{4} \sigma\right)^2 + (a\chi)^2} \\ \mu^* &= \mu - g_\omega \omega_0 \end{aligned}$$

Free Energy in "no sea approximation" N.K. Glendenning, Phys.Lett. B208, 335 (1988) & Nucl.Phys. A493, 521 (1989)

$$\Omega = \Omega_{\text{mes}}^{\text{tree}} + \sum_{i} \int \frac{d^{3}k}{2\pi^{3}} \left\{ -E_{i}(k) + [E_{i}(k) - \mu^{*}] \Theta[\mu^{*} - E_{i}(k)] \right\}$$
$$E_{N,N^{*}}(k) = \sqrt{\vec{p}^{2} + M_{N,N^{*}}^{2}}$$

Solving the gap equation - gas/liquid phase transition:



- $M_N^{\rm vac} = 939.12 \,\mathrm{MeV}$ 
  - $\mu_{c} = 923.06 \,\mathrm{MeV}$
- $\Rightarrow E_B = 16.06 \,\mathrm{MeV}$

$$\rho_0 = 0.15 \, \mathrm{fm}^{-3}$$

 $M_{N^*}^{\text{vac}} = 1535.55 \text{ MeV}$  is identified with N(1535)

#### Fermions in a magnetic field

the dispersion relations for fermions with charge q in a magnetic field are altered: Landau level quantization

$$E_{\ell}(k_z) = \sqrt{k_z^2 + M^2 + 2|q|B\ell}$$

Free energy for one charged fermion species  $\psi$  at  $\mu={\cal T}={\rm 0}$ 

$$\begin{split} \Omega_{\rm f} &= \frac{B^2}{2} - \frac{|q|B}{4\pi^2} \sum_{\ell} \left(2 - \delta_{\ell,0}\right) \int {\rm d}\, k_z \, E_\ell(k_z) = \\ &= -2 \int \frac{{\rm d}^3 k}{(2\pi)^3} \sqrt{\vec{k}^2 + M^2} + \frac{B_r^2}{2} - \frac{\left(|q_r|B_r\right)^2}{24\pi^2} \ln \frac{2 \, |q_r|B_r}{\Lambda^2 A^{12}} \\ &- \frac{\left(|q_r|B_r\right)^2}{2\pi^2} \left[\frac{x^2}{4} \left(3 - 2\ln x\right) + \frac{x}{2} \left(\ln \frac{x}{2\pi} - 1\right) + \psi^{(-2)}(x)\right], \end{split}$$

with  $x = M^2/(2|q_r|B_r)$ , A being the Glaisher constant and A is the renormalization scale

#### Magnetic catalysis in the eLSM

This contribution gives the well known magnetic catalysis effect V.P. Gusynin, V.A. Miransky, and I.A. Shovkovy, Phys.Lett., B349:477483 (1995)





Phase diagram and mass gap:



Phase diagram and mass gap:









### So far the effect of the B-dependent vacuum terms has always been neglected:

A. Broderick, M.Prakash, and J.M. Lattimer, Astrophys. J. 537, 351 (2000) & Phys.Lett. B531, 167 (2002),
M. Sinha, B. Mukhopadhyay, and A. Sedrakian, Nucl.Phys. A898, 43 (2013),
A. Rabhi, P. Panda, and C. Providencia, Phys.Rev. C84, 035803 (2011),
FP, A. Rebhan, and A. Schmitt, J.Phys. G39, 054006 (2012),
J. Dong, W. Zuo, and J. Gu, Phys.Rev. D87, 103010 (2013),
R.C.R. de Lima, S.Avancini, and C. Providencia, Phys.Rev. C89, 035804 (2013),
R. Casali, L.Castro, and D. Menezes, Phys.Rev. C89, 015805 (2014),
incorporation of magnetic catalysis in models of nuclear matter is the

first important step

#### For realistic applications we need to:

- include the anomalous magnetic moment
- allow for anisotropic condensates: (magnetic) chiral spiral

   A. Rebhan, S. Stricker, and A. Schmitt, JHEP 0905, 084 (2009),
   I. Frolov, V.C. Zhukovsky, and K. Klimenko, Phys.Rev. D82, 076002 (2010),
   A. Rebhan, S. Stricker, and A. Schmitt, JHEP 0905, 084 (2009),
   FP, A. Rebhan, and A. Schmitt, JHEP 1103, 033 (2011),
   FP, A. Rebhan, and A. Schmitt, J.Phys. G39, 054006 (2012)
- $\bullet$  implose neutrality and  $\beta\text{-equilibrium}$  conditions
- analyze the chiral phase transition in the eLSM