

# Silicene Passivation by Few-Layer Graphene

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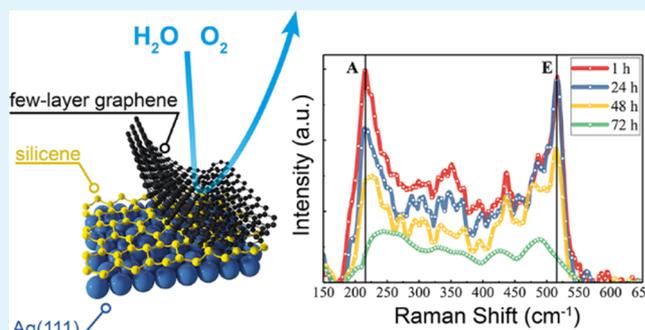
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## Supporting Information

**ABSTRACT:** The stabilization of silicene at ambient conditions is essential for its characterization, future processing, and device integration. Here, we demonstrate in situ encapsulation of silicene on Ag(111) by exfoliated few-layer graphene (FLG) flakes, allowing subsequent Raman analysis under ambient conditions. Raman spectroscopy measurements proved that FLG capping serves as an effective passivation, preventing degradation of silicene for up to 48 h. The acquired data are consistent with former in situ Raman measurements, showing two characteristic peaks, located at 216 and 515  $\text{cm}^{-1}$ . Polarization-dependent measurements allowed to identify the two modes as A and E, demonstrating that the symmetry properties of silicene are unaltered by the capping process.

**KEYWORDS:** silicene, graphene, passivation, Raman spectroscopy, 2D materials



## 1. INTRODUCTION

The outstanding electrical properties of silicene are of foremost interest for the development of high-performance two-dimensional (2D) material-based devices.<sup>1</sup> Silicene combines an ultrahigh carrier mobility<sup>2–4</sup> with the unique opportunity to tune the band gap by applying an external electric field<sup>5,6</sup> or by chemical functionalization.<sup>7,8</sup>

Because of the instability of silicene, its synthesis requires ultrahigh vacuum (UHV) conditions and a supporting substrate.<sup>3</sup> Silver still remains the most popular choice for growing silicene epitaxially over large areas because of the moderate interaction as well as the low lattice mismatch between the Ag(111) surface and silicene.<sup>9,10</sup> However, silicene grown on Ag(111) undergoes oxidation when exposed to ambient conditions,<sup>1,11</sup> breaking the hexagonal structure consisting of both  $\text{sp}^2$ - and  $\text{sp}^3$ -hybridized orbitals.<sup>12,13</sup> Therefore, exploring a reliable process for the encapsulation of silicene is fundamental.

An ideal passivation layer has to stabilize silicene as well as enable the analysis of its structural integrity after subsequent processing steps. The properties of the passivated layer can be investigated by Raman spectroscopy, which has been proven to be an extremely powerful tool for the identification and characterization of 2D materials.<sup>14–16</sup> At present, a reliable knowledge of the vibrational properties of silicene on Ag(111) has been gained through a growing number of in situ Raman experiments,<sup>15,17,18</sup> thus representing a solid model for the identification and analysis of the encapsulated silicene layer. Hence, this nondestructive technique is the method of choice

for the identification of silicene in its different phases and for verifying its structural integrity.

Air-stable 2D materials such as graphene and hexagonal boron nitride (hBN) are considered to be excellent passivation layers because of their inertness,<sup>19</sup> high flexibility,<sup>20</sup> and outstanding impermeability to even the smallest molecules.<sup>21</sup> Moreover, they are optically transparent and non-Raman active in the fingerprint region of silicene and are thus ideal candidates for its stabilization.

Here, we present an encapsulation method based on the in situ mechanical exfoliation of few-layer graphene (FLG) flakes directly on top of silicene epitaxially grown on Ag(111).

The FLG is capable of stabilizing silicene up to 48 h and allows acquiring ex situ  $\mu$ -Raman spectra that are in good agreement with in situ experiments and provide high lateral resolution.

## 2. EXPERIMENTAL SECTION

Experiments were performed in an UHV system at a base pressure of  $5 \times 10^{-11}$  mbar. The Ag(111)/mica growth substrate (MaTecK) was cleaned through several cycles of  $\text{Ar}^+$  ion sputtering (1 keV, 5 min), followed by annealing at 800 K for 20 min. Silicene synthesis was achieved by silicon electron beam evaporation (Specs, EBE-1) with a deposition rate of  $\approx 0.02$  ML/min and a substrate temperature of 535 K. The annealing and growth temperatures were controlled by an

Received: November 26, 2018

Accepted: March 13, 2019

Published: March 13, 2019

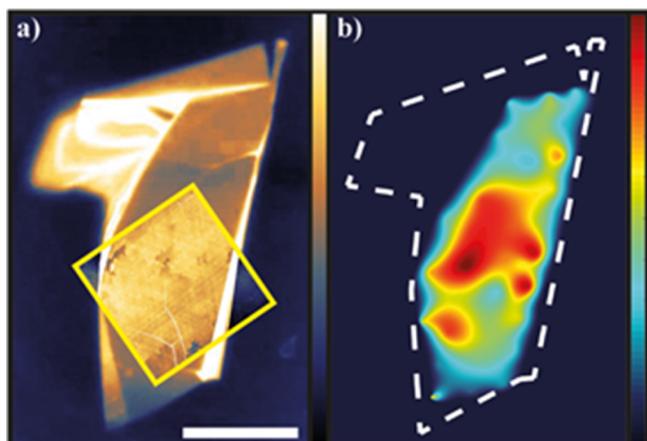
infrared pyrometer with a precision of  $\pm 2$  K. The structure and quality of the obtained silicene were verified via LEED measurements.

To passivate silicene, thin graphitic films were prepared by mechanical exfoliation<sup>22</sup> from natural graphite (HQGraphene) on top of a vacuum compatible polyimide tape (thickness of  $\approx 25 \pm 2$   $\mu\text{m}$ ) with a siliconic adhesive (HB830). The tape was inserted into a specifically designed stamping device located in a dedicated UHV chamber, directly connected to the growth chamber via a transfer line. The encapsulation process was performed by stamping the previously prepared polyimide tape with graphite flakes onto silicene at a substrate temperature of  $\approx 373$  K. The particular stamping device allows a precise control of the encapsulation procedure in order to prevent mechanically induced damages to the sample.

Ex situ Raman analysis was performed in a back-scattering geometry using a confocal  $\mu$ -Raman setup (Alpha300, WITec) with a 532 nm (2.33 eV) Nd:YAG laser and a spot size of  $\approx 720$  nm. All measurements were performed at ambient conditions, and the laser power was chosen to minimize heating effects. For further surface topography analysis, atomic force microscopy (AFM) measurements were performed in tapping mode.

### 3. RESULTS

**3.1. FLG Capping of Silicene.** The Raman measurements revealed large areas of stable passivated silicene under various FLG flakes. However, not all of them were capable of passivating silicene. A large majority of the FLG flakes that encapsulated silicene were found to be between 10 and 15 nm thick, verified by AFM as presented in Figure 1. This does not



**Figure 1.** (a) Optical microscopy image of a stamped  $10 \times 40$   $\mu\text{m}^2$  FLG flake on top of the silicene/Ag(111) surface. The yellow framed area shows a detailed AFM measurement of the capping flake (Lateral scale bar 10  $\mu\text{m}$ , logarithmic  $z$  scale 15 nm). (b) Normalized intensity heat map of the out-of-plane vibrational silicene Raman mode, interpolated from several dozen Raman point scans. The dashed white line indicates the border of the FLG layer.

necessarily imply that thinner FLG flakes are less suited for the passivation. It rather displays the well-known problem of locating thin 2D materials,<sup>23,24</sup> especially on metallic substrates,<sup>25</sup> where graphene up to a certain amount of layers is essentially invisible. Thicker FLG flakes did not allow a detailed Raman analysis of the underlying substrate. Furthermore, larger flakes showed a higher probability of an intact subjacent silicene layer because typically no signal of silicene could be detected close to the flake edges. A further issue for the integrity of silicene was the inevitably introduced contaminations from the polyimide tape. Cautious treatment of the FLG flakes and exfoliation tape in cleanroom

environment, as well as sufficient bake-out and degassing time in the UHV chamber, proved to have significant influence on the yield of capped silicene.

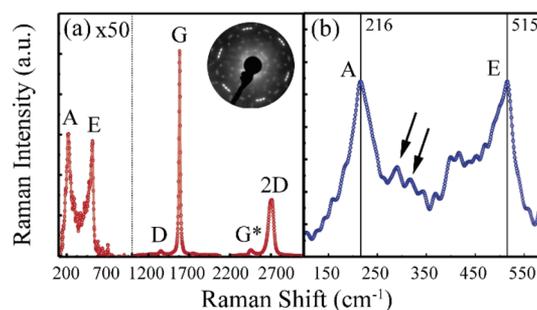
Figure 1a illustrates an optical image, including an AFM topography inset, of a  $\approx 10$  nm thick and a  $10 \times 40$   $\mu\text{m}^2$  sized FLG flake on top of the Ag(111)/silicene sample after the stamping process. A normalized intensity heat map of the most prominent silicone-related Raman signal can be seen in Figure 1b. The intensity of the out-of-plane vibrational mode was taken as the indication for the presence of silicene.

The encapsulated silicene film showed good stability for up to 48 h. The inability of permanent stabilization is a problem that also occurs when using graphene as protective capping for other materials such as black phosphorus<sup>26,27</sup> or metals.<sup>28,29</sup> Capped substrates can experience oxidation at grain boundaries and defects of the protective graphene layer.<sup>28,29</sup>

For black phosphorus encapsulated by graphene, degradation has been observed from the edges of the graphene crystals,<sup>26,27</sup> which is in good agreement with our experimental observations. Concerning silicene, a direct degradation from the edges seems unlikely because oxygen adatoms prefer to accommodate in the surface area rather than at the sides, as the dangling bonds of edge silicon atoms are stabilized by the Ag substrate.<sup>30,31</sup> Therefore, the degradation is suspected to occur because of an imperfect adhesion of graphene to the substrate, which allows water and oxygen to enter at the interface. Furthermore, there is a clear absence or marked decrease of the silicene Raman signal at macroscopic defects such as wrinkles in the FLG capping layer, indicating a preferential diffusion path for the oxidizing species.

To circumvent this problem further, encapsulating materials such as hBN have been suggested. Although thin hBN capping layers were shown to have similar issues with the long-term stability of subjacent substrates,<sup>32</sup> recent studies showed promising results for multilayered structures.<sup>33</sup> However, our passivation experiments using the same approach with hBN appeared thus far to be incapable of stabilizing silicene.

**3.2. Raman Characterization of Encapsulated Silicene.** Figure 2a shows the Raman spectrum of a FLG flake encapsulating an epitaxial  $4 \times 4$  silicene monolayer, consisting mainly of  $4 \times 4$  silicene accompanied by the ( $\sqrt{13} \times$



**Figure 2.** (a) Raman spectrum of the FLG flake encapsulating 1 ML of silicene grown on Ag(111). Alongside FLG Raman modes, two additional peaks assigned to silicene (A and E) can be clearly distinguished in the low-frequency region of the spectrum (enlarged 50 times). The small inset shows the silicene LEED pattern. (b) Enlarged view of the low-frequency region of the spectrum. Two main modes are clearly visible, located at 216 and 515  $\text{cm}^{-1}$ . Additional modes can be spotted between 250 and 350  $\text{cm}^{-1}$  (indicated by arrows), compatible with the literature reported out-of-plane optical phonon (ZO) modes.

$\sqrt{13}$ )R13.9° reconstruction, atop of the Ag(111) growth substrate. Although the  $4 \times 4$  structure was reported to be the most stable one, it coexists with other silicene reconstructions because their formation energies are similar.<sup>34,35</sup> In this work, we will refer to the mixture of  $4 \times 4$  and ( $\sqrt{13} \times \sqrt{13}$ )R13.9° silicene as silicene only, if not stated otherwise.

The spectrum has been measured at ambient conditions and room temperature, approximately 30 min after the sample was removed from the UHV chamber. The Raman signal shows the well-known peaks for exfoliated FLG, located at 1350, 1585, 2450, and 2700  $\text{cm}^{-1}$ , labeled, respectively, as D, G, G\*, and 2D, following the literature.<sup>36,37</sup>

Two additional Raman modes can be clearly distinguished at 216 and 515  $\text{cm}^{-1}$ . Previous in situ<sup>15,18</sup> and ex situ<sup>38</sup> Raman studies reliably assigned a similar vibrational signature to the silicene monolayer grown on Ag(111).

Figure 2b shows an enlarged view of the fingerprint region of silicene. Raman spectra that are acquired in situ using macro setups always show a broad peak of a-Si because of the unavoidable quantity of amorphous silicon grown on unfavorable portions of the Ag substrate. The high signal of the amorphous silicon hinders possible low-intensity Raman modes situated within the 300–500  $\text{cm}^{-1}$  spectral range. The  $\mu$ -Raman setup used in the present work, instead, is capable of analyzing sub-micrometer areas, allowing highly resolved detection of a-Si-free silicene spectra.

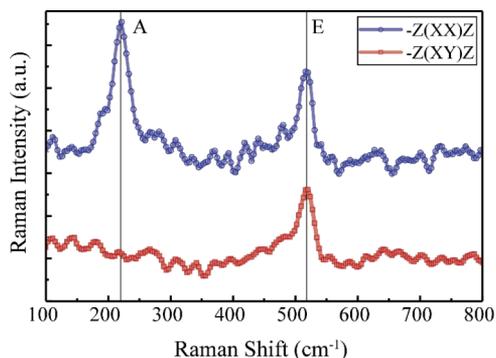
The 515  $\text{cm}^{-1}$  phonon mode is reported by both in situ and ex situ studies and has been assigned to the  $E_{2g}$  vibration mode for silicene superstructures, related to the longitudinal and transversal in-plane optical (LO and TO) modes.<sup>15,17,18,38</sup> The 216  $\text{cm}^{-1}$  feature, instead, has been assigned to mode A, related to out-of-plane optical (ZO) phonons. Interestingly, the A mode is reported only by in situ analysis.<sup>15,17,18</sup> The previous ex situ Raman study of silicene exploited a different encapsulation method, where the stabilization of the 2D layer is obtained with the deposition of a thin Al layer that oxidizes preferentially, resulting in an  $\text{Al}_2\text{O}_3$  protecting layer.<sup>38</sup> However, the encapsulating layer is responsible for the presence of a broad peak that likely hinders the detection of certain Raman features.

In addition to the two high-intensity peaks, it is also possible to distinguish two small peaks around  $\approx 295$  and  $\approx 315$   $\text{cm}^{-1}$  (indicated by arrows), which are in agreement with Raman features detected via tip enhanced Raman spectroscopy (TERS) and assigned to ZO vibrational modes.<sup>17</sup>

Polarization-dependent Raman measurements were conducted in order to confirm unambiguously the identification of silicene and to gain a deeper knowledge on the structure of the capped 2D layer. It is known that the  $4 \times 4$  silicene grown on Ag(111) shows a  $C_{6v}$  symmetry, while freestanding silicene belongs to the  $D_{3d}$  group. The minority phase ( $\sqrt{13} \times \sqrt{13}$ )R13.9°, instead, belongs to the zero symmetry group  $C_1$ . For this reason, the polarization-dependent measurements will be discussed only regarding the  $4 \times 4$  phase.<sup>34,39</sup>

The  $C_{6v}$  group possesses three possible phonon symmetries: A, E1, and E2.<sup>40</sup> The selection rules for this group predict that in a back-scattering geometry configuration, the A mode is suppressed when the spectrum is collected in crossed polarization configuration ( $-Z(XY)Z$ )—following Porto notation), whereas it is visible in parallel polarization configuration ( $-Z(XX)Z$ ).<sup>41</sup> E1 and E2 modes, instead, are always visible, regardless of the chosen polarization geometry of the Raman setup. For this reason, peaks that belong to one of these two

modes will be assigned to a more general E mode. Figure 3 illustrates the comparison between the spectrum of the capped silicene collected in crossed and parallel polarization configurations.



**Figure 3.** Raman spectrum of silicene collected in back-scattering geometry, with parallel polarization configuration [ $-Z(XX)Z$  (Porto notation)—blue] and crossed polarization configuration [ $-Z(XY)Z$  (Porto notation)—red]. The Raman mode located at 216  $\text{cm}^{-1}$  is visible only in parallel polarization configuration and can be assigned to the phonon symmetry “A”. The 515  $\text{cm}^{-1}$  peak is detected in both configurations and thus can be assigned to the “E” symmetry.

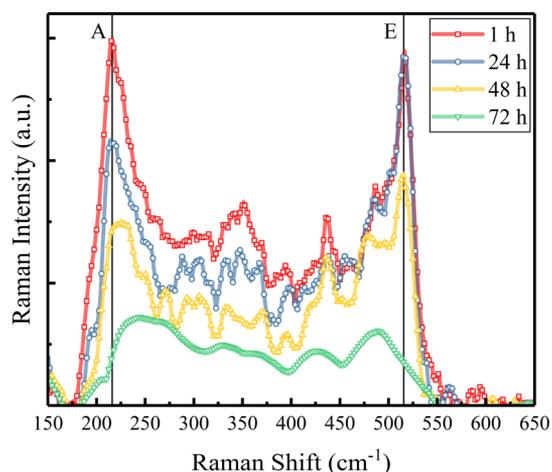
Clearly, the peak located at 216  $\text{cm}^{-1}$  vanishes when the spectrum is collected in a crossed polarization geometry, whereas the peak located at 515  $\text{cm}^{-1}$  is always visible, in perfect agreement with the results from the in situ experiment.<sup>15</sup> These observations confirm the assignment of the two peaks to phonon modes A and E. It is important to stress that the appearance of phonon mode A, which is not detected in bulk silicon structures, is a clear indication of the two-dimensionality of the capped material.<sup>42</sup>

It must be noted that the Raman signature of silicene is particularly sensitive to structural changes of the 2D layer. For example, the E mode location for freestanding silicene is simulated at 570  $\text{cm}^{-1}$  because of the shorter bond length compared to silicene on Ag(111).<sup>43</sup> Thus, the actual Raman analysis shows that silicene–FLG interaction is weak enough to preserve the particular structure of silicene. This is crucial for preserving the electrical properties of the silicene layer, as they can be altered by changes in the geometrical structure.<sup>44,45</sup>

The passivating capabilities of the FLG flake were probed by following the time evolution of the encapsulated silicene Raman spectrum while exposing the sample to ambient conditions.

Figure 4 shows the comparison between the spectra acquired 1, 24, 48, and 72 h after the removal of the sample from the UHV chamber.

The Raman spectra collected after 1 and 24 h of exposure to ambient conditions look similar, with no visible shift of the two high-intensity modes and the same behavior in the spectral region between 250 and 500  $\text{cm}^{-1}$ . In particular, the intensity of the E mode is unchanged, whereas the A mode shows a slightly lower intensity. This difference can be explained by the obvious experimental difficulty to analyze the exactly same spot of the first measurement, thus probing a slightly different silicene area. The 48 h spectrum clearly shows again the two main A and E modes of silicene, without visible Raman shift. The spectrum details in the region between the two main

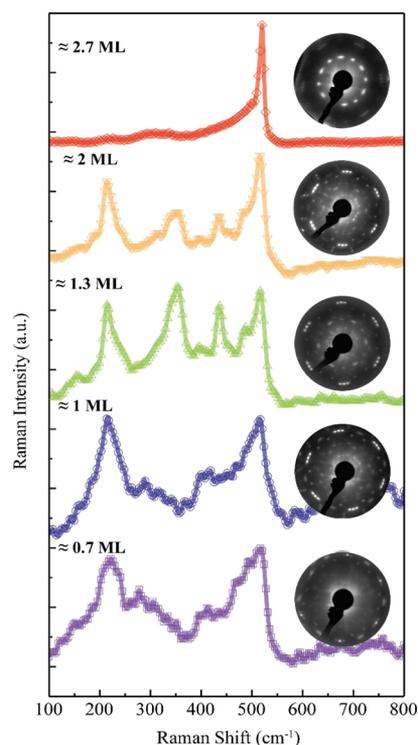


**Figure 4.** Time evolution of the Raman spectrum of encapsulated silicene exposed to ambient conditions for 1 h (red), 24 h (blue), 48 h (yellow), and 72 h (green).

modes are still similar to what was observed after 1 h of exposure to ambient conditions. A general decrease in intensity, alongside with a rounding of the A mode, can be noticed, hinting to a start of the degradation process. After 72 h of exposure to ambient conditions, no traces of the presence of silicene can be detected, demonstrating a complete degradation of the 2D layer. It must be noted that all Raman spectra acquired outside of an encapsulation area resemble the green spectrum of Figure 4, even few minutes after the removal of the sample from the UHV chamber, thus proving the capping capabilities of the FLG layer. Finally, it must be taken into consideration that the measurement process can locally heat the FLG flake.<sup>46</sup> For this reason, a repeated measurement could speed up the degradation process, as the contact between silicene and the capping flake deteriorates because of thermal shrinking and consequent expansion of the FLG layer.<sup>47</sup> To prove these assumptions, we examined a sample with two FLG flakes successfully passivating silicene. One flake was analyzed approximately after 1 h of exposure to ambient conditions, whereas the second flake was measured only once after 36 h. As provided in the Supporting Information, the two spectra show the same Raman intensities. Thus, we conclude that the Raman measurement itself accelerates the degradation.

The passivation method was applied to study the Raman spectra of samples characterized by different Si coverages ( $\Theta$ ). By varying the evaporation time, it was possible to obtain samples with  $\Theta$  of  $\approx 0.7$  ML,  $\approx 1$  ML,  $\approx 1.3$  ML,  $\approx 2$  ML, and  $\approx 2.7$  ML. Each sample was then passivated following the method described above, and the Raman spectra were acquired ex situ within 1 h after the removal from the UHV chamber.

Figure 5 shows the comparison between the normalized spectra of the respective samples. The samples with  $\Theta \leq 1$  ML show a Raman spectrum mainly dominated by the A and E phonon modes, as previously discussed. The respective LEED patterns in the insets show the typical reconstruction for monolayer silicene grown on Ag(111) at 530 K with a prominent signal from the  $4 \times 4$  phase mixed with a low signal for the minority  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  reconstruction.<sup>35</sup> The LEED pattern for the sample with  $\Theta \approx 1$  ML is clearly more intense than the pattern observed for the sample with  $\Theta \approx 0.7$  ML because of the presence of larger silicene domains.



**Figure 5.** Normalized Raman spectra and LEED pattern of samples characterized by different silicon coverages ( $\Theta$ ). From bottom to top: spectrum of  $\approx 0.7$  ML silicene (purple) characterized by the two main modes A and E of monolayer silicene, with respective LEED picture showing a  $4 \times 4$  silicene reconstruction with a minority  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  phase; spectrum of  $\approx 1$  ML silicene (blue), with respective LEED picture showing a high-intensity signal from a  $4 \times 4$  silicene reconstruction with a minority  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  phase; spectrum of  $\approx 1.3$  ML silicene (green) characterized by the appearance of additional peaks between the A and E modes, with respective LEED picture showing a mixture of  $4 \times 4$  and  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  reconstruction because of the monolayer regions, and  $\frac{4}{\sqrt{3}} \times \frac{4}{\sqrt{3}}R30^\circ$ , attributed to the presence of regions where  $\Theta$  exceeds 1 ML; spectrum of  $\approx 2$  ML silicene (orange) characterized by a lower intensity of the additional peaks between the A and E modes, with respective LEED picture showing a mixture of  $4 \times 4$ ,  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$ , and  $\frac{4}{\sqrt{3}} \times \frac{4}{\sqrt{3}}R30^\circ$ ; spectrum of  $\approx 2.7$  ML silicene (red) characterized by a high-intensity peak located at  $\approx 520$  cm<sup>-1</sup>, with respective LEED picture showing an intense signal from the  $\frac{4}{\sqrt{3}} \times \frac{4}{\sqrt{3}}R30^\circ$  phase, together with a faint signal from the  $4 \times 4$  and  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  reconstructions.

When  $\Theta$  exceeds 1 ML, the Raman spectrum clearly changes. While A and E modes are still visible, two intense and sharp peaks appear around  $\approx 350$  and  $\approx 450$  cm<sup>-1</sup>, respectively. Similar Raman features were previously reported by an in situ Raman study for multilayer silicene samples.<sup>18</sup> These additional Raman peaks were tentatively assigned to edge-activated phonon modes, similar to that observed in graphene flakes, in which defects are responsible for double-resonating Raman scattering processes.<sup>48</sup> The LEED pattern changes, now showing a mixture of  $4 \times 4$  and  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  silicene, because of the monolayer regions, and  $\frac{4}{\sqrt{3}} \times \frac{4}{\sqrt{3}}R30^\circ$ , attributed to the presence of regions where  $\Theta$  exceeds 1 ML.<sup>49,50</sup> When  $\Theta$  is further increased toward 2 ML, the additional Raman peaks are still visible, but their

intensity is lower with respect to the intensity of the A and E modes. The LEED pattern still shows a mixed reconstruction of  $4 \times 4$ ,  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$ , and  $\frac{4}{\sqrt{3}} \times \frac{4}{\sqrt{3}}R30^\circ$  phases. It can be assumed that the completion of the second layer leads to a reduction in the density of domain edges. This observation supports the identification of the new Raman modes as defect-activated. It must be noted, however, that the TERS experiment clearly shows that the edges of the first silicene layer are not activating any of these additional Raman modes,<sup>17</sup> thus being related only to the edges of the topmost layer. As the existence of multilayer silicene is still under debate,<sup>50–52</sup> one cannot exclude the possibility that these additional Raman modes are, instead, appearing because of a change in the crystalline structure of the analyzed layer, promoted by the higher density of Si atoms.

The sample with  $\Theta \approx 2.7$  ML is characterized by a drastically different Raman spectrum and LEED pattern. The spectrum shows a single Raman peak, centered at  $\approx 520$   $\text{cm}^{-1}$ , with a much higher intensity than the silicene A and E peaks (50x more intense in the non-normalized data). Remarkably, the same spectrum can also be detected by analyzing areas that are outside of the capping flake, demonstrating the great air stability of the layer. The LEED pattern changed as well. The pattern is now prevalently  $\frac{4}{\sqrt{3}} \times \frac{4}{\sqrt{3}}R30^\circ$ , mixed with a very low-intensity  $4 \times 4$  and  $(\sqrt{13} \times \sqrt{13})R13.9^\circ$  pattern. It has been shown that this peculiar LEED pattern could be related to a surface reconstruction of Ag atoms that floated on top of the Si layer during the evaporation process.<sup>50,53</sup>

In situ Raman studies of multilayer silicene have reported a slightly different spectrum than the one shown here, characterized by an intense peak located at  $\approx 523$   $\text{cm}^{-1}$  and a broad band located between 430 and 500  $\text{cm}^{-1}$ .<sup>54</sup> The  $\approx 3$   $\text{cm}^{-1}$  shift of the main peak can be attributed to the presence of stacking faults that cause the in-plane lattice constant to be smaller.<sup>53</sup> The increased presence of stacking faults in the mentioned in situ study can be explained by the choice of a particularly low growth temperature (470 K), with respect to the one of the present works.<sup>52</sup> The broad band, instead, can be assigned to the presence of amorphous silicon.

## 4. CONCLUSIONS

In conclusion, we present a novel method for the stabilization of silicene on Ag(111) up to 48 h. In UHV, an FLG flake is mechanically exfoliated atop of epitaxially grown silicene, forming an encapsulation layer. This allows to perform high-quality, spatially resolved, ex situ Raman analysis of the silicene layer, thanks to the Raman inactivity of the FLG flake in the spectral region of interest. The observed vibrational signature matches with the in situ Raman spectra reported in the literature, proving that the encapsulation layer does not alter the overall structure of silicene. The method has been applied to study the Raman signature of silicene samples characterized by a Si coverage spanning between 0.7 and 2.7 ML. The present work also confirms the compatibility between graphene and silicene, representing a step forward toward the integration of silicene into 2D heterostructures.

## ■ ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.8b20751.

Raman spectra showing the influence of the measurement process on the degradation of silicene under a FLG flake (PDF)

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### Author Contributions

V.R., J.G. and D.N. contributed equally to this work.

### Notes

The authors declare no competing financial interest.

## ■ ACKNOWLEDGMENTS

This work was funded by the Fonds zur Förderung der Wissenschaftlichen Forschung (FWF), Austria (project P29244-N27). We thank the center for micro- and nanostructures (ZMNS) at the TU Vienna for access to the cleanroom facilities. We would like to thank Harold J. W. Zandvliet from the University of Twente for support in the synthesis of silicene and Alessandro Molle from the Institute for Microelectronics and Microsystems for the fruitful discussions on the encapsulation process.

## ■ REFERENCES

- (1) Tao, L.; Cinquanta, E.; Chiappe, D.; Grazianetti, C.; Fanciulli, M.; Dubey, M.; Molle, A.; Akinwande, D. Silicene Field-Effect Transistors Operating at Room Temperature. *Nat. Nanotechnol.* **2015**, *10*, 227–231.
- (2) Le Lay, G. Silicene Transistors. *Nat. Nanotechnol.* **2015**, *10*, 202–203.
- (3) Zhao, J.; Liu, H.; Yu, Z.; Quhe, R.; Zhou, S.; Wang, Y.; Liu, C. C.; Zhong, H.; Han, N.; Lu, J.; Yao, Y.; Wu, K. Rise of Silicene: A Competitive 2D Material. *Prog. Mater. Sci.* **2016**, *83*, 24–151.
- (4) Voon, L. C. L. Y. Electronic Structure of Silicene. *Chin. Phys. B* **2015**, *24*, 087309.
- (5) Drummond, N. D.; Zólyomi, V.; Fal'ko, V. I. Electrically Tunable Band Gap in Silicene. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2012**, *85*, 075423.
- (6) Ni, Z.; Liu, Q.; Tang, K.; Zheng, J.; Zhou, J.; Qin, R.; Gao, Z.; Yu, D.; Lu, J. Tunable Bandgap in Silicene and Germanene. *Nano Lett.* **2012**, *12*, 113–118.
- (7) Gao, N.; Zheng, W. T.; Jiang, Q. Density Functional Theory Calculations for Two-Dimensional Silicene with Halogen Functionalization. *Phys. Chem. Chem. Phys.* **2012**, *14*, 257–261.
- (8) Dávila, M. E.; Marele, A.; De Padova, P.; Montero, I.; Hennies, F.; Pietzsch, A.; Shariati, M. N.; Gómez-Rodríguez, J. M.; Le Lay, G. Comparative Structural and Electronic Studies of Hydrogen Interaction with Isolated versus Ordered Silicon Nanoribbons Grown on Ag(110). *Nanotechnology* **2012**, *23*, 385703.
- (9) Vogt, P.; De Padova, P.; Quaresima, C.; Avila, J.; Frantzeskakis, E.; Asensio, M. C.; Resta, A.; Ealet, B.; Le Lay, G. Silicene: Compelling Experimental Evidence for Graphenelike Two-Dimensional Silicon. *Phys. Rev. Lett.* **2012**, *108*, 155501.
- (10) Chun-Liang, L.; Ryuichi, A.; Maki, K.; Noriaki, T. Comparison of Electronic Structure between Monolayer Silicenes on Ag (111). *Chin. Phys. B* **2015**, *24*, 087307.
- (11) Friedlein, R.; Van Bui, H.; Wiggers, F. B.; Yamada-Takamura, Y.; Kovalgin, A. Y.; de Jong, M. P. Interaction of Epitaxial Silicene with Overlayers Formed by Exposure to Al Atoms and O<sub>2</sub> Molecules. *J. Chem. Phys.* **2014**, *140*, 204705.

- (12) Cahangirov, S.; Sahin, H.; Le Lay, G.; Rubio, A. *Freestanding Silicene. In Introduction to the Physics of Silicene and Other 2D Materials*; Springer International Publishing, 2017.
- (13) Cinquanta, E.; Scalise, E.; Chiappe, D.; Grazianetti, C.; van den Broek, B.; Houssa, M.; Fanciulli, M.; Molle, A. Getting through the Nature of Silicene: An sp<sup>2</sup>–sp<sup>3</sup> Two-Dimensional Silicon Nanosheet. *J. Phys. Chem. C* **2013**, *117*, 16719–16724.
- (14) Jose, D.; Nijamudheen, A.; Datta, A. Tip Enhanced Raman Spectroscopy (TERS) as a Probe for the Buckling Distortion in Silicene. *Phys. Chem. Chem. Phys.* **2013**, *15*, 8700.
- (15) Solonenko, D.; Gordan, O. D.; Lay, G. L.; Şahin, H.; Cahangirov, S.; Zahn, D. R. T.; Vogt, P. 2D Vibrational Properties of Epitaxial Silicene on Ag(111). *2D Mater.* **2016**, *4*, 015008.
- (16) Yan, J.-A.; Stein, R.; Schaefer, D. M.; Wang, X.-Q.; Chou, M. Y. Electron-Phonon Coupling in Two-Dimensional Silicene and Germanene. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2013**, *88*, 121403.
- (17) Sheng, S.; Wu, J.; Cong, X.; Li, W.; Gou, J.; Zhong, Q.; Cheng, P.; Tan, P.; Chen, L.; Wu, K. Vibrational Properties of a Monolayer Silicene Sheet Studied by Tip-Enhanced Raman Spectroscopy. *Phys. Rev. Lett.* **2017**, *119*, 196803.
- (18) Zhuang, J.; Xu, X.; Du, Y.; Wu, K.; Chen, L.; Hao, W.; Wang, J.; Yeoh, W. K.; Wang, X.; Dou, S. X. Investigation of Electron-Phonon Coupling in Epitaxial Silicene by In Situ Raman Spectroscopy. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2015**, *91*, 161409.
- (19) Meric, I.; Dean, C. R.; Petrone, N.; Wang, L.; Hone, J.; Kim, P.; Shepard, K. L. Graphene Field-Effect Transistors Based on Boron–Nitride Dielectrics. *Proc. IEEE* **2013**, *101*, 1609–1619.
- (20) Akinwande, D.; Petrone, N.; Hone, J. Two-Dimensional Flexible Nanoelectronics. *Nat. Commun.* **2014**, *5*, 5678.
- (21) Bunch, J. S.; Verbridge, S. S.; Alden, J. S.; van der Zande, A. M.; Parpia, J. M.; Craighead, H. G.; McEuen, P. L. Impermeable Atomic Membranes from Graphene Sheets. *Nano Lett.* **2008**, *8*, 2458–2462.
- (22) Novoselov, K. S.; Geim, A. K.; Morozov, S. V.; Jiang, D. A.; Zhang, Y.; Dubonos, S. V.; Grigorieva, I. V.; Firsov, A. A. Electric Field Effect in Atomically Thin Carbon Films. *Science* **2004**, *306*, 666–669.
- (23) Blake, P.; Hill, E. W.; Castro Neto, A. H.; Novoselov, K. S.; Jiang, D.; Yang, R.; Booth, T. J.; Geim, A. K. Making Graphene Visible. *Appl. Phys. Lett.* **2007**, *91*, 063124.
- (24) Gorbachev, R. V.; Riaz, I.; Nair, R. R.; Jalil, R.; Britnell, L.; Belle, B. D.; Hill, E. W.; Novoselov, K. S.; Watanabe, K.; Taniguchi, T.; Geim, A. K.; Blake, P. Hunting for Monolayer Boron Nitride: Optical and Raman Signatures. *Small* **2011**, *7*, 465–468.
- (25) Velický, M.; Hendren, W. R.; Donnelly, G. E.; Katzen, J. M.; Bowman, R. M.; Huang, F. Optimising the Visibility of Graphene and Graphene Oxide on Gold with Multilayer Heterostructures. *Nanotechnology* **2018**, *29*, 275205.
- (26) Kim, D.; Song, S.; Choi, B.-Y.; Chang, Y. J.; Choi, E. J. Long-Term Stability Study of Graphene-Passivated Black Phosphorus under Air Exposure. *Curr. Appl. Phys.* **2016**, *16*, 165–169.
- (27) Doganov, R. A.; O'Farrell, E. C. T.; Koenig, S. P.; Yeo, Y.; Ziletti, A.; Carvalho, A.; Campbell, D. K.; Coker, D. F.; Watanabe, K.; Taniguchi, T.; Neto, A. H. C.; Özyilmaz, B. Transport Properties of Pristine Few-Layer Black Phosphorus by van Der Waals Passivation in an Inert Atmosphere. *Nat. Commun.* **2015**, *6*, 6647.
- (28) Dedkov, Y. S.; Fonin, M.; Laubschat, C. A Possible Source of Spin-Polarized Electrons: The Inert Graphene/Ni(111) System. *Appl. Phys. Lett.* **2008**, *92*, 052506.
- (29) Chen, S.; Brown, L.; Levendorf, M.; Cai, W.; Ju, S.-Y.; Edgeworth, J.; Li, X.; Magnuson, C. W.; Velamakanni, A.; Piner, R. D.; Kang, J.; Park, J.; Ruoff, R. S. Oxidation Resistance of Graphene-Coated Cu and Cu/Ni Alloy. *ACS Nano* **2011**, *5*, 1321–1327.
- (30) Du, Y.; Zhuang, J.; Liu, H.; Xu, X.; Eilers, S.; Wu, K.; Cheng, P.; Zhao, J.; Pi, X.; See, K. W.; Pelecks, G.; Wang, X.; Dou, S. X. Tuning the Band Gap in Silicene by Oxidation. *ACS Nano* **2014**, *8*, 10019–10025.
- (31) Ali, M.; Ni, Z.; Cottenier, S.; Liu, Y.; Pi, X.; Yang, D. Formation, Structures and Electronic Properties of Silicene Oxides on Ag(111). *J. Mater. Sci. Technol.* **2017**, *33*, 751–757.
- (32) Island, J. O.; Steele, G. A.; van der Zant, H. S. J.; Castellanos-Gomez, A. Environmental Instability of Few-Layer Black Phosphorus. *2D Mater.* **2015**, *2*, 011002.
- (33) Li, L.; Kim, J.; Jin, C.; Ye, G. J.; Qiu, D. Y.; da Jornada, F. H.; Shi, Z.; Chen, L.; Zhang, Z.; Yang, F.; Watanabe, K.; Taniguchi, T.; Ren, W.; Louie, S. G.; Chen, X. H.; Zhang, Y.; Wang, F. Direct Observation of the Layer-Dependent Electronic Structure in Phosphorene. *Nat. Nanotechnol.* **2016**, *12*, 21–25.
- (34) Pflugradt, P.; Matthes, L.; Bechstedt, F. Silicene-Derived Phases on Ag(111) Substrate versus Coverage: Ab Initio Studies. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2014**, *89*, 035403.
- (35) Moras, P.; Mentess, T. O.; Sheverdyayeva, P. M.; Locatelli, A.; Carbone, C. Coexistence of Multiple Silicene Phases in Silicon Grown on Ag(111). *J. Phys.: Condens. Matter* **2014**, *26*, 185001.
- (36) Ferrari, A. C.; Meyer, J. C.; Scardaci, V.; Casiraghi, C.; Lazzeri, M.; Mauri, F.; Piscanec, S.; Jiang, D.; Novoselov, K. S.; Roth, S.; Geim, A. K. Raman Spectrum of Graphene and Graphene Layers. *Phys. Rev. Lett.* **2006**, *97*, 187401.
- (37) Malard, L. M.; Pimenta, M. A.; Dresselhaus, G.; Dresselhaus, M. S. Raman Spectroscopy in Graphene. *Phys. Rep.* **2009**, *473*, 51–87.
- (38) Molle, A.; Grazianetti, C.; Chiappe, D.; Cinquanta, E.; Cianci, E.; Tallarida, G.; Fanciulli, M. Hindering the Oxidation of Silicene with Non-Reactive Encapsulation. *Adv. Funct. Mater.* **2013**, *23*, 4340–4344.
- (39) Ribeiro-Soares, J.; Almeida, R. M.; Cançado, L. G.; Dresselhaus, M. S.; Jorio, A. Group Theory for Structural Analysis and Lattice Vibrations in Phosphorene Systems. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2015**, *91*, 205421.
- (40) Dresselhaus, M. S.; Dresselhaus, G.; Jorio, A. (Eds.). *Group Theory: Application to the Physics of Condensed Matter*; Springer-Verlag, 2008.
- (41) Harris, D. C.; Bertolucci, M. D. *Symmetry and Spectroscopy: An Introduction to Vibrational and Electronic Spectroscopy*; Dover Publications, 1989.
- (42) Richter, H.; Wang, Z. P.; Ley, L. The One Phonon Raman Spectrum in Microcrystalline Silicon. *Solid State Commun.* **1981**, *39*, 625–629.
- (43) Houssa, M.; Pourtois, G.; Afanas'ev, V. V.; Stesmans, A. Can Silicon Behave like Graphene? A First-Principles Study. *Appl. Phys. Lett.* **2010**, *97*, 112106.
- (44) Lin, C.-L.; Arafune, R.; Kawahara, K.; Kanno, M.; Tsukahara, N.; Minamitani, E.; Kim, Y.; Kawai, M.; Takagi, N. Substrate-Induced Symmetry Breaking in Silicene. *Phys. Rev. Lett.* **2013**, *110*, 076801.
- (45) Liu, C.-C.; Feng, W.; Yao, Y. Quantum Spin Hall Effect in Silicene and Two-Dimensional Germanium. *Phys. Rev. Lett.* **2011**, *107*, 076802.
- (46) Everall, N. J.; Lumsdon, J.; Christopher, D. J. The Effect of Laser-Induced Heating upon the Vibrational Raman Spectra of Graphites and Carbon Fibres. *Carbon* **1991**, *29*, 133–137.
- (47) Yoon, D.; Son, Y.-W.; Cheong, H. Negative Thermal Expansion Coefficient of Graphene Measured by Raman Spectroscopy. *Nano Lett.* **2011**, *11*, 3227–3231.
- (48) Basko, D. M. Theory of Resonant Multiphonon Raman Scattering in Graphene. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2008**, *78*, 125418.
- (49) Salomon, E.; El Ajjouri, R.; Lay, G. L.; Angot, T. Growth and Structural Properties of Silicene at Multilayer Coverage. *J. Phys.: Condens. Matter* **2014**, *26*, 185003.
- (50) Shirai, T.; Shirasawa, T.; Hirahara, T.; Fukui, N.; Takahashi, T.; Hasegawa, S. Structure Determination of Multilayer Silicene Grown on Ag(111) Films by Electron Diffraction: Evidence for Ag Segregation at the Surface. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2014**, *89*, 241403.
- (51) Kamal, C.; Chakrabarti, A.; Banerjee, A.; Deb, S. K. Silicene beyond Mono-Layers—different Stacking Configurations and Their Properties. *J. Phys.: Condens. Matter* **2013**, *25*, 085508.

(52) De Padova, P.; Generosi, A.; Paci, B.; Ottaviani, C.; Quaresima, C.; Olivieri, B.; Salomon, E.; Angot, T.; Le Lay, G. Multilayer Silicene: Clear Evidence. *2D Mater.* **2016**, *3*, 031011.

(53) Curcella, A.; Bernard, R.; Borensztein, Y.; Lazzeri, M.; Resta, A.; Garreau, Y.; Prévot, G. Multilayer Silicene: Clear Evidence of Ag-Terminated Bulk Silicon. *2D Mater.* **2017**, *4*, 025067.

(54) De Padova, P.; Ottaviani, C.; Quaresima, C.; Olivieri, B.; Imperatori, P.; Salomon, E.; Angot, T.; Quagliano, L.; Romano, C.; Vona, A.; Muniz-Miranda, M.; Generosi, A.; Paci, B.; Le Lay, G. 24 h Stability of Thick Multilayer Silicene in Air. *2D Mater.* **2014**, *1*, 021003.